

On the largest eigenvalue of a Hermitian random matrix model with spiked external source I. Rank one case

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Abstract

Consider a Hermitian matrix model under an external potential with spiked external source. When the external source is of rank one, we compute the limiting distribution of the largest eigenvalue for general, regular, analytic potential for all values of the external source. There is a transitional phenomenon, which is universal for convex potentials. However, for non-convex potentials, new types of transition may occur. The higher rank external source is analyzed in the subsequent paper.

1 Introduction and results

1.1 Introduction

Fix an $n \times n$ Hermitian matrix \mathbf{A}_n and consider the following density function on the set \mathcal{H}_n of $n \times n$ Hermitian matrices:

$$p_n(M) = \frac{1}{Z_n} e^{-n \operatorname{Tr}(V(M) - \mathbf{A}_n M)} \quad (1)$$

where Z_n is the normalization constant. Here the ‘external potential’ $V(x)$ is a real-valued function which decays fast enough as $|x| \rightarrow \infty$ so that Z_n is convergent. The matrix \mathbf{A}_n is called the external source: see e.g. [11], [12], [26], [27], [9], [10], [2]. Note that the distribution of eigenvalues of M is unchanged if \mathbf{A}_n is replaced by $U\mathbf{A}_n U^{-1}$ for any unitary matrix U . Since we are only concerned on eigenvalues of M , we assume without loss of generality that \mathbf{A}_n is a diagonal matrix.

A special case is when for all n , the external source has a fixed number \mathbf{m} , called the rank of \mathbf{A}_n , of fixed non-zero eigenvalues. In this case, the sequence of probability spaces (\mathcal{H}_n, p_n) is called a Hermitian matrix model with spiked external source, spiked source model for short. In this paper we only consider the case when $\mathbf{m} = 1$. The higher rank case when $\mathbf{m} > 1$ will be analyzed in the

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upcoming companion paper. Throughout this paper, we assume that $n \geq 1$ and

$$\mathbf{A}_n = \text{diag}(a, \underbrace{0, \dots, 0}_{n-1}), \quad (2)$$

where a is a real number, independent of n .

There are two important special cases. When $V(x) = x^2/2$, the spiked source model is called the GUE spiked model. The density $p_n(M)$ is that of $M = H + \mathbf{A}_n$ where H is an $n \times n$ GUE (Gaussian unitary ensemble) matrix. When $V(x) = ((1+c)x - c \log x)\chi_{(0,\infty)}(x)$, $c = (m-n)/n \geq 0$, the spiked source model is the complex Wishart spiked model. In this case, setting $\Sigma := (1 - (1+c)^{-1}\mathbf{A}_n)^{-1}$, the density $p_n(M)$ is that of $M = \Sigma^{1/2} X X^\dagger \Sigma^{1/2}$ where X is an $n \times m$ complex rectangular matrix with i.i.d. standard complex Gaussian entries.¹ For these two cases, the limit of the largest eigenvalue $\xi_{\max}(n)$ of M was studied in great detail in [3] and [23]. An important feature is the following phase transition phenomenon. Let \mathbf{e} denote the right-end point of the limiting empirical distribution of the eigenvalues of the Hermitian matrix model with no external source (see (10) below).² It was shown in both the GUE and the complex Wishart spiked models that as $n \rightarrow \infty$, with probability 1,

$$\xi_{\max}(n) \rightarrow \begin{cases} \mathbf{e}, & \text{if } a \leq \frac{1}{2}V'(\mathbf{e}), \\ x_0(a), & \text{if } a > \frac{1}{2}V'(\mathbf{e}), \end{cases} \quad (3)$$

for some continuous, increasing function $x_0(a)$ in $a \in (\frac{1}{2}V'(\mathbf{e}), \infty)$ satisfying $\lim_{a \downarrow \frac{1}{2}V'(\mathbf{e})} x_0(a) = \mathbf{e}$. Moreover, there exists $\beta > 0$ (see (22) below) such that for each $T \in \mathbb{R}$,

$$\mathbb{P}_n((\xi_{\max}(n) - \mathbf{e})\beta n^{2/3} \leq T) \rightarrow \begin{cases} F_0(T), & \text{if } a < \frac{1}{2}V'(\mathbf{e}), \\ F_1(T), & \text{if } a = \frac{1}{2}V'(\mathbf{e}), \end{cases} \quad (4)$$

and there exists $\gamma(a)$ such that for each $T \in \mathbb{R}$,

$$\mathbb{P}_n((\xi_{\max}(n) - x_0(a))\gamma(a)n^{1/2} \leq T) \rightarrow G(T), \quad \text{if } a > \frac{1}{2}V'(\mathbf{e}). \quad (5)$$

Here the function $G(T) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^T e^{-\frac{1}{2}\xi^2} d\xi$ is the cumulative distribution function of the standard normal distribution, and F_0 and $\sqrt{F_1}$ are the GUE and GOE Tracy-Widom distribution functions, respectively. They are defined in (16) and (20) below, respectively. A limit theorem was also proven for the double scaling case when $a = \frac{1}{2}V'(\mathbf{e}) + \frac{\alpha}{n^{1/3}}$.

The purpose of this paper is to extend the results (3)–(5) to the spiked source model with general potential V . It turns out that if $V(x)$ is convex in the interval $x \in (\mathbf{e}, \infty)$, then all of (3)–(5) still hold. Especially, the ‘critical value’ of a is again given by $\frac{1}{2}V'(\mathbf{e})$. However, if V is not convex in (\mathbf{e}, ∞) , new features may occur. Two key new features are the followings.

¹For the complex Wishart spiked model, V is not real analytic at $x = 0$. Throughout this paper, we only consider V which is real analytic in the whole line. However, the method can be generalized to the Wishart-type potentials in a straightforward way.

²The limiting empirical distribution in the spiked source model is the same as the Hermitian model with no external source.

- The critical value of a may be smaller than $\frac{1}{2}V'(\mathbf{e})$. See Lemma 1.2 and Theorem 1.2. For such a case, when a equals this critical value, $\xi_{\max}(n)$ does not converge with probability 1. Instead it converges to two or more values, each with non-zero probability. In this case, the fluctuation of $\xi_{\max}(n)$ is generically F_0 at the smallest limiting value and G at the larger limiting values. See Theorem 1.3.
- There may be a discrete set of ‘secondary critical values’ of a , which are greater than the critical value. If a is at a secondary critical value, then $\xi_{\max}(n)$ converge to two or more values, each with non-zero probability. In this case, the fluctuation of $\xi_{\max}(n)$ is generically G at each of the limiting values. See Theorem 1.4.

The exact assumptions on the potential V is given in Subsection 1.2. The universality result for convex potentials is in Subsection 1.3. In Subsection 1.4 we define the critical and the secondary critical values for non-convex potentials. The limit laws for the non-convex potentials are given in Subsection 1.5.

While we were preparing for this paper and the companion paper for the higher rank case, we learned that M. Bertola, R. Buckingham, S. Y. Lee and V. Pierce were also working on the spiked source models (see [8] for the first part of their work). While we focus, especially in the second paper, on the limit laws when $a_1, \dots, a_{\mathbf{m}}$ are distinct, Bertola, Buckingham, Lee and Pierce focus on the case when $a_1 = \dots = a_{\mathbf{m}}$ and $\mathbf{m} \rightarrow \infty$ slower than n . Also we use the asymptotics of usual orthogonal polynomials but Bertola, Buckingham, Lee and Pierce use asymptotics of multiple orthogonal polynomials via Riemann-Hilbert problem of size larger than 2.

Before closing this subsection, we mention that the spiked real symmetric matrix model is much more difficult. Even for the GOE and the real Wishart case, the limiting distribution at the critical value is not yet known. For the quaternionic case, the limiting distribution is obtained when the rank $\mathbf{m} = 1$ (see [25] for the Wishart model; Gaussian model is also similar).

We also mention that there are several results for the spiked Wigner ensembles and spiked sample covariance matrices. See, for example, [5], [22], [19], [13], [21], [7] and [6].

1.2 Assumptions on external potential V .

Throughout this paper, we assume the following three conditions on V :

$$V(x) \text{ is real analytic in } \mathbb{R}, \tag{6}$$

$$\frac{V(x)}{\sqrt{x^2 + 1}} \rightarrow +\infty \text{ as } |x| \rightarrow \infty, \tag{7}$$

$$V \text{ is ‘regular’}. \tag{8}$$

The second condition is to ensure the convergence of the density function: compare this with the condition on V in [17]. The third condition on being ‘regular’ is a technical condition as defined in [17]. We need a few definitions to state it.

First, recall the equilibrium measure and the so-called \mathbf{g} -function. General references are [24] and [15]. For a given potential V , the empirical distribution of the eigenvalues of the matrix model with no external source converges to the associated equilibrium measure μ . The equilibrium measure is characterized by a certain variational problem. If V is real analytic, μ is supported on a finite union of intervals,

$$J = \bigcup_{j=0}^N (b_j, a_{j+1}), \quad \text{with } b_0 < a_1 < b_1 < \cdots < a_{N+1}, \quad (9)$$

for some $N \geq 0$. We denote the right-most edge of the support by

$$\mathbf{e} := a_{N+1}. \quad (10)$$

On J , $d\mu$ has the form $d\mu = \Psi(x)dx$,

$$\Psi(x) = \frac{1}{2\pi i} R_+^{1/2}(x)h(x), \quad \text{where } R(z) = \prod_{j=0}^N (z - b_j)(z - a_{j+1}). \quad (11)$$

The function $R(z)^{1/2}$ is defined to be analytic in $\mathbb{C} \setminus J$ and satisfy $R(z)^{1/2} \sim z^N$ as $z \rightarrow \infty$. The notation $R_+^{1/2}(x)$ for $x \in J$ denotes the limit of $R^{1/2}(z)$, $z \in \mathbb{C}_+$, as $z \rightarrow x$ from above. The function $h(x)$ is real analytic in \mathbb{R} and is given by [17, Formula (3.18)].

The equilibrium measure, $d\mu(x) = \Psi(x)dx$, is characterized by the following conditions: there is a constant (called the Robin constant) ℓ such that

$$2 \int_J \log|x - s| \Psi(s) ds - V(x) = \ell \quad \text{for } x \in \bar{J}, \quad (12)$$

$$2 \int_J \log|x - s| \Psi(s) ds - V(x) \leq \ell \quad \text{for } x \in \mathbb{R} \setminus \bar{J}. \quad (13)$$

The so-called \mathbf{g} -function is defined by

$$\mathbf{g}(z) := \int_J \log(z - s) \Psi(s) ds, \quad \text{for } z \in \mathbb{C} \setminus (-\infty, \mathbf{e}). \quad (14)$$

The potential V is said to be regular (see [17]) if

- $h(x) \neq 0$ for $x \in \bar{J}$,
- the inequality in (13) is strict.

The first condition implies that the function $\Psi(x) > 0$ for all $x \in J$, and also that $\Psi(x)$ vanishes like a square-root at each end of the interval of the support. This in turn implies, in particular, that for the model with $\mathbf{A}_n = 0$, the largest eigenvalue has the limiting distribution given by F_0 (see e.g. [14] for the non-varying weight; varying weight case is similar using the analysis of [17].) Note that the second condition restricted to the domain $x > \mathbf{e}$ implies that

$$2\mathbf{g}(x) - V(x) < \ell, \quad x > \mathbf{e}. \quad (15)$$

We will use this fact later.

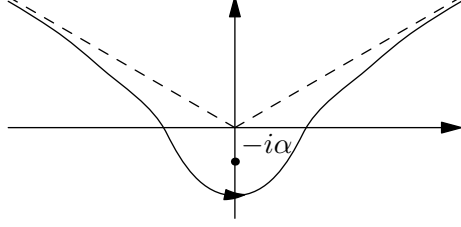


Figure 1: The contour from $\infty e^{5\pi i/6}$ to $\infty e^{\pi i/6}$.

1.3 Statement of results: convex potentials

Let $F_0(T)$ be the GUE Tracy-Widom distribution defined by

$$F_0(T) := \det(1 - \chi_{[T, \infty)} K_{\text{Airy}} \chi_{[T, \infty)}), \quad (16)$$

where $\chi_{[T, \infty)}$ denotes the projection operator on $[T, \infty)$, and K_{Airy} is the Airy operator defined by the kernel

$$K_{\text{Airy}}(x, y) = \frac{\text{Ai}(x) \text{Ai}'(y) - \text{Ai}'(x) \text{Ai}(y)}{x - y}. \quad (17)$$

Here Ai is the Airy function.

For $\alpha \in \mathbb{R}$, define the function

$$C_\alpha(\xi) := \frac{1}{2\pi} \int e^{i\frac{1}{3}z^3 + i\xi z} \frac{dz}{\alpha + iz}, \quad (18)$$

where the contour is from $\infty e^{5\pi i/6}$ to $\infty e^{\pi i/6}$ and the pole $z = -i\alpha$ lies above the contour in the complex plane: see Figure 1.

Define

$$F_1(T; \alpha) := F_0(T) \cdot \left(1 - \langle (1 - \chi_{[T, \infty)} K_{\text{Airy}} \chi_{[T, \infty)})^{-1} C_\alpha, \text{Ai} \rangle_{[T, \infty)} \right), \quad (19)$$

where $\langle f, g \rangle_E$ denotes the real inner product over E , $\int_E f(x)g(x)dx$. (See [3, Definition 1.3].) When $\alpha = 0$,

$$F_1(T) := F_1(T; 0) \quad (20)$$

equals the square of the GOE Tracy-Widom distribution (see [3, Formula (24)]).

Fix a potential V satisfying the assumptions (6)–(8). In the companion paper on the higher rank case, we need to consider the spiked source model of rank 1 whose density function is same as in (1) but with the change that the matrix M is now of size $n - j + 1$ and \mathbf{A}_n is replaced by \mathbf{A}_{n-j+1} for fixed j :

$$\frac{1}{Z_{n-j+1, n}} e^{-n \text{Tr}(V(M) - \mathbf{A}_{n-j+1} M)}, \quad M \in \mathcal{H}_{n-j+1}. \quad (21)$$

Note that the factor n in front of the potential is unchanged. For a subset $E \subset \mathbb{R}$, let $\mathbb{P}_{n-j+1, n}(a; E)$ denote the gap probability that there are no eigenvalues of M in the set E , where a represents the unique non-zero eigenvalue of \mathbf{A}_{n-j+1} . Hence $\mathbb{P}_{n-j+1, n}(a; [t, \infty))$ is the probability that the largest eigenvalue of M is less than t .

Let \mathbf{e} be as in (10). Set (recall (11))

$$\beta := \left(\lim_{x \uparrow \mathbf{e}} \frac{\pi \Psi(x)}{\sqrt{\mathbf{e} - x}} \right)^{2/3} = \left(\frac{h(\mathbf{e})}{2} \right)^{2/3} \left(\frac{R(z)}{z - \mathbf{e}} \right)^{1/3} \Big|_{z=\mathbf{e}} \quad (22)$$

so that $\Psi(x) \sim \frac{\beta^{3/2}}{\pi} \sqrt{\mathbf{e} - x}$ for $x \uparrow \mathbf{e}$. For $T \in \mathbb{R}$, define the intervals

$$I_n^T := \left[\mathbf{e} + \frac{T}{\beta n^{2/3}}, \infty \right) \quad (23)$$

and

$$J_n^T(x_*) := \left[x_* + \frac{T}{\sqrt{(V''(x_*) - \mathbf{g}''(x_*))n}}, \infty \right) \quad (24)$$

for $x_* > \mathbf{e}$, if $V''(x_*) - \mathbf{g}''(x_*) > 0$. Note that if $V(x)$ is convex in $x \geq \mathbf{e}$, then $V''(x) - \mathbf{g}''(x) > 0$ for all $x > \mathbf{e}$. For later reference we note that $V''(x_*) - \mathbf{g}''(x_*) = -\mathbf{G}''(x_*)$ in terms of the notation (32) that is defined below.

The following is the first main result of this paper. Let $V(x)$ be a potential that is convex in $x \in (\mathbf{e}, \infty)$. For $a > \frac{1}{2}V'(\mathbf{e})$, let $x_0(a)$ be the unique maximizer of the function $\mathbf{g}(x) - V(x) + ax$ in $x \in (\mathbf{e}, \infty)$. Such a maximizer exists since $\mathbf{g}(x) - V(x)$ in $x \in (\mathbf{e}, \infty)$ is strictly concave and $\mathbf{g}'(\mathbf{e}) - V'(\mathbf{e}) + a = -\frac{1}{2}V'(\mathbf{e}) + a > 0$ (see (30)) and $\mathbf{g}'(x) - V'(x) + a < 0$ for all large enough x . This $x_0(a)$ is same as in Lemma 1.3.

Theorem 1.1 (convex potential). *Let $V(x)$ be a potential that is convex in $x \in (\mathbf{e}, \infty)$. Set*

$$\mathbf{a}_c := \frac{1}{2}V'(\mathbf{e}). \quad (25)$$

The following holds for each $T \in \mathbb{R}$ as $n \rightarrow \infty$ and $j = O(1)$.

(a) For $a < \mathbf{a}_c$,

$$\lim_{n \rightarrow \infty} \mathbb{P}_{n-j+1, n}(a; I_n^T) = F_0(T). \quad (26)$$

(b) For

$$a = \mathbf{a}_c + \frac{\beta \alpha}{n^{1/3}}, \quad (27)$$

where α is in a compact subset of \mathbb{R} , we have

$$\lim_{n \rightarrow \infty} \mathbb{P}_{n-j+1, n}(a; I_n^T) = F_1(T; -\alpha). \quad (28)$$

(c) For $a > \mathbf{a}_c$,

$$\lim_{n \rightarrow \infty} \mathbb{P}_{n-j+1, n}(a; J_n^T(x_0(a))) = G(T). \quad (29)$$

Hence the transition phenomenon is universal for convex potentials. The next two subsections are about non-convex potentials

1.4 Critical value and secondary critical values

In this Subsection, we define critical values and the secondary critical values of a .

By definition (14), $\mathbf{g}(x)$ is real analytic in (\mathbf{e}, ∞) , is continuously differentiable in $[\mathbf{e}, \infty)$ and satisfies

$$\begin{aligned} \mathbf{g}'(x) &> 0, \quad \mathbf{g}''(x) < 0, \quad \text{for } x \in (\mathbf{e}, \infty), \\ \mathbf{g}(\mathbf{e}) &= \frac{V(\mathbf{e}) + \ell}{2}, \quad \mathbf{g}'(\mathbf{e}) = \frac{V'(\mathbf{e})}{2}, \quad \lim_{x \rightarrow \infty} \mathbf{g}'(x) = 0. \end{aligned} \quad (30)$$

Definition 1.1. For $a \in (0, \frac{1}{2}V'(\mathbf{e}))$, define $c = c(a)$ as the unique point in (\mathbf{e}, ∞) satisfying

$$\mathbf{g}'(c(a)) = a. \quad (31)$$

For $a \geq \frac{1}{2}V'(\mathbf{e})$, define $c(a) := \mathbf{e}$.

Note that $c(a)$ decreases strictly in $a \in (0, V'(\mathbf{e})/2)$ and continuous in $a \in (0, \infty)$.

Define two auxiliary functions

$$\begin{aligned} \mathbf{G}(z) &= \mathbf{G}(z; a) := \mathbf{g}(z) - V(z) + az, \\ \mathbf{H}(z) &= \mathbf{H}(z; a) := -\mathbf{g}(z) + az + \ell \end{aligned} \quad (32)$$

for $z \in \mathbb{C} \setminus (-\infty, \mathbf{e})$. Observe the following Lemma. The proof follows straightforwardly from the definition of \mathbf{g} , the variational condition (12), the assumption (7) on V and (15). We omit the details.

Lemma 1.1. *Let $a > 0$. We have the following properties.*

- (a) $\mathbf{H}(x)$ is a convex function in $x \in [\mathbf{e}, \infty)$ with the unique minimum attained at $x = c(a)$.
- (b) $\mathbf{H}(x) > \mathbf{G}(x)$ for all $x \in (\mathbf{e}, \infty)$.
- (c) $\mathbf{H}(\mathbf{e}) = \mathbf{G}(\mathbf{e}) = -\frac{1}{2}V(\mathbf{e}) + a\mathbf{e} + \frac{1}{2}\ell$.
- (d) $\lim_{x \downarrow \mathbf{e}} \mathbf{H}'(x) = \lim_{x \downarrow \mathbf{e}} \mathbf{G}'(x) = a - \frac{1}{2}V'(\mathbf{e})$.
- (e) As $x \rightarrow +\infty$, $\mathbf{H}(x) \rightarrow +\infty$, $\mathbf{H}(x)/x \rightarrow a$, $\mathbf{G}(x) \rightarrow -\infty$ and $\mathbf{G}(x)/x \rightarrow -\infty$.

See Figures 2, 3 and 4 for a few examples of the graphs of \mathbf{G} and \mathbf{H} .

Define the set

$$\mathcal{A}_V := \{a \in (0, \infty) \mid \text{there exists } \bar{x} \in (c(a), \infty) \text{ such that } \mathbf{G}(\bar{x}; a) > \mathbf{H}(c(a); a)\}. \quad (33)$$

Definition 1.2. The critical value for the spiked source model with potential V is defined as

$$\mathbf{a}_c := \inf \mathcal{A}_V. \quad (34)$$

Lemma 1.2. *We have the following properties.*

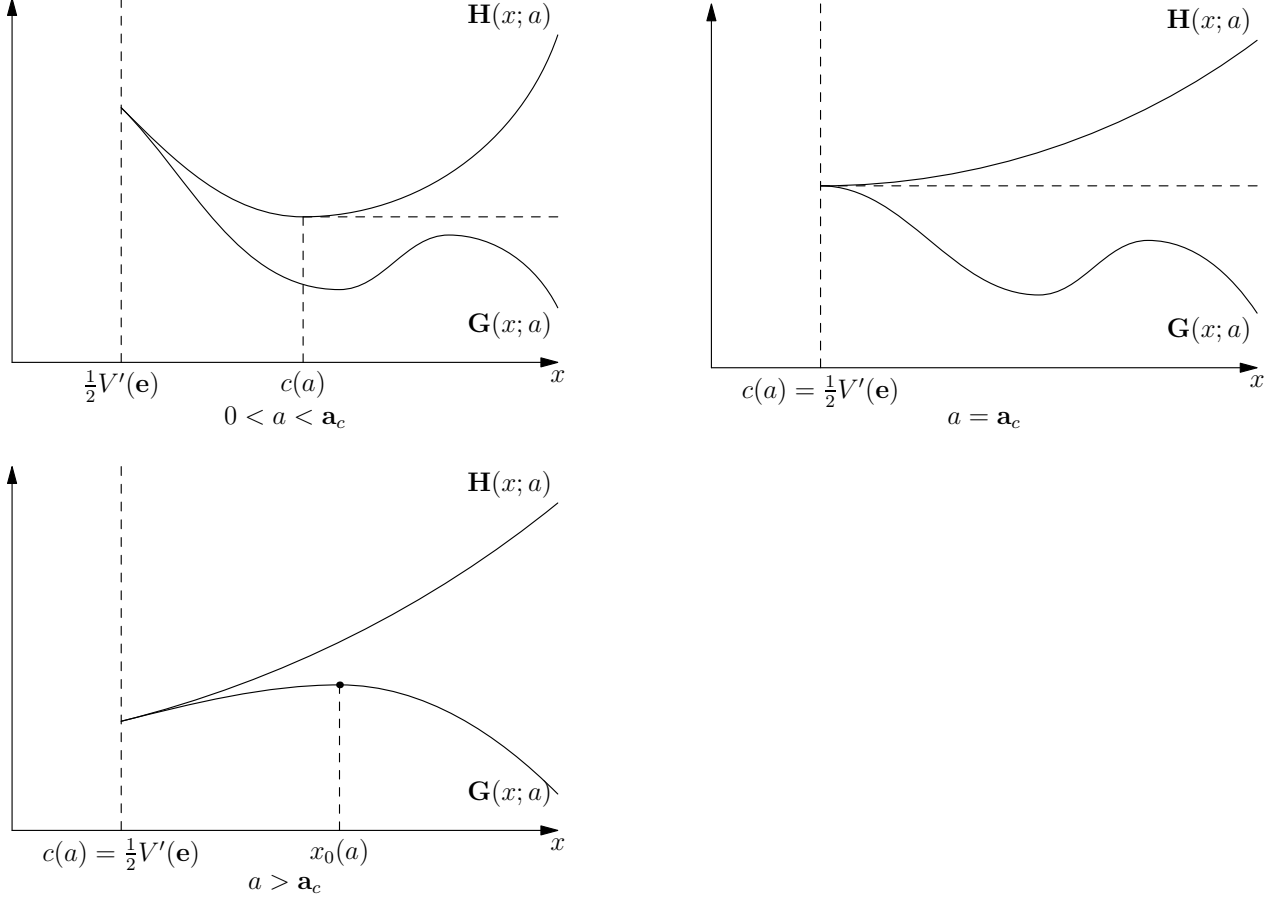


Figure 2: Schematic graphs of the functions $\mathbf{H}(x; a)$ and $\mathbf{G}(x; a)$ for a potential V such that $\mathbf{a}_c = \frac{1}{2}V'(\mathbf{e})$, assuming that $a \notin \mathcal{J}_V$.

- (a) $(\frac{1}{2}V'(\mathbf{e}), \infty) \subset \mathcal{A}_V$. Hence $\mathbf{a}_c \leq \frac{1}{2}V'(\mathbf{e})$.
- (b) The set \mathcal{A}_V is an open, semi-infinite interval. Hence $\mathcal{A}_V = (\mathbf{a}_c, \infty)$.
- (c) $\mathbf{a}_c > 0$.
- (d) For $0 < a < \mathbf{a}_c$, we have $\mathbf{G}(x; a) < \mathbf{H}(c(a); a)$ for all $x \in (c(a), \infty)$.
- (e) If the potential $V(x)$ is convex for $x \geq \mathbf{e}$, then $\mathbf{a}_c = \frac{1}{2}V'(\mathbf{e})$, and $\mathbf{G}(x; \mathbf{a}_c) < \mathbf{H}(\mathbf{e}; \mathbf{a}_c)$ for all $x > \mathbf{e}$. (Note that $c(\mathbf{a}_c) = \mathbf{e}$ and $\mathbf{G}(\mathbf{e}; \mathbf{a}_c) = \mathbf{H}(\mathbf{e}; \mathbf{a}_c)$.)
- (f) If the potential V is such that $\mathbf{a}_c < \frac{1}{2}V'(\mathbf{e})$, then $\mathbf{G}(x; \mathbf{a}_c) \leq \mathbf{H}(c(\mathbf{a}_c); \mathbf{a}_c)$ for all $x \in (c(a), \infty)$, and the equality is attained at least at one point.

Proof. (a) Let $a \in (\frac{1}{2}V'(\mathbf{e}), \infty)$. Since $\lim_{x \downarrow \mathbf{e}} \mathbf{G}'(x) = a - \frac{1}{2}V'(\mathbf{e}) > 0$, there is $\bar{x} > \mathbf{e}$ such that $\mathbf{G}(\bar{x}) > \mathbf{G}(\mathbf{e}) = \mathbf{H}(\mathbf{e})$. Thus $a \in \mathcal{A}_V$.

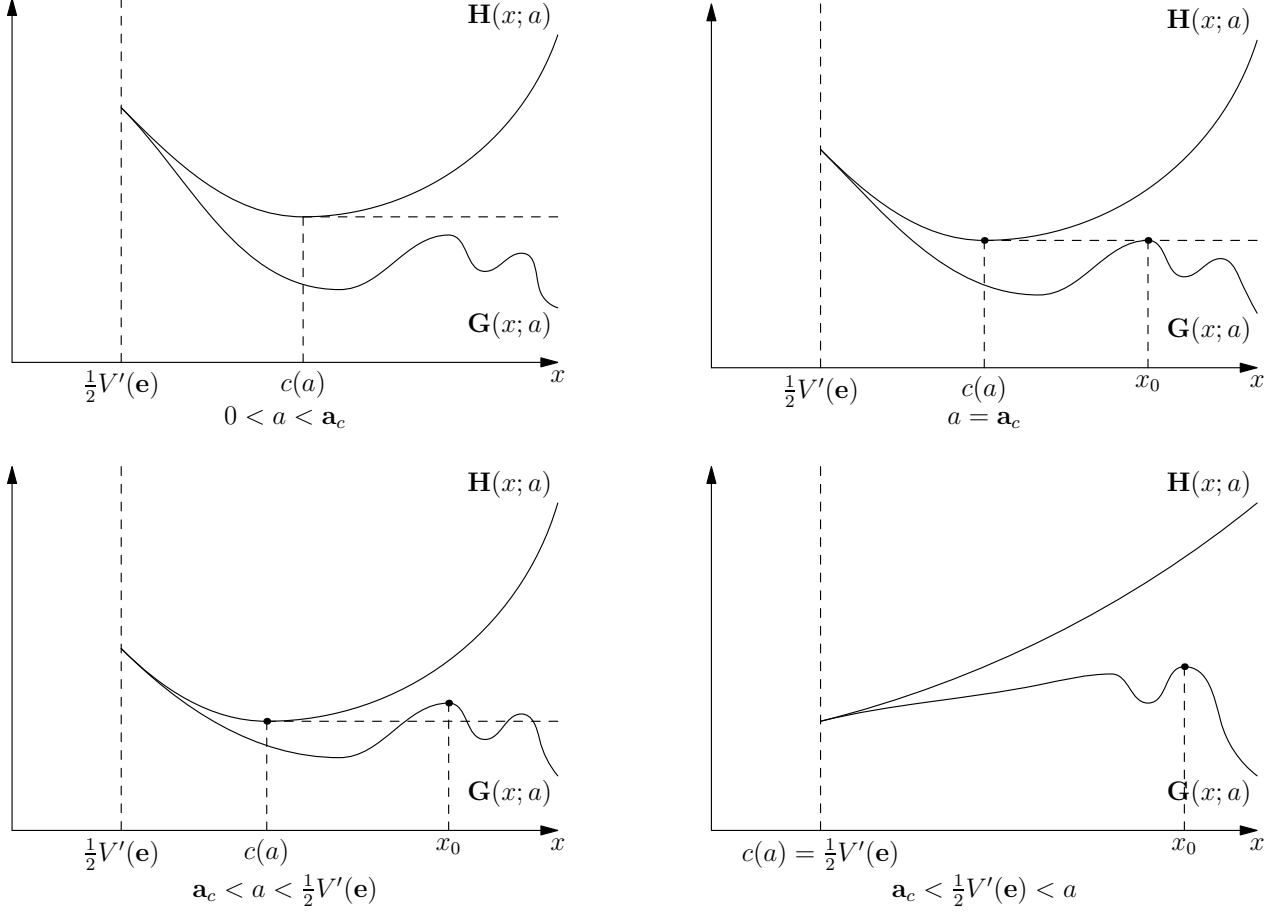


Figure 3: Schematic graphs of the functions $\mathbf{H}(x; a)$ and $\mathbf{G}(x; a)$ for a potential V such that $\mathbf{a}_c < \frac{1}{2}V'(\mathbf{e})$, assuming that $a \notin \mathcal{I}_V$.

- (b) The continuity of \mathbf{G} and \mathbf{H} in a implies that \mathcal{A}_V is an open set. Now we show that \mathcal{A} is a semi-infinite interval. Suppose that $a \in \mathcal{A}_V$ and $a < \frac{1}{2}V'(\mathbf{e})$. Let $\bar{x} \in (c(a), \infty)$ be the point such that $\mathbf{G}(\bar{x}; a) > \mathbf{H}(c(a); a)$. Let $a' \in (a, \frac{1}{2}V'(\mathbf{e})]$. From Definition 1.1 of $c(a)$, we see that $c(a') < c(a)$, and hence $\bar{x} \in (c(a'), \infty)$. Moreover,

$$\mathbf{G}(\bar{x}; a') - \mathbf{H}(c(a'); a') = [\mathbf{G}(\bar{x}; a) - \mathbf{H}(c(a); a)] + [\mathbf{H}(c(a); a') - \mathbf{H}(c(a'); a')] + [(a' - a)(\bar{x} - c(a))] \quad (35)$$

is strictly positive since each term in bracket is strictly positive. Thus $a' \in \mathcal{A}_V$, and this, together with (a), implies that \mathcal{A}_V is a semi-infinite interval.

- (c) We have $\mathbf{G}(x) - \mathbf{H}(c(a)) = -V(x) + a(x - c(a)) + \mathbf{g}(x) + \mathbf{g}(c(a)) - \ell \leq -V(x) + ax + 2\mathbf{g}(x) - \ell$. This tends to $-\infty$ as $x \rightarrow +\infty$ due to the growth condition (7) on V . Also $c(a) \rightarrow +\infty$ as $a \downarrow 0$. Therefore when a is close to 0, $\mathbf{G}(x) - \mathbf{H}(c(a)) < 0$ for $x > c(a)$. Hence, $a \notin \mathcal{A}_V$ if a is small enough.

- (d) Let $0 < a < \mathbf{a}_c$. Suppose that there is $\bar{x} \in (c(a), \infty)$ such that $\mathbf{G}(\bar{x}; a) = \mathbf{H}(c(a); a)$. For

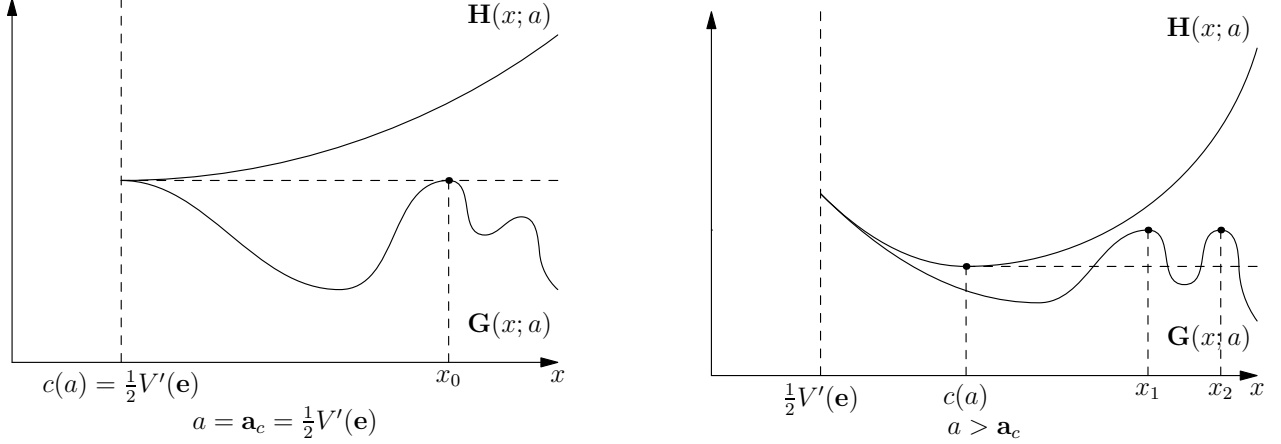


Figure 4: Schematic graphs of functions $\mathbf{H}(x; a)$ and $\mathbf{G}(x; a)$ when $a = \mathbf{a}_c \in \mathcal{J}_V$ and $\mathbf{a}_c > a \in \mathcal{J}_V$.

any $a' \in (a, \mathbf{a}_c)$, we have $c(a) > c(a')$ since $a < a' < \frac{1}{2}V'(\mathbf{e})$. Thus we find from (35) that $\mathbf{G}(\bar{x}; a') - \mathbf{H}(c(a'); a') > 0$. This implies that $a' \in \mathcal{A}_V$ which is a contradiction.

(e) Let $0 < a < \frac{1}{2}V'(\mathbf{e})$. We will show that $a \notin \mathcal{A}$. Since V is convex, $\mathbf{G}(x)$ is concave in $x \in (\mathbf{e}, \infty)$. As $\mathbf{G}'(\mathbf{e}; a) < 0$, this implies that $\mathbf{G}(x)$ is decreasing in $x \in (\mathbf{e}, \infty)$. Thus for $x \in (c(a), \infty)$, $\mathbf{G}(x) < \mathbf{G}(c(a)) < \mathbf{H}(c(a))$. Hence $a \notin \mathcal{A}_V$. When $a = \frac{1}{2}V'(\mathbf{e})$, a similar argument implies that $\mathbf{G}(x) < \mathbf{H}(x)$ for all $x > \mathbf{e}$.

(f) This follows from the continuity of \mathbf{G} and \mathbf{H} in a and the fact that $\mathbf{a}_c = \inf \mathcal{A}_V$. □

See typical graphs of \mathbf{G} and \mathbf{H} for $\mathbf{a}_c = \frac{1}{2}V'(\mathbf{e})$ in Figure 2, and typical graphs of \mathbf{G} and \mathbf{H} for $0 < \mathbf{a}_c < \frac{1}{2}V'(\mathbf{e})$ in Figure 3.

Remark 1.1. When V is non-convex, there may exist $\bar{x} > \mathbf{e}$ such that $\mathbf{G}(\bar{x}; \mathbf{a}_c) = \mathbf{H}(\mathbf{e}, \mathbf{a}_c)$ even if $\mathbf{a}_c = \frac{1}{2}V'(\mathbf{e})$.

By Definition 1.2 of \mathbf{a}_c , when $a > \mathbf{a}_c$, $\mathbf{G}(x; a) > \mathbf{H}(c(a); a)$ for some $x > c(a)$. The point x at which $\mathbf{G}(x; a)$ attains its maximum plays an important role. Indeed, we will show in the below that if the maximum is attained at a unique point, then $\xi_{\max}(n)$ converges to this point (see Theorem 1.2). However, it may happen that for some a 's, the function $\mathbf{G}(x; a)$ attain its maximum at more than one point. Let

$$\mathbf{G}_{\max}(a) := \max_{x \in [c(a), \infty)} \mathbf{G}(x; a), \quad (36)$$

and define

$$\mathcal{J}_V := \{a \in [\mathbf{a}_c, \infty) \mid \mathbf{G}_{\max}(a) \text{ is attained at more than one point}\}. \quad (37)$$

This set is discrete since $\mathbf{G}(x; a)$ is analytic in both x and a . Note that when V is convex, $\mathbf{a}_c \notin \mathcal{J}_V$ from Lemma 1.2(e). For a non-convex V , as indicated in Remark 1.1, \mathbf{a}_c may or may not be in \mathcal{J}_V no matter if $\mathbf{a}_c = \frac{1}{2}V'(\mathbf{e})$. See typical graphs of \mathbf{G} and \mathbf{H} for $a \in \mathcal{J}_V$ in Figure 4.

We have the following Lemma.

Lemma 1.3. (a) For $a \in [\mathbf{a}_c, \infty)$ such that $a \notin \mathcal{J}_V$, let $x_0(a)$ be the unique point in $[c(a), \infty)$ at which $\mathbf{G}(x; a)$ attains its maximum. Then $x_0(a)$ is a continuous, strictly increasing function in $a \in [\mathbf{a}_c, \infty) \setminus \mathcal{J}_V$.

(b) If $a_0 \in \mathcal{J}_V$ and $a_0 > \mathbf{a}_c$, then

$$\lim_{a \uparrow a_0} x_0(a) < \lim_{a \downarrow a_0} x_0(a). \quad (38)$$

Note that if $a \in \mathcal{J}_V$ satisfies $a > \mathbf{a}_c$ or $a = \mathbf{a}_c < \frac{1}{2}V'(\mathbf{e})$, then there exist points $x_1(a) < x_2(a) < \dots < x_r(a)$ in $(c(a), \infty)$, for some $r \geq 2$, such that

$$\mathbf{G}_{\max}(a) = \mathbf{G}(x_1(a); a) = \dots = \mathbf{G}(x_r(a); a). \quad (39)$$

On the other hand, if V is a potential such that $\mathbf{a}_c = \frac{1}{2}V'(\mathbf{e})$ and $\mathbf{a}_c \in \mathcal{J}_V$, then there exist, for some $r \geq 1$, $x_1(\mathbf{a}_c) < x_2(\mathbf{a}_c) < \dots < x_r(\mathbf{a}_c)$ in (\mathbf{e}, ∞) such that

$$\mathbf{G}(\mathbf{e}; a) = \mathbf{G}_{\max}(a) = \mathbf{G}(x_1(\mathbf{a}_c); \mathbf{a}_c) = \dots = \mathbf{G}(x_r(\mathbf{a}_c); \mathbf{a}_c). \quad (40)$$

Proof of Lemma 1.3. The continuity of $x_0(a)$ for $a \notin \mathcal{J}_V$ is a direct consequence of the continuity of $\mathbf{G}(x; a)$ in both x and a . Let $\mathbf{a}_c \leq a_1 < a_2$ and $a_1, a_2 \notin \mathcal{J}_V$. If we assume $x_0(a_1) \geq x_0(a_2)$, then since $\mathbf{G}(x_0(a_1); a_1) > \mathbf{G}(x_0(a_2); a_1)$, we have

$$\mathbf{G}(x_0(a_1); a_2) = \mathbf{G}(x_0(a_1); a_1) + (a_2 - a_1)x_0(a_1) > \mathbf{G}(x_0(a_2); a_1) + (a_2 - a_1)x_0(a_2) = \mathbf{G}(x_0(a_2); a_2). \quad (41)$$

This is contradictory to the assumption that $x_0(a_2)$ is the maximizer of $\mathbf{G}(x; a_2)$. Thus $x_0(a_1) < x_0(a_2)$.

If $a_0 \in \mathcal{J}_V$ and $a_0 > \mathbf{a}_c$, then $\mathbf{G}(x; a_0)$ attains its maximum in $[c(a_0), \infty)$ at $x_1(a_0), \dots, x_r(a_0)$ for some $r \geq 2$ as in (39). It is easy to check from the continuity of $\mathbf{G}(x; a)$ in a that $\lim_{a \uparrow a_0} x_0(a) = x_1(a_0)$ and $\lim_{a \downarrow a_0} x_0(a) = x_r(a_0)$. \square

Definition 1.3. The secondary critical values for the spiked model are defined as the points $a \in \mathcal{J}_V$ such that $a > \mathbf{a}_c$.

Remark 1.2. For a potential V such that $V(x)$ is convex for $x \geq \mathbf{e}$, $\mathcal{J}_V = \emptyset$ since $\mathbf{G}'(x; a)$ is a decreasing function in $x \geq \mathbf{e}$. Hence there is no secondary critical value.

1.5 Statement of results: non-convex potentials

Let $V(x)$ be a potential satisfying the conditions (6)–(8). Let I_n^T and $J_n^T(x_*)$ be the intervals defined in (23) and (24), respectively.

Theorem 1.2 (away from critical values). *The following holds for each $T \in \mathbb{R}$ as $n \rightarrow \infty$ and $j = O(1)$.*

(a) For $a < \mathbf{a}_c$,

$$\lim_{n \rightarrow \infty} \mathbb{P}_{n-j+1, n}(a; I_n^T) = F_0(T). \quad (42)$$

(b) For $a > \mathbf{a}_c$ such that $a \notin \mathcal{J}_V$, if $\mathbf{G}''(x_0(a)) \neq 0$, then

$$\lim_{n \rightarrow \infty} \mathbb{P}_{n-j+1,n}(a; J_n^T(x_0(a))) = G(T). \quad (43)$$

When a is at or near the critical value \mathbf{a}_c , we have the following result. The case when $a = \mathbf{a}_c$ is attained by setting $\alpha = 0$.

Theorem 1.3 (at or near the critical value). *We have the following for each $T \in \mathbb{R}$.*

(a) Suppose that V is a potential such that $\mathbf{a}_c = \frac{1}{2}V'(\mathbf{e})$ and $\mathbf{a}_c \notin \mathcal{J}_V$. Then for

$$a = \mathbf{a}_c + \frac{\beta\alpha}{n^{1/3}}, \quad (44)$$

where α is in a compact subset of \mathbb{R} , we have

$$\lim_{n \rightarrow \infty} \mathbb{P}_{n-j+1,n}(a; I_n^T) = F_1(T; -\alpha). \quad (45)$$

(b) Let V be a potential such that $\mathbf{a}_c < \frac{1}{2}V'(\mathbf{e})$. If $\mathbf{a}_c \notin \mathcal{J}_V$ and $\mathbf{G}''(x_0(\mathbf{a}_c); \mathbf{a}_c) \neq 0$, then for

$$a = \mathbf{a}_c + \frac{\alpha}{n}, \quad (46)$$

where α is in a compact subset of \mathbb{R} , we have, as $n \rightarrow \infty$,

$$\mathbb{P}_{n-j+1,n}(a; I_n^T) = p_{j,n}(\alpha)F_0(T) + o(1) \quad (47)$$

and

$$\mathbb{P}_{n-j+1,n}(a; J_n^T(x_0(\mathbf{a}_c))) = p_{j,n}(\alpha) + (1 - p_{j,n}(\alpha))G(T) + o(1), \quad (48)$$

where the constant $p_{j,n}(\alpha) \in (0, 1)$ is defined by (145). As a function of α , $p_{j,n}(\alpha)$ is decreasing and satisfies $p_{j,n}(\alpha) \rightarrow 0$ as $\alpha \rightarrow \infty$ and $p_{j,n}(\alpha) \rightarrow 1$ as $\alpha \rightarrow -\infty$ for each fixed n . Also for each fixed α , $p_{j,n}(\alpha)$ lies in a compact subset of $(0, 1)$ for all large n .

Remark 1.3. When the potential $V(x)$ is convex for $x \geq \mathbf{e}$, then $\mathbf{a}_c = \frac{1}{2}V'(\mathbf{e})$, $\mathcal{J}_V = \emptyset$ (see Remark 1.2) and $\mathbf{G}''(x) < 0$ for all $x > \mathbf{e}$. Hence Theorem 1.2 and Theorem 1.3(a) imply Theorem 1.1.

For a at or near the secondary critical values $\mathcal{J}_V \setminus \{\mathbf{a}_c\}$, we have the following result.

Theorem 1.4 (at or near the secondary critical values). *Let V be a potential such that $\mathcal{J}_V \neq \emptyset$. Let $a_0 \in \mathcal{J}_V \setminus \{\mathbf{a}_c\}$ be a secondary critical point. If $\mathbf{G}(x; a_0)$ attains its maximum $\mathbf{G}_{\max}(a_0)$ at two points $x_1(a_0) < x_2(a_0)$ in $(c(a), \infty)$ and if $\mathbf{G}''(x_1(a_0); a_0) \neq 0$ and $\mathbf{G}''(x_2(a_0); a_0) \neq 0$, then for*

$$a = a_0 + \frac{\alpha}{n}, \quad (49)$$

where α is in a compact subset of \mathbb{R} , we have

$$\mathbb{P}_{n-j+1,n}(a; J_n^T(x_1(a_0))) = p_{j,n}^{(1)}(\alpha)G(T) + o(1) \quad (50)$$

and

$$\mathbb{P}_{n-j+1,n}(a; J_n^T(x_2(a_0))) = p_{j,n}^{(1)}(\alpha) + p_{j,n}^{(2)}(\alpha)G(T) + o(1), \quad (51)$$

where $p_{j,n}^{(1)}(\alpha)$ and $p_{j,n}^{(2)}(\alpha) \in (0, 1)$ are defined in (164) and (158), and $p_{j,n}^{(1)}(\alpha) + p_{j,n}^{(2)}(\alpha) = 1$. As a function of α , $p_{j,n}^{(1)}(\alpha)$ is decreasing and satisfies $p_{j,n}^{(1)}(\alpha) \rightarrow 0$ as $\alpha \rightarrow \infty$ and $p_{j,n}^{(1)}(\alpha) \rightarrow 1$ as $\alpha \rightarrow -\infty$ for each fixed n . Also for each fixed α , $p_{j,n}^{(1)}(\alpha)$ is in a compact subset of $(0, 1)$ independent of n .

The above three theorems describe the ‘generic’ cases. The next part describes the three ‘exceptional cases’.

As the first exceptional case, suppose that in Theorem 1.4, the maximum of \mathbf{G} is attained at more than two points. Let $x_1(a) < x_2(a) < \dots < x_r(a)$ be these maximizers. If $\mathbf{G}''(x_k(a_0); a_0) \neq 0$ for all $k = 1, \dots, r$, then we have, for each $k = 1, \dots, r$,

$$\mathbb{P}_{n-j+1,n}(a; J_n^T(x_k(a_0))) = \sum_{i=1}^{k-1} p_{j,n}^{(i)}(\alpha) + p_{j,n}^{(k)}(\alpha)G(T) + o(1), \quad (52)$$

for some $p_{j,n}^{(i)}(\alpha) \in (0, 1)$ such that $p_{j,n}^{(1)}(\alpha) + \dots + p_{j,n}^{(r)}(\alpha) = 1$. Explicitly, $p_{j,n}^{(i)}(\alpha) := \frac{A_i(\alpha)}{A_1(\alpha) + \dots + A_r(\alpha)}$ where $A_i(\alpha)$ is defined in (158). The situation when $\mathbf{a}_c \in \mathcal{J}_V$ in Theorem 1.3(b) is similar. In this case, the maximum of $\mathbf{G}(x; \mathbf{a}_c)$ in $(c(\mathbf{a}_c), \infty)$ is attained at $x_1(\mathbf{a}_c) < x_2(\mathbf{a}_c) < \dots < x_r(\mathbf{a}_c)$ for some $r \geq 2$ (see (39)). Assume that $\mathbf{G}''(x_i(\mathbf{a}_c); \mathbf{a}_c) \neq 0$ for all $i = 1, \dots, r$. Then with $C_i(\alpha)$, $i = 1, \dots, r$, defined by (144) with $x_0(\mathbf{a}_c)$ replaced by $x_i(\mathbf{a}_c)$, set $p_{j,n}^{(i)}(\alpha) := \frac{C_i(\alpha)}{C_0 + C_1(\alpha) + \dots + C_r(\alpha)}$, $i = 1, \dots, r$, where C_0 is defined by (143). Then (47) holds with $p_{j,n}(\alpha)$ replaced by $p_{j,n}^{(0)}(\alpha) := 1 - p_{j,n}^{(1)}(\alpha) - \dots - p_{j,n}^{(r)}(\alpha)$. The limit in (48) is replaced by

$$\mathbb{P}_{n-j+1,n}(a; J_n^T(x_k(a_0))) = \sum_{i=0}^{k-1} p_{j,n}^{(i)}(\alpha) + p_{j,n}^{(k)}(\alpha)G(T) + o(1) \quad (53)$$

for $k = 1, \dots, r$.

The second exceptional case is when $\mathbf{a}_c \in \mathcal{J}_V$ in Theorem 1.3, case 1. This case is given in the following Theorem. In this case, there are two natural scalings in a .

Theorem 1.5. *Let V be a potential such that $\mathbf{a}_c = \frac{1}{2}V'(\mathbf{e})$. Suppose that $\mathbf{a}_c \in \mathcal{J}_V$. Assume that $\mathbf{G}(x; \mathbf{a}_c)$ attains its maximum at the unique point $x_0(\mathbf{a}_c) \in (c(\mathbf{a}_c), \infty)$ and $\mathbf{G}''(x_0(\mathbf{a}_c); \mathbf{a}_c) \neq 0$. Then the following holds.*

(a) For

$$a = \mathbf{a}_c + \frac{\beta\alpha}{n^{1/3}}, \quad (54)$$

where α is in a compact subset of $(-\infty, 0)$, we have

$$\lim_{n \rightarrow \infty} \mathbb{P}_n(a; I_n^T) = F_1(T; -\alpha). \quad (55)$$

(b) For

$$a = \mathbf{a}_c + \frac{\alpha'}{n}, \quad (56)$$

where α' is in a compact subset of \mathbb{R} , we have

$$\mathbb{P}_{n-j+1,n}(a; I_n^T) = p_{j,n}(\alpha')F_1(T) + o(1) \quad (57)$$

where $F_1(T) = F_1(T; 0)$, and

$$\mathbb{P}_{n-j+1,n}(a; J_n^T(x_0(\mathbf{a}_c))) = p_{j,n}(\alpha') + (1 - p_{j,n}(\alpha'))G(T) \quad (58)$$

where $p_{j,n}(\alpha')$ is defined in (280). As a function of α , $p_{j,n}(\alpha')$ is decreasing and satisfies $p_{j,n}(\alpha') \rightarrow 0$ as $\alpha' \rightarrow \infty$ and $p_{j,n}(\alpha') \rightarrow 1$ as $\alpha' \rightarrow -\infty$ for each fixed n . Also for each fixed α , $p_{j,n}(\alpha)$ is in a compact subset of $(0, 1)$ independent of n .

If the maximum of $\mathbf{G}(x; \mathbf{a}_c)$ is attained at more than one point, then (58) should be changed in a natural way as in (53).

The third exceptional case is when the double derivative of $\mathbf{G}(x; a)$ vanishes at its maximizers. Then the function $G(x)$ is replaced by its higher analogue, and the scalings in the interval and a are also changed accordingly. Concretely, in Theorem 1.2(b), if $\mathbf{G}''(x_0(a)) = 0$, then since $x_0(a)$ is the maximum point, there exists $k > 1$ such that $\mathbf{G}^{(2k)}(x_0(a)) < 0$ and $\mathbf{G}^{(j)}(x_0(a)) = 0$ for all $j = 1, 2, \dots, 2k - 1$. Then (43) is changed to

$$\lim_{n \rightarrow \infty} \mathbb{P}_{n-j+1,n}(a; \hat{J}_n^T(x_0(a); k)) = \frac{\int_{-\infty}^T e^{-x^{2k}} dx}{\int_{-\infty}^{\infty} e^{-x^{2k}} dx}, \quad (59)$$

where the interval $\hat{J}_n^T(x_*; k)$ is defined by

$$\hat{J}_n^T(x_*; k) := \left[x_* + \left(\frac{n(V^{(2k)}(x_*) - \mathbf{g}^{(2k)}(x_*))}{(2k)!} \right)^{-1/(2k)} T, \infty \right). \quad (60)$$

In Theorem 1.3(b), Theorem 1.4 and Theorem 1.5, a is scaled as $a = a_0 + \frac{\alpha}{n}$. When $\mathbf{G}''(x_i(a_0); a_0) = 0$, ($i = 0$ in Theorems 1.3(b) and 1.5, and $i = 1, 2$ in Theorem 1.4,) then this scaling also needs to be changed. For example, Theorem 1.4 is changed to the following theorem.

Theorem 1.6. *Let $a_0 \in \mathcal{J}_V \setminus \{\mathbf{a}_c\}$ be a secondary critical value. Assume that $\mathbf{G}(x; a_0)$ attains its maximum $\mathbf{G}_{\max}(a_0)$ at two points $x_1(a_0) < x_2(a_0)$ in $(c(a), \infty)$. Suppose that $\mathbf{G}''(x_1(a_0); a_0) \neq 0$, and for some $k > 1$, and suppose that $\mathbf{G}^{(2k)}(x_2(a_0); a_0) \neq 0$ and $\mathbf{G}^{(i)}(x_2(a_0); a_0) = 0$ for all $i = 1, \dots, 2k - 1$. Then for*

$$a = a_0 - q \frac{\log n}{n} + \frac{\alpha}{n}, \quad \text{where } q := \frac{\frac{1}{2} - \frac{1}{2k}}{x_2(a_0) - x_1(a_0)} \quad (61)$$

where α is in a compact subset of \mathbb{R} , we have

$$\mathbb{P}_{n-j+1,n}(a; J_n^T(x_1(a_0))) = p_{j,n}^{(1)}(\alpha)G(T) + o(1), \quad (62)$$

and

$$\mathbb{P}_{n-j+1,n} \left(a; \hat{J}_n^T(x_2(a_0); k) \right) = p_{j,n}^{(1)}(\alpha) + p_{j,n}^{(2)}(\alpha) \frac{\int_{-\infty}^T e^{-x^{2k}} dx}{\int_{-\infty}^{\infty} e^{-x^{2k}} dx} + o(1), \quad (63)$$

where the interval $\hat{J}_n^T(x; k)$ is defined by (60), $p_{j,n}^{(1)}(\alpha)$ and $p_{j,n}^{(2)}(\alpha)$ are defined in (182), and $p_{j,n}^{(1)}(\alpha) + p_{j,n}^{(2)}(\alpha) = 1$. As a function of α , $p_{j,n}^{(1)}(\alpha)$ is decreasing and satisfies $p_{j,n}^{(1)}(\alpha) \rightarrow 0$ as $\alpha \rightarrow \infty$ and $p_{j,n}^{(1)}(\alpha) \rightarrow 1$ as $\alpha \rightarrow -\infty$ for each fixed n . Also for each fixed α , $p_{j,n}^{(1)}(\alpha)$ is in a compact subset of $(0, 1)$ independent of n .

The changes needed for Theorem 1.3(b), and Theorem 1.5 are analogous. Also it may happen that the two or more of the exceptional cases occur simultaneously. Then one needs simply combine the results together in a straightforward way, and we skip the details.

We remark that if the support J of the equilibrium measure is of one interval i.e. $N = 0$ (see (9)), then the probabilities $p_{j,n}$ and $p_{j,n}^{(j)}$ in Theorem 1.3(b), 1.4, 1.5(b) and 1.6 do not depend on n . This follows from Remark 6.1 on $\mathcal{M}_{j,n}$, $\tilde{\mathcal{M}}_{j,n}$ and $\mathcal{B}_{j,n}(\mathbf{e})$ and the definition of these probabilities. When $N > 0$, the dependence of these probabilities on n is from the theta function formula of $\mathcal{M}_{j,n}$, $\tilde{\mathcal{M}}_{j,n}$ and $\mathcal{B}_{j,n}(\mathbf{e})$ in Section 6, and is in a quasi-periodic way.

We also remark that one can obtain the convergence in probability 1 as in (3) from the above theorems together with the fact that all of the limiting distributions decay rapidly at the tails.

An explicit example of a potential such that $\mathbf{a}_c < \frac{1}{2}V'(\mathbf{e})$ can be constructed as follows. We use the potential defined in [18, Formula (4.14)] (we change the original notation e into \bar{e} here):

$$V_{\bar{e},\epsilon}(x) = \frac{1-\epsilon}{1+\bar{e}\tilde{e}} \left(\frac{1}{4}x^4 - \frac{\bar{e}+\tilde{e}}{3}x^3 + \frac{\bar{e}\tilde{e}-2}{2}x^2 + 2(\bar{e}+\tilde{e})x \right), \quad (64)$$

where ϵ is a very small positive number, \bar{e} is any number > 2 and \tilde{e} is determined by e from the condition that

$$\int_2^{\bar{e}} (x-\bar{e})(x-\tilde{e})\sqrt{x^2-4}dx = 0. \quad (65)$$

From results in [18, Section 4], it is known that $V_{\bar{e},\epsilon}$ is a regular potential with the support of the equilibrium measure given by $[-2+O(\epsilon), 2+O(\epsilon)]$. For all $x > \mathbf{e} = 2+O(\epsilon)$, (15) holds. However, at $x = \bar{e}$,

$$2\mathbf{g}(\bar{e}) - V(\bar{e}) - \ell = -E(\epsilon) \quad (66)$$

for some $E(\epsilon)$ satisfying $E(\epsilon) = O(\epsilon)$ and $E(\epsilon) > 0$. Hence for any a , $\mathbf{G}(\bar{e}; a) - \mathbf{H}(\bar{e}; a) = O(\epsilon)$. Now there exists $a \in (0, \frac{1}{2}V'(\mathbf{e}))$ such that $c(a) \in (\mathbf{e}, \bar{e})$ since the minimizer $c(a)$ of $\mathbf{H}(x; a)$ is continuous in $a \in (0, \infty)$, decreases strictly in $a \in (0, \frac{1}{2}V'(\mathbf{e}))$, $\lim_{a \downarrow 0} c(a) = +\infty$ and $c(\frac{1}{2}V'(\mathbf{e})) = \mathbf{e}$ (see the sentence after Definition 1.1). Since $\mathbf{H}(c(a); a) < \mathbf{H}(\bar{e}; a)$, we have $\mathbf{G}(\bar{e}; a) > \mathbf{H}(c(a); a)$ if ϵ is small enough. Then $a \in \mathcal{A}_V$ and $\mathbf{a}_c < \frac{1}{2}V'_{\bar{e},\epsilon}(\mathbf{e})$ for each fixed $\bar{e} > 2$ if $\epsilon > 0$ is small enough.

The paper is organized as follows. The outline of the proof of theorems is given in Section 2. The results on the orthonormal polynomials and the kernel $K_{n-j,n}$ are summarized in Section 6.

The proofs of the theorems are given in Sections 3, 4 and 5. We consider three cases, $\mathbf{a}_c < \frac{1}{2}V'(\mathbf{e})$, $\mathbf{a}_c > \frac{1}{2}V'(\mathbf{e})$ and $\mathbf{a}_c = \frac{1}{2}V'(\mathbf{e})$ separately. Throughout this paper we only consider $a > 0$. The $a < 0$ case is discussed briefly in the end of Section 2.

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2 Outline of the proof

Let

$$p_j(x; n) = \gamma_j(n)x^j + \dots \quad (67)$$

be the orthonormal polynomial of degree j with respect to the weight $e^{-nV(x)}$. Here take $\gamma_j(n) > 0$ to make p_j unique. Set

$$\psi_j(x; n) := p_j(x; n)e^{-\frac{n}{2}V(x)}, \quad \varphi_j(x; n) := p_j(x; n)e^{-nV(x)}. \quad (68)$$

Let

$$K_{j,n}(x, y) := \sum_{i=0}^{j-1} \psi_i(x; n)\psi_i(y; n) = \frac{\gamma_{j-1}(n)}{\gamma_j(n)} \frac{\psi_j(x; n)\psi_{j-1}(y; n) - \psi_{j-1}(x; n)\psi_j(y; n)}{x - y} \quad (69)$$

be the Christoffel-Darboux kernel. Define the constant

$$\Gamma_j(a; n) := \langle e^{n(ax-V(x)/2)}, \psi_j(x; n) \rangle_{\mathbb{R}}, \quad (70)$$

and the function

$$\begin{aligned} \tilde{\psi}_j(x; a; n) &:= \frac{1}{\Gamma_j(a; n)} \left(e^{n(ax-V(x)/2)} - \sum_{i=0}^{j-1} \Gamma_i(a; n)\psi_i(x; n) \right) \\ &= \frac{1}{\Gamma_j(a; n)} \left(e^{n(ax-V(x)/2)} - \int_{\mathbb{R}} K_{j,n}(x, y)e^{n(ay-V(y))} dy \right). \end{aligned} \quad (71)$$

It is easy to check that $\tilde{\psi}_j(x; a; n)$ is characterized by the orthonormality conditions

$$\langle \tilde{\psi}_j(x; a; n), \psi_k(x; n) \rangle_{\mathbb{R}} = \delta_{jk}, \quad \text{for } k = 0, 1, \dots, j, \quad (72)$$

in the vector space spanned by $\{e^{-\frac{n}{2}V(x)}, xe^{-\frac{n}{2}V(x)}, \dots, x^{j-1}e^{-\frac{n}{2}V(x)}, e^{n(ax-V(x)/2)}\}$. In the multiple orthogonal polynomial theory, $\tilde{\psi}_j(x; a; n)e^{nV(x)/2}$ is the multiple orthogonal polynomial of type I with potentials $e^{-nV(x)}$ and e^{ax} , see e.g. [10]. From this follows the well-definedness of $\tilde{\psi}_j(x; a; n)$, i.e., $\Gamma_j(a; n) \neq 0$.

We sometimes drop the dependence on n or a in $\psi_j(x; n)$, $\Gamma_j(a; n)$ and $\tilde{\psi}_j(x; a; n)$ and write $\psi_j(x)$, $\Gamma_j(a)$ and $\tilde{\psi}_j(x)$ for convenience.

The starting point of our analysis is the following basic result in the theory of Hermitian matrix model with external source, specialized to the spiked source model of rank 1 (see e.g. [26], [10]): for any $E \subset \mathbb{R}$,

$$\mathbb{P}_{n-j+1,n}(a; E) = \det \left(1 - \chi_E \tilde{K}_{n-j+1,n} \chi_E \right), \quad (73)$$

where

$$\tilde{K}_{n-j+1,n} := K_{n-j,n} + \tilde{\psi}_{n-j} \otimes \psi_{n-j}. \quad (74)$$

Here χ_E denotes the projection operator on E and $K_{j,n}$ is the operator on $L^2(\mathbb{R})$, defined by the kernel (69). Note that the only term in (74) that depends on a is $\tilde{\psi}_{n-j}(x; a; n)$. The kernel $K_{n-j,n}(x, y)$ is precisely the reproducing kernel in the Hermitian random matrix model of size $n-j$ with weight e^{-nV} with no external source. Hence for the rank 1 spiked Hermitian model, the reproducing kernel is a rank 1 perturbation of $K_{n-j,n}$. For the higher rank spiked Hermitian model, the reproducing kernel is a rank r perturbation of $K_{n-j,n}$ which will be studied in the subsequent paper.

For the asymptotic result for this paper, we need asymptotics of $K_{n-j,n}$, ψ_{n-j} and $\tilde{\psi}_{n-j}$. The asymptotics of orthogonal polynomials and the Christoffel-Darboux kernel with respect to a varying weight have been studied extensively. Most notably precise strong asymptotics were obtained for a general class of potentials using the Deift-Zhou steepest-descent method for the associated Riemann-Hilbert problem (see [17] and [14]). We use the results of [17] extensively. However, in [17], only the case of $j = 0, 1$ are stated explicitly. For more general $j = O(1)$, the same analysis of [17] can be carried out after a few changes. These have been studied in various other papers (see e.g. [4] for the discrete orthogonal polynomials case). We summarize the necessary changes and state the explicit asymptotic formulas in Section 6 below.

The main part of this paper is the asymptotic analysis of $\tilde{\psi}_{n-j}(x; a; n)$ defined in (71). For this purpose, we first evaluate the asymptotics of $\mathbf{\Gamma}_{n-j}(a; n)$. This can be achieved in principle by plugging in the asymptotics of $\psi_{n-j}(x; n)$ in the definition (70) of $\mathbf{\Gamma}_{n-j}(a; n)$, and then evaluating the inner product asymptotically. However, the oscillatory nature of $\psi_{n-j}(x; n)$ in the support of the equilibrium measure makes it cumbersome to evaluate the inner product in this way. Instead we re-express $\mathbf{\Gamma}_{n-j}(a; n)$ in terms of a sum of integrals involving both $\psi_{n-j}(x; n)$ and its Cauchy transform (see (85)). This removes the oscillation and the asymptotic analysis becomes more straightforward. A similar trick is also used in evaluating $\int_{\mathbb{R}} K_{n-j,n}(x, y) e^{n(ay-V(y))} dy$ (see Lemma 3.1). The analysis is divided into several cases depending on the location of the critical points. Each of these cases correspond to the theorems in Introduction.

For the set $E = I_n^T$ or $E = J_n^T(x_*)$ of interest in each theorem, we can show that $\tilde{\psi}_{n-j}(x; a; n)$ is in $L^2(E)$ for each n , and hence $\chi_E \tilde{\psi}_{n-j} \otimes \psi_{n-j} \chi_E$ in (74) is a trace class operator. However, the L^2 -norm of $\tilde{\psi}_{n-j}$ is not uniformly bounded in n nor there is a simple estimate on the L^2 norm. For example, we will find that in (124) that $\tilde{\psi}_{n-j}(x; a; n) = O(\sqrt{n} e^{n(\mathbf{G}(x) + \mathbf{H}(x) - 2\mathbf{H}(c))/2})$ for $x > c$ when $a < \min\{\mathbf{a}_c, \frac{1}{2}V'(\mathbf{e})\}$. The function $\mathbf{G}(x) + \mathbf{H}(x) - 2\mathbf{H}(c) = -V(x) + 2a(x-c) + 2\mathbf{g}(c) - \ell$ tends to $-\infty$ as $x \rightarrow +\infty$, but it may be positive for some value $x > c$. This implies that we do not have a good trace norm of $\chi_E \tilde{\psi}_{n-j} \otimes \psi_{n-j} \chi_E = \|\tilde{\psi}_{n-j}\|_{L^2(E)} \|\psi_{n-j}\|_{L^2(E)}$, and we cannot compare the size of $\chi_E K_{n-j,n} \chi_E$ and $\chi_E \tilde{\psi}_{n-j} \otimes \psi_{n-j} \chi_E$. However, the rapid decay of the operator $(1 - \chi_E K_{n-j,n} \chi_E)^{-1}$ can be used to control the estimates. We proceed as follows.

From (73) and (74),

$$\begin{aligned}
\mathbb{P}_{n-j+1,n}(a; E) &= \det(1 - \chi_E \tilde{K}_{n-j+1,n} \chi_E) \\
&= \det(1 - \chi_E K_{n-j,n} \chi_E) \cdot \det(1 - (1 - \chi_E K_{n-j,n} \chi_E)^{-1} \chi_E \tilde{\psi}_{n-j} \otimes \psi_{n-j} \chi_E) \\
&= \det(1 - \chi_E K_{n-j,n} \chi_E) \cdot [1 - \langle (1 - \chi_E K_{n-j,n} \chi_E)^{-1} \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_E] \\
&= \det(1 - \chi_E K_{n-j,n} \chi_E) \cdot [1 - \langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_E \\
&\quad - \langle (1 - \chi_E K_{n-j,n} \chi_E)^{-1} \chi_E K_{n-j,n} \chi_E \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_E].
\end{aligned} \tag{75}$$

The advantage of using this formula is that $\tilde{\psi}_{n-j}$ appears in the inner product $\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_E$ and the function $(K_{n-j,n} \chi_E \tilde{\psi}_{n-j})(x)$. We will see that we have good estimates on both of these quantities due to the fast decay of $\psi_{n-j}(y; n)$ and $K_{n-j,n}(x, y)$ as $n \rightarrow \infty$ for all $y \in E$.

We study two kinds of intervals $E = I_n^T$ and $E = J_n^T(x_*)$, $x_* > \mathbf{e}$.

- (a) For $E = J_n^T(x_*)$ where x_* , which may depend on n , is in a compact subset of (\mathbf{e}, ∞) , from (346), $\det(1 - \chi_{J_n^T(x_*)} K_{n-j,n} \chi_{J_n^T(x_*)}) \rightarrow 1$. For the last inner product in (75), note that the operator norm of $(1 - \chi_{J_n^T(x_*)} K_{n-j,n} \chi_{J_n^T(x_*)})^{-1}$ is uniformly bounded from (344), and $\psi_{n-j}(x) \rightarrow 0$ in $L^2(J_n^T(x_*))$ from (333). We will show that the $L^2(J_n^T(x_*))$ norm of $K_{n-j,n} \chi_{J_n^T(x_*)} \tilde{\psi}_{n-j}$ is uniformly bounded. Hence we will have

$$\mathbb{P}_{n-j+1,n}(a; J_n^T(x_*)) = 1 - \langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{J_n^T(x_*)} + o(1) \tag{76}$$

Therefore, we need

- (i) asymptotic evaluation of $\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{J_n^T(x_*)}$,
 - (ii) uniform boundedness of $L^2(J_n^T(x_*))$ -norm of $K_{n-j,n} \chi_{J_n^T(x_*)} \tilde{\psi}_{n-j}$.
- (b) For $E = I_n^T$, from (345), $\det(1 - \chi_{I_n^T} K_{n-j,n} \chi_{I_n^T}) \rightarrow F_0(T)$. For the cases in Section 3, we will show that $K_{n-j,n} \chi_{I_n^T} \tilde{\psi}_{n-j}$ is uniformly bounded in $L^2(I_n^T)$ but for the cases in Section 5, we will see that the $L^2(I_n^T)$ -norm of $K_{n-j,n} \chi_{I_n^T} \tilde{\psi}_{n-j}$ is $O(n^{1/6})$. We here state the necessary estimates separately.

- (b1) For Sections 3, we need

- (i) asymptotic evaluation of $\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{I_n^T}$,
- (ii) uniform $L^2(I_n^T)$ boundedness of $K_{n-j,n} \chi_{I_n^T} \tilde{\psi}_{n-j}$.

Then it follows that, since the operator norm of $(1 - \chi_{I_n^T} K_{n-j,n} \chi_{I_n^T})^{-1}$ is uniformly bounded from Corollary 6.3 (c), and $\psi_{n-j} \rightarrow 0$ in $L^2(I_n^T)$ from (332), that

$$\mathbb{P}_{n-j+1,n}(a; I_n^T) = (F_0(T) + o(1)) \cdot [1 - \langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{I_n^T} + o(1)]. \tag{77}$$

- (b2) For Section 5, we need

- (i) asymptotic evaluation of $\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{I_n^T}$,

(ii) asymptotics evaluation of

$$u_{j,n}(\xi) := \frac{1}{\sqrt{n}}(K_{n-j,n}\chi_{I_n^T}\tilde{\psi}_{n-j})(\mathbf{e} + \beta^{-1}n^{-2/3}\xi) \quad (78)$$

in $L^2([T, \infty))$.

Then since $(1 - \chi_{[T,\infty)}\mathcal{K}_{n-j,n}\chi_{[T,\infty)})^{-1} \rightarrow (1 - \chi_{[T,\infty)}K_{\text{Airy}}\chi_{[T,\infty)})^{-1}$ in operator norm from (343) and $n^{-1/6}\psi_{n-j}(\mathbf{e} + \beta^{-1}n^{-2/3}\xi) - \mathcal{B}_{j,n}(\mathbf{e}) \text{Ai}(\xi) \rightarrow 0$ in $L^2([T, \infty))$ by Corollary 6.1(d), we find that if $u_{j,n} - u_n \rightarrow 0$ in $L^2([T, \infty))$, then

$$\begin{aligned} \mathbb{P}_{n-j+1,n}(a; I_n^T) &= (F_0(T) + o(1)) \cdot [1 - \langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{I_n^T} \\ &\quad - \frac{1}{\beta} \mathcal{B}_{j,n}(\mathbf{e}) \langle (1 - \chi_{[T,\infty)}K_{\text{Airy}}\chi_{[T,\infty)})^{-1}u_n, \text{Ai} \rangle_{[T,\infty)} + o(1)]. \end{aligned} \quad (79)$$

In Sections 3, 4 and 5 we only consider $a > 0$. When $a < 0$, the largest eigenvalue in the spiked source model defined by (21) has the same distribution as the negative value of the smallest eigenvalue of the spiked source model that is defined by the same formula but with the potential function $\hat{V}(x) = V(-x)$ and the external source matrix $-\mathbf{A}_{n-j+1}$. Since $\hat{V}(x)$ is regular as long as $V(x)$ is, and the non-zero eigenvalue of \mathbf{A}_{n-j+1} is positive, the analysis in this paper applies for that spiked source model. We need to keep track of the smallest eigenvalue in the new spiked source model, and it can be done in the same way that we analyze the largest one. It can be checked that the limiting distribution of the smallest eigenvalue is not affected by the positive external source eigenvalue a , corresponding to the $a < 0$ case of Theorem 1.2(a). We skip any further remarks.

3 When $0 < a < \frac{1}{2}V'(\mathbf{e})$

As outlined in Section 2, we need to show that the L^2 norm of $K_{n-j,n}\chi_{I_n^T}\tilde{\psi}_{n-j}$ is uniformly bounded in n , and need to evaluate $\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_E$ asymptotically for appropriate choices of the interval E . The key part is the asymptotic evaluation of the function $\tilde{\psi}_{n-j}(x)$. In Subsection 3.1 we first evaluate $\Gamma_{n-1}(a)$ asymptotically and then use this in Subsection 3.2 to evaluate $\tilde{\psi}_{n-j}(x)$. The remaining subsections are devoted to the proof of the main theorems in each sub-case.

3.1 Asymptotic evaluation of $\Gamma_{n-j}(a) := \Gamma_{n-j}(a; n)$

From the definition (70),

$$\Gamma_{n-j}(a) = \int_{-\infty}^c \varphi_{n-j}(y)e^{nay}dy + \int_c^{\infty} \varphi_{n-j}(y)e^{nay}dy \quad (80)$$

for any $c \in \mathbb{R}$. We take $c = c(a)$ as in Definition 1.1. Note that $c(a) > \mathbf{e}$ since $a < \frac{1}{2}V'(\mathbf{e})$. The reason that we split the integral at $y = c$ will be clear in the below, particularly the paragraph before (89).

Let

$$(Cf)(z) := \frac{1}{2\pi i} \int_{\mathbb{R}} \frac{f(y)}{y - z} dy \quad (81)$$

denote the Cauchy transform of function $f \in L^2(\mathbb{R})$. Using $C_+ - C_- = 1$ and noting that φ_{n-j} is analytic, we have

$$\int_{-\infty}^c \varphi_{n-j}(y) e^{nay} dy = \int_{-\infty}^c ((C_+ \varphi_{n-j})(y) - (C_- \varphi_{n-j})(y)) e^{nay} dy. \quad (82)$$

Note that $(C\varphi_{n-j})(z) e^{naz} \rightarrow 0$ exponentially as $\Re(z) \rightarrow -\infty$ since $a > 0$. Therefore, we can deform the contour and obtain

$$\int_{-\infty}^c \varphi_{n-j}(y) e^{nay} dy = - \int_{\Gamma_+ \cup \Gamma_-} (C\varphi_{n-j})(z) e^{naz} dz, \quad (83)$$

where, with a constant $C_\Gamma > \frac{1}{2a}$,

$$\begin{aligned} \Gamma_+ &:= \{c + it \mid 0 < t < C_\Gamma\} \cup \{c + iC_\Gamma - t \mid t \geq 0\}, \\ \Gamma_- &:= \text{complex conjugation of } \Gamma_+. \end{aligned} \quad (84)$$

The contours are oriented as indicated in Figure 5. Therefore we find

$$\mathbf{\Gamma}_{n-j}(a) = - \int_{\Gamma_+ \cup \Gamma_-} (C\varphi_{n-j})(z) e^{naz} dz + \int_c^\infty \varphi_{n-j}(y) e^{nay} dy. \quad (85)$$

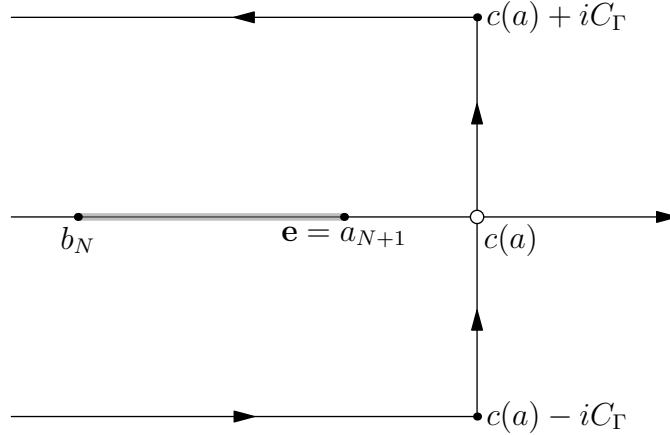


Figure 5: The contours Γ_+ and Γ_- .

Let δ be given in Proposition 6.1. Let $\epsilon < \min\{c(a) - \mathbf{e}, 2\delta\}$ be a small enough positive constant, independent of n , such that all maximizers of $\mathbf{G}(x; a)$ in $[c, \infty)$ are in the interval $(\mathbf{e} + \epsilon, \infty)$. Recall the asymptotics of $(C\varphi_{n-j})(z)$ summarized in Section 6. Since the contours Γ_\pm lie in B_δ (in Figure 10) in Section 6, from the asymptotic formula (320) for $(C\varphi_{n-j})(z)$,

$$\int_{\Gamma_+ \cup \Gamma_-} (C\varphi_{n-j})(z) e^{naz} dz = \int_{\Gamma_+ \cup \Gamma_-} \tilde{M}_{j,n}(z) e^{n(\mathbf{H}(z;a) - \ell/2)} dz. \quad (86)$$

Here we recalled the definition (32), $\mathbf{H}(z; a) := -\mathbf{g}(z) + az + \ell$.

We now use the method of steepest-descent to evaluate the integral asymptotically. By Lemma 1.1, $\mathbf{H}'(c(a); a) = 0$ and $\mathbf{H}''(c(a); a) > 0$. It is straightforward to check, with the help of the formula of $\mathbf{g}(x)$ in (14), that for $z(t) = c + it$, $t > 0$, the function $\Re \mathbf{H}(z(t); a)$ in t satisfies

$$\frac{d}{dt} \Re \mathbf{H}(z(t); a) = - \int \frac{t}{(c-s)^2 + t} \Psi(s) ds < 0. \quad (87)$$

Also for $z(t) = c + iC_\Gamma - t \in \Gamma_+$, $t \geq 0$,

$$\frac{d}{dt} \Re \mathbf{H}(z(t); a) = - \int \frac{t - c + s}{(t - c + s)^2 + C_\Gamma^2} \Psi(s) ds - a, \quad (88)$$

is negative for all $t \geq 0$ if $C_\Gamma > 1/(2a)$. Hence $\Re \mathbf{H}(z; a)$ decreases as z moves along Γ_+ counterclockwise. Similarly $\Re \mathbf{H}(z; a)$ increases as z moves along Γ_- counterclockwise. Therefore $\overline{\Gamma_+ \cup \Gamma_-}$ is a curve of steep-descent for \mathbf{H} with the saddle point at $z = c$. The fact that $z = c$ is a saddle point of \mathbf{H} is the reason that we have split the integral in (80) at c .

From Proposition 6.1(b), $\tilde{M}_{j,n} = \tilde{\mathcal{M}}_{j,n}(z)(1 + O(n^{-1}))$ and $\tilde{\mathcal{M}}_{j,n}(z)$ is analytic in B_δ . Moreover, $\tilde{M}_{j,n}$, $\tilde{M}'_{j,n}$ and $1/\tilde{M}_{j,n}$ are uniformly $O(1)$ in a neighborhood of $c(a)$ and $\tilde{M}_{j,n} = O(z^{j-1})$ uniformly in n as $z \rightarrow \infty$. Thus the method of steepest-descent can be applied to (86) and we obtain

$$\int_{\Gamma_+ \cup \Gamma_-} (C\varphi_{n-j})(z) e^{naz} dz = \frac{i\sqrt{2\pi} \tilde{\mathcal{M}}_{j,n}(c) e^{n(\mathbf{H}(c(a); a) - \ell/2)}}{\sqrt{n\mathbf{H}''(c(a); a)}} (1 + o(1)). \quad (89)$$

Now consider the second integral in (85). From the asymptotics (318) for φ_{n-j} in B_δ ,

$$\int_c^\infty \varphi_{n-j}(y) e^{nay} dy = \int_c^\infty M_{j,n}(y) e^{n(\mathbf{G}(y; a) - \ell/2)} dy \quad (90)$$

Using that $M_{j,n}$ is uniformly bounded in any compact subset of $[\mathbf{e} + \epsilon, \infty)$, and $M_{j,n}(y) = O(y^{-j})$ as $y \rightarrow \infty$, and using that $\mathbf{G}(y; a) \rightarrow -\infty$ at least linearly, we obtain the trivial estimate that

$$\int_c^\infty \varphi_{n-j}(y) e^{nay} dy = O(e^{n(\mathbf{G}_{\max}(a) - \ell/2)}) \quad (91)$$

where $\mathbf{G}_{\max}(a) := \max\{\mathbf{G}(y; a) \mid y \geq c(a)\}$. Together with (89), we obtain the following result. Recall the properties of \mathbf{G} and \mathbf{H} in Subsection 1.4.

- Suppose that $a < \mathbf{a}_c$. Then $\mathbf{H}(c; a) > \mathbf{G}_{\max}(a)$. Therefore, (89) is exponentially larger than (91) and we obtain

$$\Gamma_{n-j}(a) = \sqrt{\frac{2\pi}{n}} e^{-n\ell/2} \frac{-i\tilde{\mathcal{M}}_{j,n}(c(a))}{\sqrt{\mathbf{H}''(c(a); a)}} e^{n\mathbf{H}(c(a); a)} (1 + o(1)). \quad (92)$$

- Suppose that $\mathbf{a}_c < a < \frac{1}{2}V'(\mathbf{e})$ (assuming that V is such that $\mathbf{a}_c < \frac{1}{2}V'(\mathbf{e})$). Then $\mathbf{G}_{\max}(a) > \mathbf{H}(c; a)$ and hence (91) is exponentially larger than (89). Suppose that $a \notin \mathcal{J}_V$ and let $x_0 = x_0(a) \in (c(a), \infty)$ be the unique point $\mathbf{G}_{\max}(a)$ is attained. If $\mathbf{G}''(x_0; a) \neq 0$, using the

Laplace's method applied to (90) (using the properties of $M_{j,n}$ in Proposition 6.1 (a)), we obtain

$$\Gamma_{n-j}(a) = \sqrt{\frac{2\pi}{-n\mathbf{G}''(x_0; a)}} \mathcal{M}_{j,n}(x_0) e^{n\mathbf{G}(x_0; a) - n\ell/2} (1 + o(1)). \quad (93)$$

If $\mathbf{G}''(x_0; a) = 0$ and $\mathbf{G}^{(2k)}(x_0; a) \neq 0$, $\mathbf{G}^{(j)}(x_0; a) = 0$ for $j = 1, \dots, 2k-1$, then the Laplace's method implies that

$$\Gamma_{n-j}(a) = \left(\frac{(2k)!}{-n\mathbf{G}^{(2k)}(x_0; a)} \right)^{1/(2k)} \int_{-\infty}^{\infty} e^{-x^{2k}} dx \mathcal{M}_{j,n}(x_0) e^{n\mathbf{G}(x_0; a) - n\ell/2} (1 + o(1)). \quad (94)$$

When $a \in \mathcal{J}_V$, the contributions to (91) at each maximizer should be added. Examples of this case are in (155) and (174).

- If $a = \mathbf{a}_c$ (note that since we assumed $a < \frac{1}{2}V'(\mathbf{e})$, this implies that V is such that $\mathbf{a}_c < \frac{1}{2}V'(\mathbf{e})$), then $\mathbf{H}(c; \mathbf{a}_c) = \mathbf{G}_{\max}(\mathbf{a}_c)$. If we further assume that $\mathbf{a}_c \notin \mathcal{J}_V$ and $\mathbf{G}''(x_0(\mathbf{a}_c); \mathbf{a}_c) \neq 0$, (89) and (91) are of same order. Then by Laplace's method applied to (91),

$$\Gamma_{n-j}(\mathbf{a}_c) = \sqrt{\frac{2\pi}{n}} e^{-n\ell/2} \left(\frac{\mathcal{M}_{j,n}(x_0(\mathbf{a}_c))}{\sqrt{-\mathbf{G}''(x_0(\mathbf{a}_c))}} + \frac{-i\tilde{\mathcal{M}}_{j,n}(c(\mathbf{a}_c))}{\sqrt{\mathbf{H}''(c(\mathbf{a}_c))}} \right) e^{n\mathbf{G}(x_0(\mathbf{a}_c); \mathbf{a}_c)} (1 + o(1)). \quad (95)$$

We can consider a double scaling case when

$$a = \mathbf{a}_c + \frac{\alpha}{n} \quad (96)$$

where α is in a compact subset of \mathbb{R} . By the definition of c and x_0 , a direct computation shows that

$$\mathbf{G}(x_0(a); a) = \mathbf{G}(x_0(\mathbf{a}_c); \mathbf{a}_c) + \frac{\alpha x_0(\mathbf{a}_c)}{n} + O(n^{-2}), \quad (97)$$

$$\mathbf{H}(c(a); a) = \mathbf{H}(c(\mathbf{a}_c); \mathbf{a}_c) + \frac{\alpha c(\mathbf{a}_c)}{n} + O(n^{-2}). \quad (98)$$

This implies that

$$\Gamma_{n-j}(a) = \sqrt{\frac{2\pi}{n}} e^{-n\ell/2} \left(\frac{\mathcal{M}_{j,n}(x_0(\mathbf{a}_c))}{\sqrt{-\mathbf{G}''(x_0(\mathbf{a}_c); \mathbf{a}_c)}} + \frac{-i\tilde{\mathcal{M}}_{j,n}(c(\mathbf{a}_c))}{\sqrt{\mathbf{H}''(c(\mathbf{a}_c); \mathbf{a}_c)}} e^{\alpha(c(\mathbf{a}_c) - x_0(\mathbf{a}_c))} \right) e^{n\mathbf{G}(x_0(a); a)} (1 + o(1)). \quad (99)$$

If $\mathbf{G}''(x_0(\mathbf{a}_c)) = 0$, then the term (91) is greater than (89) by a fractional power of n (see (94)), and hence the term involving $\tilde{\mathcal{M}}_{j,n}(c(\mathbf{a}_c))$ disappears in the expression of $\Gamma_{n-j}(\mathbf{a}_c)$. On the other hand, if $\mathbf{a}_c \in \mathcal{J}_V$, then there are more than one maximizers of $\mathbf{G}(y; \mathbf{a}_c)$ making contributions in (91). We do not state the formulas explicitly here but instead state them in the appropriate subsections where they arise.

3.2 Asymptotic evaluation of $\tilde{\psi}_{n-j}(x) := \tilde{\psi}_{n-j}(x; a)$

3.2.1 Algebraic formula

From (71),

$$\Gamma_{n-j}(a)\tilde{\psi}_{n-j}(x) = e^{n(ax-V(x)/2)} - \sum_{i=0}^{n-j-1} \psi_i(x) \int_{-\infty}^{\infty} \varphi_i(y) e^{nay} dy. \quad (100)$$

This can be written in the following way. Let c be any constant such that $c > \mathbf{e}$. We will take $c = c(a)$ as in Definition 1.1 in the subsequent sections for asymptotic analysis, but the following result holds for any $c > \mathbf{e}$.

Lemma 3.1. *For $0 < a < \frac{1}{2}V'(\mathbf{e})$, we have, with Γ_{\pm} given in (84), for $x \in \mathbb{R} \setminus \{c\}$,*

$$\begin{aligned} \Gamma_{n-j}(a)\tilde{\psi}_{n-j}(x) &= e^{n(ax-V(x)/2)} 1_{(c,\infty)}(x) \\ &+ \int_{\Gamma_+ \cup \Gamma_-} \mathbb{K}_{n-j,n}(x, z) e^{naz} dz - \int_c^{\infty} K_{n-j,n}(x, y) e^{n(ay-V(y)/2)} dy \end{aligned} \quad (101)$$

where

$$\mathbb{K}_{k,n}(x, z) := \frac{\gamma_{k-1} \psi_k(x) (C\varphi_{k-1})(z) - \psi_{k-1}(x) (C\varphi_k)(z)}{\gamma_k (x - z)}, \quad x \neq z \quad (102)$$

and $K_{k,n}(x, y)$ is defined in (69).

Proof. By the same calculation that leads to (85), (100) equals

$$\begin{aligned} \Gamma_{n-j}(a; n)\tilde{\psi}_{n-j}(x) &= e^{n(ax-V(x)/2)} \\ &+ \sum_{i=0}^{n-j-1} \psi_i(x) \int_{\Gamma_+ \cup \Gamma_-} (C\varphi_i)(z) e^{naz} dz - \sum_{i=0}^{n-j-1} \psi_i(x) \int_c^{\infty} \varphi_i(y) e^{nay} dy. \end{aligned} \quad (103)$$

We exchange the sum and the integral in both terms. The second sum can be simplified by using the Christoffel-Darboux formula and becomes the last integral in (101). To analyze the first sum, we first take $x \in \mathbb{C} \setminus (\mathbb{R} \cup \Gamma_+ \cup \Gamma_-)$. By using the definition of Cauchy operator and from the Christoffel-Darboux formula, we have

$$\begin{aligned} \sum_{i=0}^{n-j-1} \psi_i(x) (C\varphi_i)(z) &= \frac{1}{2\pi i} \int_{\mathbb{R}} \frac{1}{w - z} \sum_{i=0}^{n-j-1} \psi_i(x) \varphi_i(w) dw \\ &= \frac{\gamma_{n-j-1}}{\gamma_{n-j}} \frac{1}{2\pi i} \int_{\mathbb{R}} \frac{\psi_{n-j}(x) \varphi_{n-j-1}(w) - \psi_{n-j-1}(x) \varphi_{n-j}(w)}{(w - z)(x - w)} dw. \end{aligned}$$

Using the partial fraction formula and the definition of the Cauchy transformation again, this equals

$$\begin{aligned}
& \frac{\gamma_{n-j-1}}{\gamma_{n-j}} \frac{1}{x-z} \frac{1}{2\pi i} \left[\int_{\mathbb{R}} \frac{\psi_{n-j}(x)\varphi_{n-j-1}(w) - \psi_{n-j-1}(x)\varphi_{n-j}(w)}{w-z} dw \right. \\
& \quad \left. - \int_{\mathbb{R}} \frac{\psi_{n-j}(x)\varphi_{n-j-1}(w) - \psi_{n-j-1}(x)\varphi_{n-j}(w)}{w-x} dw \right] \\
&= \frac{\gamma_{n-j-1}}{\gamma_{n-j}} \left[\frac{\psi_{n-j}(x)(C\varphi_{n-j-1})(z) - \psi_{n-j-1}(x)(C\varphi_{n-j})(z)}{x-z} \right. \\
& \quad \left. - \frac{\psi_{n-j}(x)(C\varphi_{n-j-1})(x) - \psi_{n-j-1}(x)(C\varphi_{n-j})(x)}{x-z} \right] \\
&= \frac{\gamma_{n-j-1}}{\gamma_{n-j}} \left[\frac{\psi_{n-j}(x)(C\varphi_{n-j-1})(z) - \psi_{n-j-1}(x)(C\varphi_{n-j})(z)}{x-z} \right] \\
& \quad + \frac{1}{2\pi i} \frac{e^{-nV(x)/2}}{x-z}, \tag{104}
\end{aligned}$$

where the identity (326) is used in the last line. Hence the first sum on the right-hand-side of (103) satisfies

$$\begin{aligned}
& \sum_{i=0}^{n-j-1} \psi_i(x) \int_{\Gamma_+ \cup \Gamma_-} (C\varphi_i)(z) e^{naz} dz \\
&= \frac{\gamma_{n-j-1}}{\gamma_{n-j}} \int_{\Gamma_+ \cup \Gamma_-} \frac{\psi_{n-j}(x)(C\varphi_{n-j-1})(z) - \psi_{n-j-1}(x)(C\varphi_{n-j})(z)}{x-z} e^{naz} dz \\
& \quad + \frac{e^{-nV(x)/2}}{2\pi i} \int_{\Gamma_+ \cup \Gamma_-} \frac{e^{naz}}{x-z} dz. \tag{105}
\end{aligned}$$

Note that this was proven for $x \in \mathbb{C} \setminus (\mathbb{R} \cup \Gamma_+ \cup \Gamma_-)$, but the identity holds for $x \in \mathbb{R} \setminus \{c\}$ as well by analytic continuation. The last integral equals $-2\pi i e^{nax}$ for $x \in (-\infty, c)$ and equals 0 for $x \in (c, \infty)$ by Cauchy integral formula. Therefore we obtain (101). \square

3.2.2 For $x \geq \mathbf{e} + \epsilon$:

Take $c = c(a)$ as in Definition 1.1 in the formula of Lemma 3.1. Fix $\epsilon > 0$ small enough so that $[c, \infty) \subset [\mathbf{e} + \epsilon, \infty)$.

Lemma 3.2. For $x \in [\mathbf{e} + \epsilon, \infty)$,

$$\tilde{\psi}_{n-j}(x) = e^{n(\mathbf{G}(x) - \mathbf{H}(x))/2} \left\{ \frac{1}{e^{n\ell/2} \mathbf{\Gamma}_{n-j}(a)} e^{n\mathbf{H}(x)} \mathbf{1}_{(c, \infty)}(x) + O(\sqrt{n}(1 + |x|)^{-j}) \right\} \tag{106}$$

as $n \rightarrow \infty$ and $j = O(1)$.

Proof. Fix $\epsilon' \in (0, c - \mathbf{e} - \epsilon)$. Assume that x satisfies $|x - c| \geq \epsilon'$. Noting $ax - V(x)/2 = (\mathbf{G}(x) + \mathbf{H}(x) - \ell)/2$, by Lemma 3.1 we have

$$\tilde{\psi}_{n-j}(x) = e^{n(\mathbf{G}(x) - \mathbf{H}(x))/2} \left\{ \frac{1}{e^{n\ell/2} \mathbf{\Gamma}_{n-j}(a)} e^{n\mathbf{H}(x)} \mathbf{1}_{(c, \infty)}(x) + Q_{j,n}(x) \right\}, \tag{107}$$

where

$$Q_{j,n}(x) = \frac{1}{\Gamma_{n-j}(a)} \left[\int_{\Gamma_+ \cup \Gamma_-} e^{n(\mathbf{H}(x) - \mathbf{G}(x))/2} \mathbb{K}_{n-j,n}(x, z) e^{naz} dz - \int_c^\infty e^{n(\mathbf{H}(x) - \mathbf{G}(x))/2} K_{n-j,n}(x, y) e^{n(ay - V(y)/2)} dy \right]. \quad (108)$$

From (336), the second integral over (c, ∞) in (108) is

$$O\left((1 + |x|)^{-j} \int_c^\infty e^{n(\mathbf{G}(y) - \ell/2)} (1 + |y|)^{-j} dy\right). \quad (109)$$

On the other hand, for the integral over $\Gamma_+ \cup \Gamma_-$, (318) and (320) imply that

$$\mathbb{K}_{n-j,n}(x, z) e^{naz} = \frac{(M_{j,n}(x) \tilde{M}_{j+1,n}(z) - M_{j+1,n}(x) \tilde{M}_{j,n}(z))}{x - z} e^{n(\mathbf{G}(x) - \mathbf{H}(x))/2} e^{n(\mathbf{H}(z) - \ell/2)}. \quad (110)$$

By Proposition 6.1(a)(b),

$$\frac{(M_{j,n}(x) \tilde{M}_{j+1,n}(z) - M_{j+1,n}(x) \tilde{M}_{j,n}(z))}{x - z} = O\left(\frac{(1 + |x|)^{-j} (1 + |z|)^j}{|x - z|}\right). \quad (111)$$

Thus using the fact that c is the saddle point of $\mathbf{H}(z)$, we have that the first integral over $\Gamma_+ \cup \Gamma_-$ in (108) is

$$O\left((1 + |x|)^{-j} \int_{\Gamma_+ \cup \Gamma_-} e^{n(\mathbf{H}(z) - \ell/2)} \frac{(1 + |z|)^j}{|x - z|} dz\right) = O\left((1 + |x|)^{-j} n^{-1/2} e^{n\mathbf{H}(c)}\right) \quad (112)$$

On the other hand, for $\Gamma_{n-j}(a)$, we have from (85), (86), (89) and (90) that

$$\Gamma_{n-j}(a) = O\left(n^{-1/2} e^{n\mathbf{H}(c)}\right) + O\left(\int_c^\infty e^{n(\mathbf{G}(y) - \ell/2)} (1 + |y|)^{-j} dy\right). \quad (113)$$

Note that we also have a matching lower bound. Comparing the estimate (113) of $\Gamma_{n-j}(a)$ and two estimates (109) and (112), we find that $Q_{j,n}(x) = O((1 + |x|)^{-j})$ uniformly for $x \geq \mathbf{e} + \epsilon$ if $|x - c| \geq \epsilon'$ for a positive constant ϵ' . Note that in this case the error term in (106) does not contain \sqrt{n} .

Now let x satisfy $|x - c| < \epsilon'$. In this case, we start with the formula (101) with a different choice of c . We replace c by $c \pm n^{-1/2}$ and let $(\Gamma_+)^{\pm} \cup (\Gamma_-)^{\pm}$ be a contour deformed from $\Gamma_+ \cup \Gamma_-$ by a semicircle of radius $n^{-1/2}$ to the right/left, respectively, as illustrated in Figures 6 and 7. Here we take the $+$ sign if $x - c \geq 0$ and take the $-$ sign if $x - c < 0$. Then

$$\tilde{\psi}_{n-j}(x) = e^{n(\mathbf{G}(x) - \mathbf{H}(x))/2} \left\{ \frac{1}{e^{n\ell/2} \Gamma_{n-j}(a)} e^{n\mathbf{H}(x)} 1_{(c \pm n^{-1/2}, \infty)}(x) + Q_{j,n}^{\pm}(x) \right\}, \quad (114)$$

and

$$Q_{j,n}^{\pm}(x) = \frac{1}{\Gamma_{n-j}(a)} \left[\int_{(\Gamma_+)^{\pm} \cup (\Gamma_-)^{\pm}} e^{n(\mathbf{H}(x) - \mathbf{G}(x))} \mathbb{K}_{n-j,n}(x, z) e^{naz} dz - \int_{c \mp n^{-1/2}}^\infty e^{n(\mathbf{H}(x) - \mathbf{G}(x))} K_{n-j,n}(x, y) e^{n(ay - V(y)/2)} dy \right]. \quad (115)$$

The second integral has the same estimation as that in (108). For the first integral, note that $|x - z| \geq n^{-1/2}$. Using this, and by recalling the asymptotics of a Cauchy-type integral $\int_{(\Gamma_+)^{\pm} \cup (\Gamma_-)^{\pm}} \frac{1}{x-z} e^{n\mathbf{H}(z)} dz = O(e^{n\mathbf{H}(c)})$ for such x , we find that the first integral of (115) is $O((1 + |x|)^{-j} e^{n\mathbf{H}(c)})$ instead of $O((1 + |x|)^{-j} n^{-1/2} e^{n\mathbf{H}(c)})$ in the case when $|x - c| \geq \epsilon'$. Hence we obtain (106) for $|x - c| < \epsilon$ by noting that $e^{n\mathbf{H}(c)} / (e^{n\ell} \Gamma_{n-j}) = O(\sqrt{n})$ for $|x - c| \leq n^{-1/2}$ from (113). \square

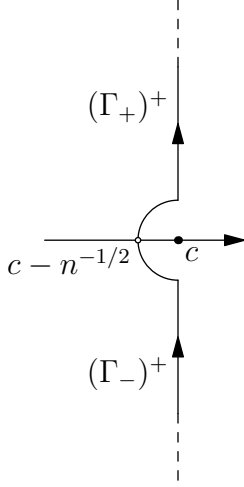


Figure 6: The deformed $(\Gamma_+)^+ \cup (\Gamma_-)^+$ for $x - c \in [0, \epsilon']$.

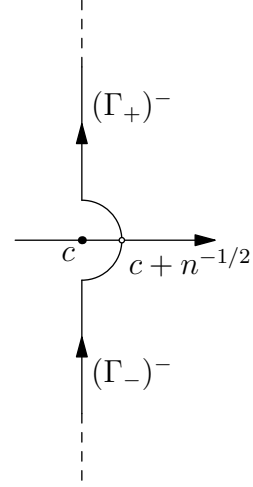


Figure 7: The deformed $(\Gamma_+)^- \cup (\Gamma_-)^-$ for $x - c \in (-\epsilon', 0)$.

3.2.3 For x near \mathbf{e} :

Let T be a fixed constant and let ϵ be a small positive constant such that $0 < \epsilon < \min\{c - \mathbf{e}, \delta_0\}$ where δ_0 is the constant in Proposition 6.1 and its corollaries in Section 6. Define the interval

$$E_{T,\epsilon} := I_n^T \setminus (\mathbf{e} + \epsilon, \infty) = [\mathbf{e} + \beta^{-1} n^{-2/3} T, \mathbf{e} + \epsilon]. \quad (116)$$

For a given $x \in E_{T,\epsilon}$, define ξ by the relation

$$x := \mathbf{e} + \beta^{-1} n^{-2/3} \xi. \quad (117)$$

Lemma 3.3. *We have for all $0 < a < V'(\mathbf{e})/2$,*

$$\tilde{\psi}_{n-j}(x) = O(n^{1/6} e^{-\frac{3}{5}|\xi|^{3/2}}), \quad (118)$$

uniformly in $x \in E_{T,\epsilon}$ and in n .

Proof. We use the formula (101). From (339) and (330), the integral over (c, ∞) in (101) is

$$O\left(n^{1/6} e^{-\frac{3}{5}|\xi|^{3/2}} \int_c^\infty e^{n(\mathbf{G}(y) - \ell/2)} (1 + |y|)^{-j} dy\right). \quad (119)$$

On the other hand, substituting (331) and (320) into (102), the integrand of the first integral over $\Gamma_+ \cup \Gamma_-$ in (101) is

$$\mathbb{K}_{n-j,n}(x, z)e^{naz} = O(n^{1/6}e^{-\frac{2}{3}|\xi|^{3/2}}(1 + |z|)^j e^{n(\mathbf{H}(z) - \ell/2)}), \quad (120)$$

for all $x \in E_{T,\epsilon}$ and $z \in \Gamma_{\pm}$ since $|x - z| > c - \mathbf{e} - \epsilon > 0$. Thus the the first integral over $\Gamma_+ \cup \Gamma_-$ in (101) is

$$O\left(n^{1/6}e^{-\frac{3}{5}|\xi|^{3/2}} \int_{\Gamma_+ \cup \Gamma_-} e^{n(\mathbf{H}(z) - \ell/2)}(1 + |z|)^j dz\right) = O\left(n^{1/6}e^{-\frac{3}{5}|\xi|^{3/2}} n^{-1/2} e^{n\mathbf{H}(c)}\right) \quad (121)$$

Substituting (119) and (121) into (101) and noting that $1_{(c,\infty)}(x) = 0$ for $x \in E_{T,\epsilon}$, we obtain

$$\mathbf{\Gamma}_{n-j}(a)\tilde{\psi}_{n-j}(x) = n^{1/6}e^{-\frac{3}{5}|\xi|^{3/2}} \left[O\left(n^{-1/2}e^{n\mathbf{H}(c)}\right) + O\left(\int_c^\infty e^{n(\mathbf{G}(y) - \ell/2)}(1 + |y|)^{-j} dy\right) \right]. \quad (122)$$

Comparing with (113) as in the previous subsection, we obtain (118). \square

3.3 Proof of Theorem 1.1(a) and Theorem 1.2(a)

Recall the outline of the proof described in Section 2. The proof proceed exactly same for both convex and non-convex potentials. The only important assumption is that $0 < a < \mathbf{a}_c$.

We first evaluate $\tilde{\psi}_{n-j}(x)$. Fix $0 < \epsilon < \delta_0$ to satisfy the conditions in Subsection 3.2 where δ_0 is the constant in Proposition 6.1 and its corollaries in Section 6. Since $0 < a < \mathbf{a}_c$, the asymptotics (92) implies that

$$\frac{1}{e^{n\ell/2}\mathbf{\Gamma}_{n-j}(a)}e^{n\mathbf{H}(x)} = O(\sqrt{n}e^{n(\mathbf{H}(x) - \mathbf{H}(c))}). \quad (123)$$

Since $\mathbf{H}(x) > \mathbf{H}(c)$ for all $x > c$ and $\mathbf{H}(x) \rightarrow \infty$ fast by Lemma 1.1, this term is larger than $O(\sqrt{n}(1 + |x|)^{-j})$. Inserting this into (106), we obtain

$$\tilde{\psi}_{n-j}(x) = \begin{cases} O(\sqrt{n}e^{n(\mathbf{G}(x) + \mathbf{H}(x) - 2\mathbf{H}(c))/2}), & x > c, \\ O(\sqrt{n}e^{n(\mathbf{G}(x) - \mathbf{H}(x))/2}), & \mathbf{e} + \epsilon \leq x \leq c. \end{cases} \quad (124)$$

On the other hand, for $x \in E_{T,\epsilon} := [\mathbf{e} + \frac{T}{\beta_{N+1}n^{2/3}}, \mathbf{e} + \epsilon]$ (see (116)), we have from Lemma 3.3 that

$$\tilde{\psi}_{n-j}(x) = O(n^{1/6}e^{-\frac{3}{5}|\xi|^{3/2}}), \quad (125)$$

where ξ is defined by (117).

Now evaluate the inner product $\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{I_n^T}$. We divide the interval I_n^T into two parts: $(\mathbf{e} + \epsilon, \infty)$ and $E_{T,\epsilon}$. From the asymptotics (124) of $\tilde{\psi}_{n-j}$ and (330) of ψ_{n-j} ,

$$\begin{aligned} \langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{(\mathbf{e} + \epsilon, \infty)} &= \int_{\mathbf{e} + \epsilon}^c O(\sqrt{n}e^{n(\mathbf{G}(x) - \mathbf{H}(x))}(1 + |x|)^{-j}) dx \\ &\quad + \int_c^\infty O(\sqrt{n}e^{n(\mathbf{G}(x) - \mathbf{H}(c))}(1 + |x|)^{-j}) dx \\ &= O(e^{-\epsilon'n}) \end{aligned} \quad (126)$$

for a constant $\epsilon' > 0$ since when $a < \mathbf{a}_c$, $\mathbf{G}(x) < \mathbf{H}(c)$ for all $x > c$ and $\mathbf{G}(x) - \mathbf{H}(x) < 0$ for all $x > \mathbf{e}$ (see Lemmas 1.1 and 1.2). On the other hand, by the asymptotics (125) of $\tilde{\psi}_{n-j}$ and (331) of ψ_{n-j} , we find, after the change of variables $x \mapsto \xi$ defined in (117), that

$$\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{E_{T,\epsilon}} = \int_T^{\beta n^{2/3}\epsilon} O(n^{1/3} e^{-\frac{6}{5}|\xi|^{3/2}}) \frac{d\xi}{\beta n^{2/3}} = O(n^{-1/3}). \quad (127)$$

Therefore,

$$\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{I_n^T} = O(n^{-1/3}). \quad (128)$$

Finally, we show the uniform boundedness of $K_{n-j,n} \chi_{I_n^T} \tilde{\psi}_{n-j}$ in $L^2(I_n^T)$. From the asymptotics of $K_{n-j,n}$ given in Corollary 6.2 and the asymptotics (124) and (125) of $\tilde{\psi}_{n-j}$, we find for $x \in E_{T,\epsilon}$ that

$$\begin{aligned} (K_{n-j,n} \chi_{I_n^T} \tilde{\psi}_{n-j})(x) &= \int_{E_{T,2\epsilon}} K_{n-j,n}(x,y) \tilde{\psi}_{n-j}(y) dy + \int_{\mathbf{e}+2\epsilon}^{\infty} K_{n-j,n}(x,y) \tilde{\psi}_{n-j}(y) dy \\ &= \int_T^{2\epsilon\beta n^{2/3}} O(n^{2/3} e^{-\frac{3}{5}|\xi|^{3/2} - \frac{3}{5}|\eta|^{3/2}} n^{1/6} e^{-\frac{3}{5}|\eta|^{3/2}}) \frac{d\eta}{\beta n^{2/3}} \\ &\quad + \int_{\mathbf{e}+2\epsilon}^c O(n^{1/6} e^{-\frac{3}{5}|\xi|^{3/2}} e^{n(\mathbf{G}(y) - \mathbf{H}(y))} (1 + |y|)^{-j}) dy \\ &\quad + \int_c^{\infty} O(n^{1/6} e^{-\frac{3}{5}|\xi|^{3/2}} e^{n(\mathbf{G}(y) - \mathbf{H}(c))} (1 + |y|)^{-j}) dy \\ &= O(n^{1/6} e^{-\frac{3}{5}|\xi|^{3/2}}) \end{aligned} \quad (129)$$

where ξ is defined by (117). Similarly, for $x \geq \mathbf{e} + \epsilon$,

$$\begin{aligned} (K_{n-j,n} \chi_{I_n^T} \tilde{\psi}_{n-j})(x) &= \int_{E_{T,\epsilon/2}} K_{n-j,n}(x,y) \tilde{\psi}_{n-j}(y) dy + \int_{\mathbf{e}+\epsilon/2}^{\infty} K_{n-j,n}(x,y) \tilde{\psi}_{n-j}(y) dy \\ &= \int_T^{\frac{1}{2}\epsilon\beta n^{2/3}} O(n^{1/6} e^{n(\mathbf{G}(x) - \mathbf{H}(x))/2} e^{-\frac{3}{5}|\eta|^{3/2}} n^{1/6} e^{-\frac{3}{5}|\eta|^{3/2}}) \frac{d\eta}{\beta n^{2/3}} \\ &\quad + \int_{\mathbf{e}+\epsilon/2}^c O(\sqrt{n} e^{n(\mathbf{G}(x) - \mathbf{H}(x))/2} e^{n(\mathbf{G}(y) - \mathbf{H}(y))}) dy \\ &\quad + \int_c^{\infty} O(\sqrt{n} e^{n(\mathbf{G}(x) - \mathbf{H}(x))/2} e^{n(\mathbf{G}(y) - \mathbf{H}(c))}) dy \\ &= O(e^{n(\mathbf{G}(x) - \mathbf{H}(x))/2}). \end{aligned} \quad (130)$$

From these, we find that

$$\|K_{n-j,n} \chi_{I_n^T} \tilde{\psi}_{n-j}\|_{L^2(I_n^T)}^2 = \int_T^{\infty} O(n^{1/3} e^{-\frac{6}{5}|\xi|^{3/2}}) \frac{d\xi}{\beta n^{2/3}} + \int_{\mathbf{e}+\epsilon}^{\infty} O(e^{n(\mathbf{G}(x) - \mathbf{H}(x))}) dx = O(n^{-1/3}). \quad (131)$$

Hence from (77), we obtain

$$\lim_{n \rightarrow \infty} \det \left(1 - \chi_{I_n^T} \tilde{K}_{n-j+1,n} \chi_{I_n^T} \right) = F_0(T). \quad (132)$$

Theorem 1.1(a) and Theorem 1.2(a) are proved.

3.4 Proof of Theorem 1.2(b) when $a < \frac{1}{2}V'(\mathbf{e})$

The proof of Theorem 1.2(b) is divided into three cases, $a < \frac{1}{2}V'(\mathbf{e})$, $a > \frac{1}{2}V'(\mathbf{e})$ and $a = \frac{1}{2}V'(\mathbf{e})$. The first case is in this subsection. The second case is in Section 4. The third case is discussed at the beginning of Section 5.

We assume that $a \in (\mathbf{a}_c, \frac{1}{2}V'(\mathbf{e}))$ and $a \notin \mathcal{J}_V$. Let $x_0 = x_0(a)$ be the unique maximizer of $\mathbf{G}(x)$ in (c, ∞) as in Lemma 1.3. We assume that $\mathbf{G}''(x_0) \neq 0$. See Remark 3.1 at the end of this subsection for a discussion when $\mathbf{G}''(x_0) = 0$ (see (59)).

Recall the definition of the interval $J_n^T(x_0)$ in (24). Note that $J_n^T(x_0) \subset (c, \infty)$. We first evaluate $\tilde{\psi}_{n-j}(x)$. By using (93), we have

$$\frac{1}{e^{n\ell/2}\Gamma_{n-j}(a)}e^{n\mathbf{H}(x)} = O(\sqrt{n}e^{n(\mathbf{H}(x)-\mathbf{G}(x_0))}), \quad x > c. \quad (133)$$

Lemma 1.1(a) and (b) imply that $\mathbf{H}(x)$ increases monotonically in $x > c$ and $\mathbf{H}(x) > \mathbf{G}(x)$ for all $x > \mathbf{e}$. Hence there exists $\epsilon' > 0$ such that $\mathbf{H}(x) > \mathbf{G}(x_0)$ for all $x > x_0 - \epsilon'$. In particular, $\mathbf{H}(x) > \mathbf{G}(x_0)$ for $x \in J_n^T(x_0)$. Therefore (106) yields, noting that $\mathbf{H}(x) \rightarrow \infty$ fast enough by Lemma 1.1(e),

$$\tilde{\psi}_{n-j}(x) = e^{n(\mathbf{G}(x)-\mathbf{H}(x))/2} \frac{e^{n\mathbf{H}(x)}}{e^{n\ell/2}\Gamma_{n-j}(a)}(1 + o(1)), \quad x \in J_n^T(x_0). \quad (134)$$

Inserting the explicit asymptotics (93) for $\Gamma_{n-j}(a)$ into (134), we have for $x > x_0 - \epsilon'$ where ϵ' is the positive constant mentioned above, and in particular for $x \in J_n^T(x_0)$ that

$$\tilde{\psi}_{n-j}(x) = \sqrt{\frac{-n\mathbf{G}''(x_0)}{2\pi}} \frac{1}{\mathcal{M}_{j,n}(x_0)} e^{n(\mathbf{G}(x)+\mathbf{H}(x)-2\mathbf{G}(x_0))/2}(1 + o(1)). \quad (135)$$

The inner product $\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{J_n^T(x_0)}$ is evaluated by using (135) and (330). For $x \in J_n^T(x_0)$,

$$\tilde{\psi}_{n-j}(x)\psi_{n-j}(x) = \sqrt{\frac{-n\mathbf{G}''(x_0)}{2\pi}} \frac{M_{j,n}(x)}{\mathcal{M}_{j,n}(x_0)} e^{n(\mathbf{G}(x)-\mathbf{G}(x_0))}(1 + o(1)). \quad (136)$$

From the assumptions for the Theorem 1.2, $\mathbf{G}(x)$ in (c, ∞) has the unique maximum at $x = x_0$ and $\mathbf{G}(x) = \mathbf{G}(x_0) + \frac{1}{2}\mathbf{G}''(x_0)(x - x_0)^2 + O(|x - x_0|^3)$ for x close to x_0 where $\mathbf{G}''(x_0) < 0$. Also $M_{j,n}(x)$, $M'_{j,n}(x)$ and $1/M_{j,n}(x)$ are bounded uniformly in n for x in a compact subset of (\mathbf{e}, ∞) and $M_{j,n}(x) = \mathcal{M}_{j,n}(x)(1 + o(1))$ from Proposition 6.1(a). Hence the standard Laplace's method applies and we obtain

$$\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{J_n^T(x_0)} = \frac{1}{\sqrt{2\pi}} \int_T^\infty e^{-\frac{1}{2}\xi^2} d\xi + o(1). \quad (137)$$

We now show that $K_{n-j,n}\chi_{J_n^T(x_0)}\tilde{\psi}_{n-j}$ is uniformly bounded in $L^2(J_n^T(x_0))$. From (135) and the part (a) of Corollary 6.2, for $x \in J_n^T(x_0)$,

$$\begin{aligned} & (K_{n-j,n}\chi_{J_n^T(x_0)}\tilde{\psi}_{n-j})(x) \\ &= \int_{J_n^T(x_0)} O(\sqrt{n}e^{n(\mathbf{G}(x)-\mathbf{H}(x))/2}(1 + |x|)^{-j}e^{n(\mathbf{G}(y)-\mathbf{G}(x_0))}(1 + |y|)^{-j})dy \\ &= O(e^{n(\mathbf{G}(x)-\mathbf{H}(x))/2}(1 + |x|)^{-j}). \end{aligned} \quad (138)$$

Therefore,

$$\|K_{n-j,n}\chi_{J_n^T(x_0)}\tilde{\psi}_{n-j}\|_{L^2(J_n^T(x_0))} = O(e^{-\epsilon'n}), \quad (139)$$

for some $\epsilon' > 0$.

Therefore, from (76), we obtain

$$\lim_{n \rightarrow \infty} \det \left(1 - \chi_{J_n^T(x_0)} \tilde{K}_{n-j+1,n} \chi_{J_n^T(x_0)} \right) = 1 - \frac{1}{\sqrt{2\pi}} \int_T^\infty e^{-\frac{1}{2}\xi^2} d\xi = G(T), \quad (140)$$

and Theorem 1.2(b) is proved.

Remark 3.1. When $\mathbf{G}''(x_0) = 0$, the Gaussian function $e^{-\frac{1}{2}\xi^2}$ in (137) is replaced by a higher-order function such as $e^{-\xi^{2k}}$ ($k > 1$). The rest of the proof is very similar. The result is the limit theorem as in (59).

3.5 Proof of Theorem 1.3(b)

Let V be a potential such that $\mathbf{a}_c < \frac{1}{2}V'(\mathbf{e})$ and $\mathbf{a}_c \notin \mathcal{J}_V$. We assume that $\mathbf{G}''(x_0(\mathbf{a}_c); \mathbf{a}_c) \neq 0$. Let

$$a = \mathbf{a}_c + \frac{\alpha}{n}, \quad (141)$$

where α is in a compact subset of \mathbb{R} .

First, consider $\tilde{\psi}_{n-j}(x)$. Note that the estimates (124) and (125) still hold. However, $\mathbf{G}(x_0(\mathbf{a}_c); \mathbf{a}_c) + \mathbf{H}(x_0(\mathbf{a}_c); \mathbf{a}_c) - 2\mathbf{H}(c(\mathbf{a}_c); \mathbf{a}_c) = 0$ at $x = x_0$, so when $a = \mathbf{a}_c$, $\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{(\mathbf{e}+\epsilon, \infty)}$ is no longer exponentially small. We need an asymptotic formula of $\tilde{\psi}(x)$ like (135). By inserting the asymptotics (99) of $\Gamma_{n-1}(a)$ into (106), similar to (135) we obtain for $x > x_0 - \epsilon'$ where ϵ' is a positive constant defined similarly as the ϵ' in (135), and in particular $x \in J_n^T$, that

$$\tilde{\psi}_{n-j}(x) = \frac{C_1(\alpha)}{C_0 + C_1(\alpha)} \sqrt{\frac{-n\mathbf{G}''(x_0(a))}{2\pi}} \frac{1}{\mathcal{M}_{j,n}(x_0(a))} e^{n(\mathbf{G}(x;a) + \mathbf{H}(x;a) - 2\mathbf{G}(x_0(a);a))/2} (1 + o(1)), \quad (142)$$

where (we omit the dependence of C_0 and $C_1(\alpha)$ on n and j to make the notations simple)

$$C_0 = \frac{-i\tilde{\mathcal{M}}_{j,n}(c(\mathbf{a}_c))}{\sqrt{\mathbf{H}''(c(\mathbf{a}_c); \mathbf{a}_c)}}, \quad (143)$$

$$C_1(\alpha) = \frac{\mathcal{M}_{j,n}(x_0(\mathbf{a}_c))}{\sqrt{-\mathbf{G}''(x_0(\mathbf{a}_c); \mathbf{a}_c)}} e^{\alpha(x_0(\mathbf{a}_c) - c(\mathbf{a}_c))}. \quad (144)$$

The constants C_0 and $C_1(\alpha)$ are positive from Proposition 6.1. If α is fixed, C_0 , $C_1(\alpha)$, C_0^{-1} and $C_1^{-1}(\alpha)$ are uniformly bounded in n . Set

$$p_{j,n}(\alpha) := \frac{C_0}{C_0 + C_1(\alpha)}. \quad (145)$$

From the definition, $p_{j,n}(\alpha)$ is a decreasing function in α , $p_{j,n}(\alpha) \rightarrow 0$ as $\alpha \rightarrow \infty$ and $p_{j,n}(\alpha) \rightarrow 1$ as $\alpha \rightarrow -\infty$ for each fixed n . Also for a fixed α , $p_{j,n}(\alpha)$ is in a compact subset of $(0, 1)$ uniformly in n . Note that when the support of the equilibrium consists of one interval, then C_0 and $C_1(\alpha)$

are independent of n , and hence so is $p_{j,n}(\alpha)$. We prove formulas (48) and (47) in Theorem 1.3 separately.

The proof of (48) is similar to that in Subsubsection 3.4. For the evaluation of $\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{J_n^T(x_0)}$, we repeat the arguments of (136)–(137). Noting that the formula of $\tilde{\psi}_{n-j}$ in (142) is the same as that in (135) except for the multiplicative factor $C_1(\alpha)/(C_0 + C_1(\alpha))$, we obtain similar to (137) that

$$\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{J_n^T(x_0)} = (1 - p_{j,n}(\alpha)) \int_T^\infty e^{-\frac{1}{2}\xi^2} \frac{d\xi}{\sqrt{2\pi}} (1 + o(1)), \quad (146)$$

where $p_{j,n}^{(0)}(\alpha)$ is defined in (145). The estimate (139) follows from the same calculations in Subsection 3.4, and we obtain from (76) that

$$\det \left(1 - \chi_{J_n^T(x_0)} \tilde{K}_{n-j+1,n} \chi_{J_n^T(x_0)} \right) = p_{j,n}(\alpha) + (1 - p_{j,n}(\alpha))G(T) + o(1), \quad (147)$$

and (48) is proved.

We now prove (47). When a is given by (141), the estimate (124) still holds. Similar to (146), we obtain by estimates (142), (124) and (118) that

$$\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{I_n^T} = (1 - p_{j,n}(\alpha)) \int_{-\infty}^\infty \frac{1}{\sqrt{2\pi}} e^{-\frac{1}{2}\xi^2} d\xi (1 + o(1)) = (1 - p_{j,n}(\alpha))(1 + o(1)). \quad (148)$$

To show that $K_{n-j,n} \chi_{I_n^T} \tilde{\psi}_{n-j}$ is uniformly bounded in $L^2(I_n^T)$, we proceed as in (129) and (130), and obtain

$$(K_{n-j,n} \chi_{I_n^T} \tilde{\psi}_{n-j})(x) = \begin{cases} O(n^{1/6} e^{-\frac{3}{5}|\xi|^{3/2}}), & x \in E_{T,\epsilon} \\ O(e^{n(\mathbf{G}(x) - \mathbf{H}(x))/2}) & x \geq \mathbf{e} + \epsilon \end{cases} \quad (149)$$

where ξ is defined by (117). Hence $\|K_{n-j,n} \chi_{I_n^T} \tilde{\psi}_{n-j}\|_{L^2(I_n^T)} = O(n^{-1/6})$. Therefore, by (77), we obtain

$$\det \left(1 - \chi_{I_n^T} \tilde{K}_{n-j+1,n} \chi_{I_n^T} \right) = p_{j,n}(\alpha)F_0(T) + o(1) \quad (150)$$

and (47) is proved.

3.6 Proof of Theorem 1.4 when $a_0 \in (\mathbf{a}_c, \frac{1}{2}V'(\mathbf{e}))$

We prove Theorem 1.4 when $a_0 < \frac{1}{2}V'(\mathbf{e})$. The case when $a_0 \geq \frac{1}{2}V'(\mathbf{e})$ will be discussed in Sections 4 and 5.

Let $a_0 \in (\mathbf{a}_c, \frac{1}{2}V'(\mathbf{e}))$ and $a_0 \in \mathcal{J}_V$. Hence a_0 is a secondary critical point. In this case, the maximum of $\mathbf{G}(x; a_0)$, $x \in (c, \infty)$, is attained at more than one point. The case when the maximum of $\mathbf{G}(x; a_0)$ is attained at more than two points can be attained by a straightforward extension and this yields (52). We omit the details in that case.

Denote the two maximizers of $\mathbf{G}(x; a_0)$ by $x_1 := x_1(a_0)$ and $x_2 := x_2(a_0)$. Let $x_1(a_0) < x_2(a_0)$. Assume that

$$\mathbf{G}''(x_1(a_0); a_0) \neq 0, \quad \mathbf{G}''(x_2(a_0); a_0) \neq 0. \quad (151)$$

The case when one of the derivative vanishes is discussed in Subsection 3.7. Let

$$a = a_0 + \frac{\alpha}{n} \quad (152)$$

where α is in a compact subset of \mathbb{R} .

First we evaluate $\tilde{\psi}_{n-j}(x)$. The asymptotics of $\mathbf{\Gamma}_{n-j}(a)$ in this case is not explicitly computed in Subsection 3.1, hence we first compute this by extending the formula (93) (see the paragraph following (94)). There are two differences from the case leading to (93) in Subsection 3.1. The first is that there are two maximizers of $\mathbf{G}(x; a_0)$ and the second is that a scales in n as in (152). The first difference simply results in adding the contributions from the both maximizers since both term are of the same order due to the condition (151). Regarding the second difference, note that since $\mathbf{G}(x; a_0)$ has maximum at $x_1(a_0)$ and $x_2(a_0)$, $\mathbf{G}(x; a)$ has two local maxima at two points, denoted by $x_1(a)$ and $x_2(a)$, which are close to $x_1(a_0)$ and $x_2(a_0)$, respectively. (Indeed, one can easily check that $x_i(a) = x_i(a_0) + \frac{\alpha}{-\mathbf{G}''(x_i(a_0))n} + O(n^{-2})$.) Using the definition of \mathbf{G} and the fact that $x_j(a_0)$ is a critical value of $\mathbf{G}(x; a_0)$, we find

$$\left. \frac{d}{da} \right|_{a=a_0} \mathbf{G}(x_i(a); a) = x_i(a_0), \quad i = 1, 2. \quad (153)$$

Hence

$$\mathbf{G}(x_i(a); a) = \mathbf{G}(x_i(a_0); a_0) + x_i(a_0) \frac{\alpha}{n} + O(n^{-2}), \quad i = 1, 2. \quad (154)$$

Therefore, as in (93) we obtain as $n \rightarrow \infty$ (note that $\mathbf{G}(x_1; a_0) = \mathbf{G}(x_2; a_0)$ and $x_i := x_i(a_0)$)

$$e^{n\ell/2} \mathbf{\Gamma}_{n-j}(a) = e^{n\mathbf{G}(x_1; a_0)} \left[\sqrt{\frac{2\pi}{-n\mathbf{G}''(x_1)}} \mathcal{M}_{j,n}(x_1) e^{x_1\alpha} + \sqrt{\frac{2\pi}{-n\mathbf{G}''(x_2)}} \mathcal{M}_{j,n}(x_2) e^{x_2\alpha} \right] (1 + o(1)). \quad (155)$$

With this asymptotics of $\mathbf{\Gamma}_{n-j}(a)$, the rest of the analysis is similar to (135), and we obtain for $x \in J_n^T(x_1)$,

$$\tilde{\psi}_{n-j}(x) = \left[\sqrt{\frac{2\pi}{-n\mathbf{G}''(x_1)}} \mathcal{M}_{j,n}(x_1) e^{x_1\alpha} + \sqrt{\frac{2\pi}{-n\mathbf{G}''(x_2)}} \mathcal{M}_{j,n}(x_2) e^{x_2\alpha} \right]^{-1} \times e^{n(\mathbf{G}(x; a) + \mathbf{H}(x; a) - 2\mathbf{G}(x_1; a_0))/2} (1 + o(1)). \quad (156)$$

We now compute the inner products $\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{J_n^T(x_i)}$, $i = 1, 2$. Using (330), (cf. (136))

$$\tilde{\psi}_{n-j}(x) \psi_{n-j}(x) = \sqrt{\frac{n}{2\pi}} \frac{\mathcal{M}_{n,j}(x) e^{n(\mathbf{G}(x; a) - \mathbf{G}(x_1; a_0))}}{A_1(\alpha) + A_2(\alpha)} (1 + o(1)) \quad (157)$$

where

$$A_i(\alpha) := \frac{\mathcal{M}_{j,n}(x_i(a_0); a_0)}{\sqrt{-\mathbf{G}''(x_i(a_0); a_0)}} e^{x_i(a_0)\alpha}, \quad i = 1, 2. \quad (158)$$

Like C_0 in (143) and $C_1(\alpha)$ in (144), $A_i(\alpha)$ is positive and is of finite distance away from 0 uniformly in n . For each $i = 1, 2$, if we set

$$x = x_i + \frac{\xi}{\sqrt{-n\mathbf{G}''(x_i)}}, \quad (159)$$

then for ξ in a compact subset of \mathbb{R} , we find using the Taylor expansion in x , and (154) that

$$\mathbf{G}(x; a) = \mathbf{G}(x_i; a_0) + x_i \frac{\alpha}{n} - \frac{\xi^2}{2n} + O(n^{-3/2}). \quad (160)$$

Thus we find

$$\tilde{\psi}_{n-j}(x)\psi_{n-j}(x) = \sqrt{\frac{n\mathbf{G}''(x_i; a_0)}{2\pi}} \frac{A_i(\alpha)}{A_1(\alpha) + A_2(\alpha)} e^{-\frac{1}{2}\xi^2} (1 + o(1)) \quad (161)$$

for x given in (159) and ξ in a compact subset of \mathbb{R} , for each $i = 1, 2$. Together with an easy estimate when x is away from $x_1(a_0)$ and $x_2(a_0)$, this implies, as in Subsection 3.4, that

$$\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{J_n^T(x_2)} = p_{j,n}^{(2)}(\alpha) \int_T^\infty \frac{1}{\sqrt{2\pi}} e^{-\frac{1}{2}\xi^2} d\xi + o(1), \quad (162)$$

$$\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{J_n^T(x_1)} = p_{j,n}^{(1)}(\alpha) \int_T^\infty \frac{1}{\sqrt{2\pi}} e^{-\frac{1}{2}\xi^2} d\xi + p_{j,n}^{(2)}(\alpha) + o(1) \quad (163)$$

for any fixed T , where for $i = 1, 2$

$$p_{j,n}^{(i)}(\alpha) = \frac{A_i(\alpha)}{A_1(\alpha) + A_2(\alpha)}. \quad (164)$$

The properties of $p_{j,n}^{(i)}(\alpha)$ stated in Theorem 1.4 can be easily checked.

The L^2 norm $\|K_{n-j,n}\chi_{J_n^T(x_0)}\tilde{\psi}_{n-j}\|_{L^2(J_n^T(x_i))}$ is estimated by the same argument as that for (139) above in Subsection 3.4 by using the asymptotics (156) of $\tilde{\psi}_{n-j}$. The result is the same exponentially decaying bound.

Therefore, by (76), we obtain

$$\det\left(1 - \chi_{J_n^T(x_1)}\tilde{K}_{n-j+1,n}\chi_{J_n^T(x_1)}\right) = p_{j,n}^{(1)}(\alpha)G(T) + o(1), \quad (165)$$

$$\det\left(1 - \chi_{J_n^T(x_2)}\tilde{K}_{n-j+1,n}\chi_{J_n^T(x_2)}\right) = p_{j,n}^{(1)}(\alpha) + p_{j,n}^{(2)}(\alpha)G(T) + o(1). \quad (166)$$

Thus Theorem 1.4 when $a < \frac{1}{2}V'(\mathbf{e})$ is proved.

3.7 Proof of Theorem 1.6 when $a \in (a_c, \frac{1}{2}V'(\mathbf{e}))$

We prove Theorem 1.6 when $a_0 < \frac{1}{2}V'(\mathbf{e})$. The case when $a_0 \geq \frac{1}{2}V'(\mathbf{e})$ will be discussed in Sections 4 and 5.

Under the assumption of Theorem 1.6, for some $k > 1$

$$\begin{aligned} \mathbf{G}''(x_1(a_0); a_0) &\neq 0, \\ \mathbf{G}^{(i)}(x_2(a_0); a_0) &= 0 \quad \text{for } i = 1, \dots, 2k-1, \\ \mathbf{G}^{(2k)}(x_2(a_0); a_0) &\neq 0. \end{aligned} \quad (167)$$

We consider the double-scaling situation when

$$a = a_0 - q \frac{\log n}{n} + \frac{\alpha}{n}, \quad \text{where } q := \frac{\frac{1}{2} - \frac{1}{2k}}{x_2(a_0) - x_1(a_0)}, \quad (168)$$

for α in a compact subset of \mathbb{R} .

The analysis is similar to Subsection 3.6. For each $i = 1, 2$, we have, as in (154),

$$\mathbf{G}(x_j(a); a) = \mathbf{G}(x_j(a_0); a_0) + x_j(a_0) \frac{-q \log n + \alpha}{n} + o(n^{-1}). \quad (169)$$

For

$$x = x_1(a_0) + \frac{\xi_1}{\sqrt{-n\mathbf{G}''(x_1(a_0); a_0)}}, \quad (170)$$

we obtain, as in (160),

$$\mathbf{G}(x; a) = \mathbf{G}(x_1(a_0); a_0) + x_1(a_0) \frac{-q \log n + \alpha}{n} - \frac{\xi_1^2}{2n} + o(n^{-1}) \quad (171)$$

for ξ_1 in a compact subset of \mathbb{R} . Similarly, for

$$x = x_2(a_0) + \left(\frac{(2k)!}{-n\mathbf{G}^{(2k)}(x_2(a_0); a_0)} \right)^{1/(2k)} \xi_2, \quad (172)$$

we have (using (167))

$$\mathbf{G}(x; a) = \mathbf{G}(x_2(a_0); a_0) + x_2(a_0) \frac{-q \log n + \alpha}{n} - \frac{\xi_2^{2k}}{n} + o(n^{-1}) \quad (173)$$

for ξ_2 in a compact subset of \mathbb{R} . Therefore, as in (155) above (cf. (94)),

$$\begin{aligned} e^{n\ell/2} \mathbf{\Gamma}_{n-j}(a) &= e^{n\mathbf{G}(x_1; a_0)} \left[\sqrt{\frac{2\pi}{-n\mathbf{G}''(x_1)}} \mathcal{M}_{j,n}(x_1) n^{-x_1 q} e^{x_1 \alpha} \right. \\ &\quad \left. + \left(\frac{(2k)!}{-n\mathbf{G}^{(2k)}(x_2; a_0)} \right)^{1/(2k)} \mathcal{M}_{j,n}(x_2) n^{-x_2 q} e^{x_2 \alpha} \int_{-\infty}^{\infty} e^{-\xi^{2k}} d\xi \right] (1 + o(1)). \end{aligned} \quad (174)$$

Since

$$\frac{1}{2} + x_1 q = \frac{1}{2k} + x_2 q = \frac{\frac{x_2}{2} - \frac{x_1}{2k}}{x_2 - x_1}, \quad (175)$$

we can write (174) as

$$e^{n\ell/2} \mathbf{\Gamma}_{n-j}(a) = n^{-\frac{\frac{x_2}{2} - \frac{x_1}{2k}}{x_2 - x_1}} e^{n\mathbf{G}(x_1; a_0)} [B_1(\alpha) + B_2(\alpha)] (1 + o(1)) \quad (176)$$

where

$$B_1(\alpha) := \sqrt{\frac{2\pi}{-\mathbf{G}''(x_1)}} \mathcal{M}_{j,n}(x_1) e^{x_1 \alpha}, \quad (177)$$

$$B_2(\alpha) := \left(\frac{(2k)!}{-\mathbf{G}^{(2k)}(x_2; a)} \right)^{1/(2k)} \mathcal{M}_{j,n}(x_2) e^{x_2 \alpha} \int_{-\infty}^{\infty} e^{-\xi^{2k}} d\xi, \quad (178)$$

As in (156), for $x \in J_n^T(x_1)$,

$$\tilde{\psi}_{n-j}(x) = n^{\frac{x_2 - x_1}{2} - \frac{x_1}{2k}} \frac{e^{n(\mathbf{G}(x;a) + \mathbf{H}(x;a) - 2\mathbf{G}(x_1(a_0);a_0))/2}}{B_1(\alpha) + B_2(\alpha)} (1 + o(1)). \quad (179)$$

From this we obtain, as in (161), that

$$\tilde{\psi}_{n-j}(x)\psi_{n-j}(x) = p_{j,n}^{(1)}(\alpha) \sqrt{\frac{-n\mathbf{G}''(x_1; a_0)}{2\pi}} e^{-\frac{1}{2}\xi_1^2} (1 + o(1)) \quad (180)$$

for x given in (170) and ξ_1 in a compact subset of \mathbb{R} , and

$$\tilde{\psi}_{n-j}(x)\psi_{n-j}(x) = p_{j,n}^{(2)}(\alpha) \left(\frac{-n\mathbf{G}^{(2k)}(x_2; a_0)}{(2k)!} \right)^{1/(2k)} \frac{e^{-\xi_2^{2k}}}{\int_{-\infty}^{\infty} e^{-\xi_2^{2k}} d\xi_2} (1 + o(1)) \quad (181)$$

for x given in (172) and ξ_1 in a compact subset of \mathbb{R} , where for $i = 1, 2$

$$p_{j,n}^{(i)}(\alpha) := \frac{B_i(\alpha)}{B_1(\alpha) + B_2(\alpha)}. \quad (182)$$

From the definition, the properties of $p_{j,n}^{(i)}(\alpha)$ in Theorem 1.6 follow easily.

Thus it follows as in (162) and (163) that

$$\langle \tilde{\psi}_{n-j}(x), \psi_{n-j}(x) \rangle_{J_n^T(x_2; k)} = p_{j,n}^{(2)}(\alpha) \frac{1}{\int_{-\infty}^{\infty} e^{-\xi^{2k}} d\xi} \int_T^{\infty} e^{-\xi^{2k}} d\xi + o(1), \quad (183)$$

$$\langle \tilde{\psi}_{n-j}(x), \psi_{n-j}(x) \rangle_{J_n^T(x_1)} = p_{j,n}^{(1)}(\alpha) \int_T^{\infty} \frac{1}{\sqrt{2\pi}} e^{-\frac{1}{2}\xi^2} d\xi + p_{j,n}^{(2)}(\alpha) + o(1). \quad (184)$$

As in the proof of Theorem 1.2(b) in Subsections 3.4 and in the proof of Theorem 1.4 in Subsection 3.6, we have $\|K_{n-j,n} \chi_{J_n^T(x_i)} \tilde{\psi}_{n-j}\|_{L^2(J_n^T(x_i))} \rightarrow 0$.

Thus we obtain, from (76), that

$$\det \left(1 - \chi_{J_n^T(x_1)} \tilde{K}_{n-j+1,n} \chi_{J_n^T(x_1)} \right) = p_{j,n}^{(1)}(\alpha) G(T) + o(1), \quad (185)$$

$$\det \left(1 - \chi_{J_n^T(x_2; k)} \tilde{K}_{n-j+1,n} \chi_{J_n^T(x_2; k)} \right) = p_{j,n}^{(1)}(\alpha) + p_{j,n}^{(2)}(\alpha) \frac{\int_{-\infty}^T e^{\xi^{2k}} d\xi}{\int_{-\infty}^{\infty} e^{\xi^{2k}} d\xi} + o(1), \quad (186)$$

and Theorem 1.6 when $a_0 < \frac{1}{2}V'(\mathbf{e})$ is proven.

4 When $a > \frac{1}{2}V'(\mathbf{e})$

Note that if $a > \frac{1}{2}V'(\mathbf{e})$, then $a > \mathbf{a}_c$. In this section, we prove Theorem 1.1(b) and Theorems 1.2(b), 1.4 and 1.6 for the case when a (or a_0) $> \frac{1}{2}V'(\mathbf{e})$. After a small change at the first step, the analysis is the same as in the case when $\mathbf{a}_c < a < \frac{1}{2}V'(\mathbf{e})$ discussed in Subsection 3.4, 3.6 and 3.7. The proof of Theorem 1.1 (b) is identical to the proof of Theorem 1.2 (b).

Note that $c(a) = \mathbf{e}$ in this case (see Definition 1.1). Since $\mathbf{G}'(\mathbf{e}) > 0$ when $a > \frac{1}{2}V'(\mathbf{e})$, $\mathbf{G}_{\max}(a) := \max\{\mathbf{G}(x; a) : x \in [\mathbf{e}, \infty)\}$ satisfies $\mathbf{G}_{\max}(a) > \mathbf{G}(\mathbf{e}) = \mathbf{H}(\mathbf{e})$. Let $\epsilon > 0$ be small enough so that all the maximizers of \mathbf{G} are in $(\mathbf{e} + 2\epsilon, \infty)$ and

$$\mathbf{G}_{\max}(a) > \mathbf{H}(\mathbf{e} + 2\epsilon) > \mathbf{H}(\mathbf{e} + \epsilon). \quad (187)$$

We have the following formula of $\tilde{\psi}_{n-j}(x)$. This is the analogue of Lemma 3.2.

Lemma 4.1. *Let $a > \frac{1}{2}V'(\mathbf{e})$. As $n \rightarrow \infty$ while $j = O(1)$,*

$$\mathbf{\Gamma}_{n-j}(a) = \int_{\mathbf{e}+\epsilon}^{\infty} \varphi_{n-j}(y) e^{n a y} dy (1 + o(1)) \quad (188)$$

and

$$\tilde{\psi}_{n-j}(x) = e^{n(\mathbf{G}(x) - \mathbf{H}(x))/2} \left\{ \frac{1}{e^{n\ell/2} \mathbf{\Gamma}_{n-j}(a)} e^{n\mathbf{H}(x)} + O((1 + |x|)^{-j}) \right\} \quad (189)$$

for $x \in [\mathbf{e} + 2\epsilon, \infty)$.

Proof. As in (85), we write $\mathbf{\Gamma}_{n-j}(a) = \mathbf{\Gamma}_{n-j}(a; n)$ as

$$\mathbf{\Gamma}_{n-j}(a) = - \int_{\bar{\Gamma}_+ \cup \bar{\Gamma}_-} (C\varphi_{n-j})(z) e^{naz} dz + \int_{\mathbf{e}+\epsilon}^{\infty} \varphi_{n-j}(y) e^{n a y} dy, \quad (190)$$

where, for a large enough but fixed positive constant $C_{\bar{\Gamma}}$, (cf. the contour Γ defined in (84))

$$\bar{\Gamma}_+ := \{\mathbf{e} + \epsilon + it \mid 0 < t \leq C_{\bar{\Gamma}}\} \cup \{\mathbf{e} + \epsilon + iC_{\bar{\Gamma}} - t \mid t \geq 0\} \quad (191)$$

and $\bar{\Gamma}_-$ is the reflected image of $\bar{\Gamma}_+$ about the real axis. The contours are oriented as indicated in Figure 8.

As in (86), by using (320) for $(C\varphi_{n-j})(z)$, the contour integral over $\bar{\Gamma}$ in (190) satisfies

$$\int_{\bar{\Gamma}_+ \cup \bar{\Gamma}_-} (C\varphi_{n-j})(z) e^{naz} dz = \int_{\bar{\Gamma}_+ \cup \bar{\Gamma}_-} \tilde{M}_{j,n}(z) e^{n(\mathbf{H}(z; a) - \ell/2)} dz. \quad (192)$$

Now $\mathbf{H}(z)$ has no saddle point in $\bar{\Gamma}_{\pm}$ since $\mathbf{H}(z) \neq 0$ for all $z \in \bar{\Gamma}$. However, it is easy to check that for $z(t) = \mathbf{e} + \epsilon + it$, $0 < t < C_{\bar{\Gamma}}$ and $z(t) = \mathbf{e} + \epsilon + iC_{\bar{\Gamma}} - t$, the formulas (87) and (88) still hold verbatim except that c becomes $\mathbf{e} + \epsilon$, provided that $C_{\bar{\Gamma}}$ is large enough, say, $C_{\bar{\Gamma}} > \frac{1}{2a}$. Hence $\Re \mathbf{H}(z)$ decreases strictly as z travels along $\bar{\Gamma}_+$ in the direction of the orientation. Similarly, $\Re \mathbf{H}(z)$ increases strictly as z travels along $\bar{\Gamma}_-$ in the direction of the orientation. Noting that for fixed n , $\tilde{M}_{j,n}(z) \rightarrow z^{j-1}$ as $z \rightarrow \infty$ and for fixed z , $\tilde{M}_{j,n}(z)$ is uniformly bounded in n from Proposition 6.1, we obtain

$$\int_{\bar{\Gamma}_+ \cup \bar{\Gamma}_-} (C\varphi_{n-j})(z) e^{naz} dz = O(e^{n\mathbf{H}(\mathbf{e}+\epsilon) - n\ell/2}). \quad (193)$$

On the other hand, consider the second integral in (190):

$$\int_{\mathbf{e}+\epsilon}^{\infty} \varphi_{n-j}(y) e^{n a y} dy = \int_{\mathbf{e}+\epsilon}^{\infty} M_{j,n}(y) e^{n(\mathbf{G}(y; a) - \ell/2)} dy. \quad (194)$$

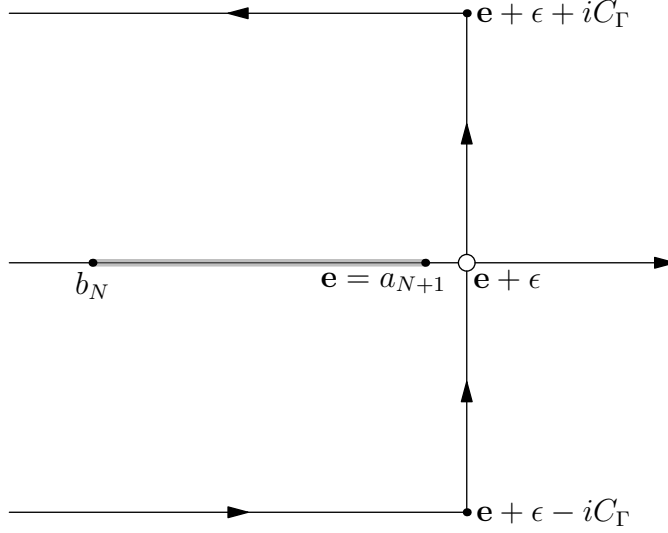


Figure 8: The contours $\bar{\Gamma}_+$ and $\bar{\Gamma}_-$

By using the Laplace's method, we find an estimate similar to (91). Hence we find that (194) is exponentially larger than (193) due to the assumption (187). Thus (188) is proven.

Now consider $\tilde{\psi}_{n-j}(x)$. Analogous to (101) in Lemma 3.1, we have, for $x \in \mathbb{R} \setminus \{\mathbf{e} + \epsilon\}$,

$$\begin{aligned} \mathbf{\Gamma}_{n-j}(a)\tilde{\psi}_{n-j}(x) &= e^{n(ax-V(x)/2)}1_{(\mathbf{e}+\epsilon,\infty)}(x) \\ &+ \int_{\bar{\Gamma}_+ \cup \bar{\Gamma}_-} \mathbb{K}_{n-j,n}(x,z)e^{naz}dz - \int_{\mathbf{e}+\epsilon}^{\infty} K_{n-j,n}(x,y)e^{n(ay-V(y)/2)}dy. \end{aligned} \quad (195)$$

Using this formula, due to the property of the $\mathbf{H}(x; a)$ on $\bar{\Gamma}_\pm$ and $\mathbf{G}(x; a)$ on $(\mathbf{e} + \epsilon, \infty)$, the analysis of the proof of Lemma 3.2 applies without any changes. If we restrict $x \geq \mathbf{e} + 2\epsilon$, then the error term $O(\sqrt{n}(1 + |x|)^{-j})$ in (106) can be replaced by $O((1 + |x|)^{-j})$ since $|x - z| \geq \epsilon$ for $z \in \bar{\Gamma}_\pm$ as in the first part of the proof of Lemma 3.1. We skip the details. \square

Since $\tilde{\psi}_{n-j}$ is the only term that depends on a and its asymptotic formula for $a > \frac{1}{2}V'(\mathbf{e})$ is same as the case when $\mathbf{a}_c < a < \frac{1}{2}V'(\mathbf{e})$ in Section 3, all the analysis in Section 3 hold without any changes. Therefore, we obtain the proof of Theorem 1.1(b) and Theorems 1.2(b), 1.4 and 1.6 for the case when a (or a_0) $> \frac{1}{2}V'(\mathbf{e})$.

5 When $a = \frac{1}{2}V'(\mathbf{e})$

First, suppose that $a = \frac{1}{2}V'(\mathbf{e}) > \mathbf{a}_c$. Then $\mathbf{G}_{\max}(a) := \max\{\mathbf{G}(x; a) \mid x \in [\mathbf{e}, \infty)\}$ satisfies $\mathbf{G}_{\max}(a) > \mathbf{H}(\mathbf{e}) = \mathbf{G}(\mathbf{e})$ (recall Definition 34 and (33)). This property is enough to prove Lemma 4.1 and the analysis of Section 4 applies without any change. Hence we obtain the proof of Theorem 1.2(b), 1.4 and 1.6 when $a = \frac{1}{2}V'(\mathbf{e}) > \mathbf{a}_c$. Combining the results of the previous two sections, we have proved all theorems except for Theorems 1.1(b), 1.3(a) and 1.5.

Theorem 1.1(b) and Theorem 1.3(a) share the same proof and this is given in Subsection 5.1. The proof of Theorem 1.5 is in Subsection 5.2.

5.1 Proof of Theorem 1.1(b) and 1.3(a)

Let V be a potential such that $\mathbf{a}_c = \frac{1}{2}V'(\mathbf{e})$. We assume that $\mathbf{a}_c \notin \mathcal{J}_V$. (This holds under the assumption of convexity of Theorem 1.1(a).) Then $\mathbf{G}(\mathbf{e}; \mathbf{a}_c) > \mathbf{G}(x; \mathbf{a}_c)$ for all $x > \mathbf{e}$. We consider a double-scaling situation when

$$a = \mathbf{a}_c + \frac{\beta\alpha}{n^{1/3}}, \quad (196)$$

where α is a real number in a compact subset of \mathbb{R} .

5.1.1 Computation of $\Gamma_{n-j}(a)$

Lemma 5.1. *We have*

$$\Gamma_{n-j}(a) = \frac{Q_n}{\beta\sqrt{n}} e^{\alpha^3/3} (\mathcal{B}_{j,n}(\mathbf{e}) + o(1)), \quad (197)$$

where

$$Q_n = Q_n(a) := e^{n(\mathbf{G}(\mathbf{e}; a) - \ell/2)}. \quad (198)$$

Here $\mathcal{B}_{j,n}(z)$ is given in (322). Note that $\mathcal{B}_{j,n}(\mathbf{e})$ is in a compact subset of $(0, \infty)$ independent of n .

Proof. In the proof of Lemma 3.1 when $a < \frac{1}{2}V'(\mathbf{e})$, we have taken the contour Γ_{\pm} to pass the point $c = c(a)$ at which $\Re\mathbf{H}(z; a)$, $z \in \Gamma_{\pm}$, takes its maximum (see (85)). Near this point, we had $\mathbf{H}(z; a) - \mathbf{H}(c(a); a) \sim \kappa(z - c(a))^2$ for some constant $\kappa > 0$. This quadratic term changes when $a = \mathbf{a}_c = \frac{1}{2}V'(\mathbf{e})$. In this case, $c(\mathbf{a}_c) = \mathbf{e}$, and (note (30))

$$\mathbf{H}'(z; \mathbf{a}_c) = -\mathbf{g}'(z) + \mathbf{a}_c = -\mathbf{g}'(z) + \mathbf{g}'(\mathbf{e}) \quad (199)$$

vanishes at $z = \mathbf{e}$. Now since $\mathbf{g}'(z)$ is the Cauchy transform of the equilibrium measure $\Psi(x)$ (see (11)), which vanishes like a square root at $x = \mathbf{e}$, we find that there is a constant $\kappa' > 0$ such that $\mathbf{H}'(z; \mathbf{a}_c) \sim \kappa'(z - \mathbf{e})^{1/2}$ for z near \mathbf{e} such that $z - \mathbf{e} \notin \mathbb{R}_-$. Hence $\mathbf{H}(z; \mathbf{a}_c) - \mathbf{H}(\mathbf{e}; \mathbf{a}_c) \sim \kappa(z - \mathbf{e})^{3/2}$ for a constant $\kappa > 0$. This 3/2-order of vanishing implies that near $z = \mathbf{e}$, $\Re\mathbf{H}(z; \mathbf{a}_c)$ decreases most rapidly in the direction of angle $2\pi/3$ and $-2\pi/3$ as z travels away from \mathbf{e} .

With the above preliminary computation in mind, we define³ the contours Σ_{\pm} as (see Figure 9)

$$\Sigma_+ = \Sigma_{2+} \cup \Sigma_{2'+} \cup \Sigma_3, \quad \Sigma_- = \Sigma_1 \cup \Sigma_{2'-} \cup \Sigma_{2-}, \quad (200)$$

³Here, the exact shape and the angle of the contour from $z = \mathbf{e}$ is not important. For example, we can use the contour that extends straightly upward from $z = \mathbf{e}$ as in Figure 5 with c replaced by \mathbf{e} . The local behavior $\mathbf{H}(z; \mathbf{a}_c) - \mathbf{H}(\mathbf{e}; \mathbf{a}_c) \sim \kappa(z - \mathbf{e})^{3/2}$ near $z = \mathbf{e}$ shows that $\Re\mathbf{H}(z)$ decays as z travels vertically away from \mathbf{e} at least locally. One can check $\Re\mathbf{H}(z)$ indeed decreases as z moves away from \mathbf{e} along on the entire curve. Our choice of the contour Σ is made for the convenience of the formulas that appear later.

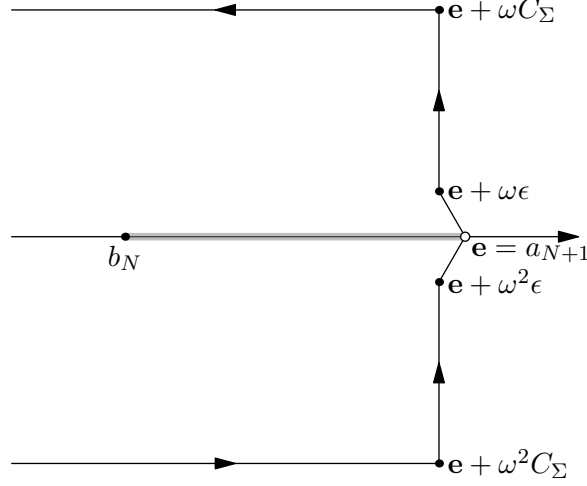


Figure 9: The contours Σ_+ and Σ_- .

and

$$\begin{aligned}
\Sigma_1 &= \{\mathbf{e} + \omega^2\epsilon - iC_\Sigma + t \mid t \leq 0\}, \\
\Sigma_{2'-} &= \{\mathbf{e} + \omega^2\epsilon + it \mid -C_\Sigma \leq t \leq 0\}, \\
\Sigma_{2-} &= \{\mathbf{e} - \omega^2t \mid -\epsilon \leq t < 0\}, \\
\Sigma_{2+} &= \{\mathbf{e} + \omega t \mid 0 < t \leq \epsilon\}, \\
\Sigma_{2'+} &= \{\mathbf{e} + \omega\epsilon + it \mid 0 \leq t \leq C_\Sigma\}, \\
\Sigma_3 &= \{\mathbf{e} + \omega\epsilon + iC_\Sigma - t \mid t \geq 0\},
\end{aligned} \tag{201}$$

where $\omega := e^{2\pi i/3}$. Here ϵ is a fixed constant chosen to satisfy the condition (215) below, and C_Σ is a positive fixed constant large enough, say, greater than $1/(2a)$. As in (85), we have

$$\mathbf{\Gamma}_{n-j}(a) = - \int_{\Sigma_+ \cup \Sigma_-} (C\varphi_{n-j})(z) e^{naz} dz + \int_{\mathbf{e}}^{\infty} \varphi_{n-j}(y) e^{nay} dy. \tag{202}$$

Let

$$\Sigma_{2\pm}^{\text{ess}} = \Sigma_{2\pm} \cap \{z \in \mathbb{C} \mid |z - \mathbf{e}| < n^{-11/21}\}. \tag{203}$$

We first consider the part of the first integral in (202) over Σ_{2+}^{ess} . Inserting the asymptotics (324) for $(C\varphi_{n-j})(z)$, there are two terms, one involving $\text{Ai}(\omega^2\Phi(z))$ and the other involving $\text{Ai}'(\omega^2\Phi(z))$. We compute each of the integrals using the change of variables $z \mapsto \xi$ defined by

$$z := \mathbf{e} + \frac{\omega}{\beta n^{2/3}} \xi. \tag{204}$$

This change of variables and the double scaling (196) imply that

$$e^{n(-\frac{1}{2}V(z)+az)} = Q_n e^{\alpha\omega\xi} (1 + o(1)) \tag{205}$$

uniformly in $z \in \Sigma_{2+}^{\text{ess}}$, where (recall Lemma 1.1(c))

$$Q_n = e^{n(-\frac{1}{2}V(\mathbf{e})+a\mathbf{e})} = e^{n(\mathbf{H}(\mathbf{e};a)-\ell/2)} = e^{n(\mathbf{G}(\mathbf{e};a)-\ell/2)}, \quad (206)$$

as defined in (198). Therefore, using the property (310) of $\Phi(z)$, and noting that $|z - \mathbf{e}| \leq n^{-11/21}$, the two integrals involving $\text{Ai}(\omega^2\Phi(z))$ and $\text{Ai}'(\omega^2\Phi(z))$ satisfy

$$\int_{\Sigma_{2+}^{\text{ess}}} \text{Ai}(\omega^2\Phi(z))B_{j,n}(z)e^{n(-\frac{1}{2}V(z)+az)}dz = \frac{\omega Q_n}{\beta n^{2/3}} \left(\mathcal{B}_{j,n}(\mathbf{e}) \int_0^\infty \text{Ai}(\xi)e^{\omega\alpha\xi}d\xi + o(1) \right) \quad (207)$$

and

$$\int_{\Sigma_{2+}^{\text{ess}}} \text{Ai}'(\omega^2\Phi(z))D_{j,n}(z)e^{n(-\frac{1}{2}V(z)+az)}dz = \frac{\omega Q_n}{\beta n^{2/3}} \left(\mathcal{D}_{j,n}(\mathbf{e}) \int_0^\infty \text{Ai}'(\xi)e^{\omega\alpha\xi}d\xi + o(1) \right). \quad (208)$$

Observe that the integrals involving $\text{Ai}(\xi)$ and $\text{Ai}'(\xi)$ are convergent as these functions decay faster than exponential functions as $\xi \rightarrow +\infty$. From these, we find that

$$\int_{\Sigma_{2+}^{\text{ess}}} (C\varphi_{n-j})(z)e^{naz}dz = -\frac{Q_n}{\beta\sqrt{n}} \left(\mathcal{B}_{j,n}(\mathbf{e}) \int_0^\infty \text{Ai}(\xi)e^{\omega\alpha\xi}d\xi + o(1) \right). \quad (209)$$

Now consider $\Sigma_{2+} \setminus \Sigma_{2+}^{\text{ess}}$. By the property (310) of $\Phi(z)$, we have that for $|z - \mathbf{e}| < 1$, there exists $c_1 > 0$ such that

$$|n^{-2/3}\Phi(z) - \beta(z - \mathbf{e})| < c_1|z - \mathbf{e}|^2. \quad (210)$$

Hence the asymptotics of $\text{Ai}(\xi)$ and $\text{Ai}'(\xi)$ as $\xi \rightarrow \infty$ ([1, 10.4.59 and 10.4.61]) imply that $z \in \Sigma_{2+}$, if $|z - \mathbf{e}| \leq \beta(1 - (3/4)^{2/3})/c_1$, then

$$\text{Ai}(\omega^2\Phi(z)) = O((1 + |\xi|)^{-1/4}e^{-\frac{1}{2}\xi^{3/2}}), \quad (211)$$

$$\text{Ai}'(\omega^2\Phi(z)) = O((1 + |\xi|)^{1/4}e^{-\frac{1}{2}\xi^{3/2}}). \quad (212)$$

Hence (324) implies that

$$(C\varphi_{n-j})(z)e^{\frac{n}{2}V(z)} = O(n^{1/6}e^{-\frac{1}{2}\xi^{3/2}}). \quad (213)$$

Also for $|z - \mathbf{e}| < 1$, there exists $c_2 > 0$ such that (cf. (205))

$$Q_n^{-1}e^{n(-\frac{1}{2}V(z)+az)} = O(e^{\alpha\omega\xi+n^{-1/3}c_2\xi^2}). \quad (214)$$

Hence if we take ϵ in (201) small enough so that

$$\epsilon < 1, \quad \epsilon < \beta c_1^{-1} \left(1 - (3/4)^{2/3} \right) \quad \text{and} \quad \epsilon < \beta^{-1} \left(\frac{1}{4} \right)^2 c_2^{-2}, \quad (215)$$

then combining (213) and (214), we have, for $n \geq (8\alpha)^{14}/\beta^7$,

$$\begin{aligned} & \int_{\Sigma_{2+} \setminus \Sigma_{2+}^{\text{ess}}} (C\varphi_{n-j})(z)e^{naz}dz \\ &= Q_n \int_{\beta n^{1/7}}^{\epsilon\beta n^{2/3}} O \left(n^{1/6} e^{-\frac{1}{2}\xi^{3/2} + \alpha\omega\xi + n^{-1/3}c_2\xi^2} \right) \frac{d\xi}{n^{2/3}} \\ &= Q_n \int_{\beta n^{1/7}}^{\epsilon\beta n^{2/3}} O \left(n^{-1/2} e^{-\frac{1}{2}\xi^{3/2} + \frac{1}{8}\xi^{3/2} + \frac{1}{4}\xi^{3/2}} \right) d\xi = Q_n O(e^{-\frac{\beta^{3/2}}{8}n^{3/14}}). \end{aligned} \quad (216)$$

For the rest of Σ_+ , by a direct calculation as in the inequalities (87) and (88), we find that $\Re(H(z; a))$ decreases strictly as z travels away from $\mathbf{e} + \omega\epsilon$ along $\Sigma_{2'+} \cup \Sigma_3$. Also by direct calculation we verify that

$$\bar{\epsilon}' := \Re \mathbf{H}(\mathbf{e}; \mathbf{a}_c) - \Re \mathbf{H}(\mathbf{e} + \omega\epsilon; \mathbf{a}_c) > 0, \quad (217)$$

where $\bar{\epsilon}'$ is a positive constant depending on ϵ . Since $\mathbf{H}(z; a) - \mathbf{H}(z; \mathbf{a}_c) = \frac{\beta\alpha}{n^{1/3}}z \rightarrow 0$ as $n \rightarrow \infty$ for a fixed z , the difference (217) with \mathbf{a}_c replaced by a is also bounded below by $\frac{1}{2}\bar{\epsilon}'$ for large enough n . Thus, from Proposition 6.1(b),

$$\begin{aligned} \int_{\Sigma_{2'+} \cup \Sigma_3} (C\varphi_{n-j})(z) e^{naz} dz &= O\left(\int_{\Sigma_{2'+} \cup \Sigma_3} |z|^{j-1} e^{n(\Re \mathbf{H}(z; a) - \ell/2)} d|z| \right) \\ &= e^{n(\mathbf{H}(\mathbf{e}; a) - \ell/2)} O(e^{-\frac{1}{4}n\bar{\epsilon}'}) \\ &= Q_n O(e^{-\frac{1}{4}n\bar{\epsilon}'}). \end{aligned} \quad (218)$$

Combining (209), (216) and (218), we obtain

$$\int_{\Sigma_+} (C\varphi_{n-j})(z) e^{naz} dz = -\frac{Q_n}{\beta\sqrt{n}} \left(\mathcal{B}_{j,n}(\mathbf{e}) \int_0^\infty \text{Ai}(\xi) e^{\omega\alpha\xi} d\xi + o(1) \right). \quad (219)$$

The integral over Σ_- can be evaluated in a similar way. Alternatively we can use the symmetry $(C\varphi_{n-j})(\bar{z}) = -\overline{(C\varphi_{n-j})(z)}$. We have

$$\int_{\Sigma_-} (C\varphi_{n-j})(z) e^{naz} dz = -\frac{Q_n}{\beta\sqrt{n}} \left(\mathcal{B}_{j,n}(\mathbf{e}) \int_0^\infty \text{Ai}(\xi) e^{\omega^2\alpha\xi} d\xi + o(1) \right). \quad (220)$$

For the integral over (\mathbf{e}, ∞) in (202), we again consider three intervals $(\mathbf{e}, \mathbf{e} + n^{-11/21}]$, $(\mathbf{e} + n^{-11/21}, \mathbf{e} + \epsilon)$ and $[\mathbf{e} + \epsilon, \infty)$, and proceed as before. We now use the asymptotics (322) for $\varphi_{n-j}(z)$ in the first two intervals and (318) for the third one. Note the similarity of (322) and (324). The calculation is similar and we obtain

$$\int_{\mathbf{e}}^\infty \varphi_{n-j}(y) e^{nay} dy = \frac{Q_n}{\beta\sqrt{n}} \left(\mathcal{B}_{j,n}(\mathbf{e}) \int_0^\infty \text{Ai}(\xi) e^{\alpha\xi} d\xi + o(1) \right). \quad (221)$$

Combining (219), (220) and (221), we find

$$\Gamma_{n-j}(a) = \frac{Q_n}{\beta\sqrt{n}} \left(\mathcal{B}_{j,n}(\mathbf{e}) \int_0^\infty \text{Ai}(\xi) (e^{\alpha\xi} + e^{\omega\alpha\xi} + e^{\omega^2\alpha\xi}) d\xi + o(1) \right). \quad (222)$$

But it is easy to check that

$$\int_0^\infty \text{Ai}(\xi) (e^{\alpha\xi} + e^{\omega\alpha\xi} + e^{\omega^2\alpha\xi}) d\xi = e^{\alpha^3/3} \quad (223)$$

for $\alpha \in \mathbb{C}$. Indeed if we denote the left-hand-side of (223) by $f(\alpha)$, then by using $\text{Ai}''(\xi) = \xi \text{Ai}(\xi)$ and $1 + \omega + \omega^2 = 0$, we find $f'(\alpha) = \alpha^2 f(\alpha)$. Recalling that $\int_0^\infty \text{Ai}(\xi) d\xi = \frac{1}{3}$ (see e.g. [1, 9.10.11]), we obtain (223). Thus we obtain (197). □

5.1.2 Evaluation of $\tilde{\psi}_{n-j}(x)$

We evaluate $\tilde{\psi}_{n-j}(x)$ for $x \in I_n^T = [\mathbf{e} + \beta^{-1}n^{-2/3}T, \infty)$ as $n \rightarrow \infty$. Let δ_0 be the constant in Section 6.

Lemma 5.2. *Let $0 < \epsilon < 2\delta_0$ be the constant in (201), satisfying the condition (215). For $x \in E_{T,\epsilon/2} := [\mathbf{e} + \beta^{-1}n^{-2/3}T, \mathbf{e} + \epsilon/2]$, we have*

$$\tilde{\psi}_{n-j}(x) = \frac{\beta\sqrt{n}}{\mathcal{B}_{j,n}(\mathbf{e})} \left[C_{-\alpha}(\xi) + e^{-\alpha^3/3} \left(\frac{1}{Q_n} e^{n(ax-V(x)/2)} - e^{\alpha\xi} \right) + o(1) \right] \quad (224)$$

where ξ is defined by the relation $x = \mathbf{e} + \beta^{-1}n^{-2/3}\xi$ as in (117), Q_n is given in (198) and $C_\alpha(\xi)$ is defined in (18). For $x \geq \mathbf{e} + \epsilon/2$, we have

$$\tilde{\psi}_{n-j}(x) = \frac{\beta\sqrt{n}e^{-\alpha^3/3}}{\mathcal{B}_{j,n}(\mathbf{e})} e^{n(\mathbf{G}(x;a) + \mathbf{H}(x;a) - 2\mathbf{G}(\mathbf{e};a))/2} (1 + o(1)). \quad (225)$$

Note that $\xi \in [T, \infty)$. Let $\bar{C} < T$ be a real number, and set $C_n := \beta^{-1}n^{-2/3}\bar{C}$. For $x \in I_n^T$ (hence $x > \mathbf{e} + C_n$), we have, as in Lemma 3.1,

$$\begin{aligned} \mathbf{\Gamma}_{n-j}(a)\tilde{\psi}_{n-j}(x) &= e^{n(ax-V(x)/2)} + \int_{(\Sigma_+ + C_n) \cup (\Sigma_- + C_n)} \mathbb{K}_{n-j,n}(x, z) e^{naz} dz \\ &\quad - \int_{\mathbf{e} + C_n}^{\infty} K_{n-j,n}(x, y) e^{n(ay-V(y)/2)} dy. \end{aligned} \quad (226)$$

Here $\Sigma_\pm + C_n$ denotes the contour Σ_\pm translated by C_n . For example, $\Sigma_+ + C_n = (\Sigma_{2+} + C_n) \cup (\Sigma_{2'+} + C_n) \cup (\Sigma_3 + C_n)$, cf. (200). We divide the proof of Lemma 5.2 into two parts.

Proof of (224). First we consider the integral over $\Sigma_+ + C_n$ in (226). For $x \in E_{T,\epsilon/2}$ and $z \in \Sigma_{2+} + C_n$, from (322) and (324),

$$\begin{aligned} \mathbb{K}_{n-j,n}(x, z) &= \left[n^{1/3} U_1(x, z) \text{Ai}(\Phi(x)) \text{Ai}(\omega^2\Phi(z)) \right. \\ &\quad + V_1(x, z) \text{Ai}(\Phi(x)) \omega^2 \text{Ai}'(\omega^2\Phi(z)) \\ &\quad + V_2(x, z) \text{Ai}'(\Phi(x)) \text{Ai}(\omega^2\Phi(z)) \\ &\quad \left. + n^{-1/3} U_2(x, z) \text{Ai}'(\Phi(x)) \omega^2 \text{Ai}'(\omega^2\Phi(z)) \right] e^{\pi i/3} e^{-nV(z)/2}, \end{aligned} \quad (227)$$

where

$$\begin{aligned} U_1(x, z) &= \frac{\gamma_{n-j-1} B_{j,n}(x) B_{j+1,n}(z) - B_{j+1,n}(x) B_{j,n}(z)}{\gamma_{n-j} (x-z)} \\ V_1(x, z) &= \frac{\gamma_{n-j-1} B_{j,n}(x) D_{j+1,n}(z) - B_{j+1,n}(x) D_{j,n}(z)}{\gamma_{n-j} (x-z)} =: \frac{W_1(x, z)}{x-z} \\ V_2(x, z) &= \frac{\gamma_{n-j-1} D_{j,n}(x) B_{j+1,n}(z) - D_{j+1,n}(x) B_{j,n}(z)}{\gamma_{n-j} (x-z)} =: \frac{W_2(x, z)}{x-z} \\ U_2(x, z) &= \frac{\gamma_{n-j-1} D_{j,n}(x) D_{j+1,n}(z) - D_{j+1,n}(x) D_{j,n}(z)}{\gamma_{n-j} (x-z)}. \end{aligned} \quad (228)$$

We now show that the main contribution to $\mathbb{K}_{n-j,n}(x, z)$ comes from the middle two terms on the right-hand side of (227). Clearly, $U_1(x, z) = O(1)$ and $U_2(x, z) = O(1)$. Since $W_1(x, x) = 1 = -W_2(x, x)$ from (329), it follows that

$$\begin{aligned} & V_1(x, z) \operatorname{Ai}(\Phi(x)) \omega^2 \operatorname{Ai}'(\omega^2 \Phi(z)) + V_2(x, z) \operatorname{Ai}'(\Phi(x)) \operatorname{Ai}(\omega^2 \Phi(z)) \\ &= \frac{\operatorname{Ai}(\Phi(x)) \omega^2 \operatorname{Ai}'(\omega^2 \Phi(z)) - \operatorname{Ai}'(\Phi(x)) \operatorname{Ai}(\omega^2 \Phi(z))}{x - z} \\ &+ \frac{W_1(x, z) - W_1(x, x)}{x - z} \operatorname{Ai}(\Phi(x)) \omega^2 \operatorname{Ai}'(\omega^2 \Phi(z)) \\ &+ \frac{W_2(x, z) - W_2(x, x)}{x - z} \operatorname{Ai}'(\Phi(x)) \operatorname{Ai}(\omega^2 \Phi(z)). \end{aligned} \quad (229)$$

Observe that $\frac{W_i(x, z) - W_i(x, x)}{x - z} = O(1)$, $i = 1, 2$. For

$$x = \mathbf{e} + \frac{\xi}{\beta n^{2/3}} \in E_{T, \epsilon/2}, \quad z = \mathbf{e} + \frac{\eta}{\beta n^{2/3}} \in \Sigma_{2+} + C_n, \quad (230)$$

using the estimates (211) and (212) for $\operatorname{Ai}(\omega^2 \Phi(z))$ and $\operatorname{Ai}'(\omega^2 \Phi(z))$, and analogous estimates for $\operatorname{Ai}(\Phi(x))$ and $\operatorname{Ai}'(\Phi(x))$, we find that (noting that the ξ in (211) and (212) are slightly different from the ξ and η in (230))

$$\mathbb{K}_{n-j,n}(x, z) = \left[\beta n^{2/3} \frac{\omega^2 \operatorname{Ai}(\xi) \operatorname{Ai}'(\omega^2 \eta) - \operatorname{Ai}'(\xi) \operatorname{Ai}(\omega^2 \eta)}{\xi - \eta} + O(n^{1/3} e^{-\frac{1}{2}(|\xi|^{3/2} + |\eta|^{3/2})}) \right] e^{\pi i/3} e^{-nV(z)/2} \quad (231)$$

for x and z in (230). For $z \in (\Sigma_{2'+} + C_n) \cup (\Sigma_{3+} + C_n)$, noting that $|x - z| \geq 1/\epsilon$, a straightforward calculation using (322) and (320) implies that

$$\mathbb{K}_{n-j,n}(x, z) = O(n^{1/6} e^{-\frac{1}{2}|\xi|^{3/2}} (1 + |z|)^j e^{n(-\mathbf{g}(z) + \ell/2)}). \quad (232)$$

Hence we obtain, by noting (205), (214) and the condition (215) satisfied by ϵ , using $-\mathbf{g}(z) + \ell/2 + az = \mathbf{H}(z; a) - \ell/2$ and noting the calculation in (218), that

$$\begin{aligned} & \int_{\Sigma_+ + C_n} \mathbb{K}_{n-j,n}(x, z) e^{naz} dz \\ &= Q_n \left[e^{\frac{\pi i}{3}} \int_{\bar{C}}^{\omega \cdot \infty} \frac{\operatorname{Ai}(\xi) \operatorname{Ai}'(\omega^2 \eta) \omega^2 - \operatorname{Ai}'(\xi) \operatorname{Ai}(\omega^2 \eta)}{\xi - \eta} e^{\alpha \eta} d\eta + O(n^{-1/3} e^{-\frac{1}{2}|\xi|^{3/2}}) \right]. \end{aligned} \quad (233)$$

The estimates for the integral over the contour $\Sigma_- + C_n$ can be obtained either by Schwarz reflection principle or by a similar calculation. We find

$$\begin{aligned} & \int_{\Sigma_- + C_n} \mathbb{K}_{n-j,n}(x, z) e^{naz} dz \\ &= Q_n \left[-e^{\frac{\pi i}{3}} \int_{\omega^2 \cdot \infty}^{\bar{C}} \frac{\operatorname{Ai}(\xi) \operatorname{Ai}'(\omega \eta) - \omega^2 \operatorname{Ai}'(\xi) \operatorname{Ai}(\omega \eta)}{\xi - \eta} e^{\alpha \eta} d\eta + O(n^{-1/3} e^{-\frac{1}{2}|\xi|^{3/2}}) \right]. \end{aligned} \quad (234)$$

For the integral over $(\mathbf{e} + C_n, \infty)$ in (226), we need asymptotics of $K_{n-j,n}(x, y)$. For $x \in E_{T, \epsilon/2}$ and $y \in (\mathbf{e} + C_n, \mathbf{e} + \epsilon)$, setting $x = \mathbf{e} + \beta^{-1}n^{-2/3}\xi$ and $y = \mathbf{e} + \beta^{-1}n^{-2/3}\eta$, we have

$$K_{n-j,n}(x, y) = \beta n^{2/3} \frac{\text{Ai}(\xi) \text{Ai}'(\eta) - \text{Ai}'(\xi) \text{Ai}(\eta)}{\xi - \eta} + O(n^{1/3} e^{-\frac{1}{2}(|\xi|^{3/2} + |\eta|^{3/2})}). \quad (235)$$

This follows from the analysis similar to that of (231). A weaker estimate is in (341), which is actually enough for our purpose. Hence for $x \in E_{T, \epsilon/2}$,

$$\begin{aligned} & \int_{\mathbf{e} + C_n}^{\mathbf{e} + \epsilon} K_{n-j,n}(x, y) e^{n(ay - \frac{1}{2}V(y))} dy \\ &= Q_n \left[\int_{\bar{C}}^{\infty} \frac{\text{Ai}(\xi) \text{Ai}'(\eta) - \text{Ai}'(\xi) \text{Ai}(\eta)}{\xi - \eta} e^{\alpha \eta} d\eta + O(n^{-1/3} e^{-\frac{1}{2}|\xi|^{3/2}}) \right]. \end{aligned} \quad (236)$$

Now consider the integral over $(\mathbf{e} + \epsilon, \infty)$. By the estimate (339) of $K_{n-j,n}(x, y)$ for $x \in E_{T, \epsilon/2}$ and $y \in (\mathbf{e} + \epsilon, \infty)$ and the identity

$$-\frac{1}{2}V(y) + ay - \frac{1}{n} \log Q_n = \frac{\mathbf{G}(y; a) + \mathbf{H}(y; a) - 2\mathbf{G}(\mathbf{e}; y)}{2}, \quad (237)$$

we obtain

$$\begin{aligned} & \int_{\mathbf{e} + \epsilon}^{\infty} K_{n-j,n}(x, y) e^{n(ay - \frac{1}{2}V(y))} dy \\ &= Q_n \int_{\mathbf{e} + \epsilon}^{\infty} O\left(n^{1/6} e^{-\frac{3}{5}|\xi|^{3/2}} (1 + |y|)^{-j} e^{n(\mathbf{G}(y; a) - \mathbf{G}(\mathbf{e}; a))}\right) dy = Q_n O(e^{-\frac{3}{5}|\xi|^{3/2} - \frac{\epsilon'}{2}n}), \end{aligned} \quad (238)$$

where

$$\epsilon' = \min_{y \geq \mathbf{e} + \epsilon} \mathbf{G}(\mathbf{e}; a) - \mathbf{G}(y; a) > 0. \quad (239)$$

Combining (236) and (238), we have

$$\begin{aligned} & \int_{\mathbf{e} + C_n}^{\infty} K_{n-j,n}(x, y) e^{n(ay - \frac{1}{2}V(y))} dy \\ &= Q_n \left[\int_{\bar{C}}^{\infty} \frac{\text{Ai}(\xi) \text{Ai}'(\eta) - \text{Ai}'(\xi) \text{Ai}(\eta)}{\xi - \eta} e^{\alpha \eta} d\eta + O(n^{-1/3} e^{-\frac{3}{5}|\xi|^{3/2}}) \right]. \end{aligned} \quad (240)$$

We now insert the results (233), (234) and (240) into (226) and evaluate $\tilde{\psi}_{n-j}(x)$ for $x \in E_{T, \epsilon}$. The combination of the three integrals can be simplified. For this purpose, note that

$$\frac{\text{Ai}(\xi) \text{Ai}'(\eta) - \text{Ai}'(\xi) \text{Ai}(\eta)}{\xi - \eta} = \int_0^{\infty} \text{Ai}(\xi + t) \text{Ai}(\eta + t) dt, \quad (241)$$

$$\frac{\omega^2 \text{Ai}(\xi) \text{Ai}'(\omega^2 \eta) - \text{Ai}'(\xi) \text{Ai}(\omega^2 \eta)}{\xi - \eta} = \int_0^{\infty} \text{Ai}(\xi + t) \text{Ai}(\omega^2(\eta + t)) dt, \quad (242)$$

$$\frac{\omega \text{Ai}(\xi) \text{Ai}'(\omega \eta) - \text{Ai}'(\xi) \text{Ai}(\omega \eta)}{\xi - \eta} = \int_0^{\infty} \text{Ai}(\xi + t) \text{Ai}(\omega(\eta + t)) dt, \quad (243)$$

where we require $\Re\eta < \Re\xi$ in (242) and (242). This can be verified by using $\text{Ai}''(z) = z \text{Ai}(z)$ and the asymptotics [1, 10.4.59 and 10.4.61] of $\text{Ai}(z)$ and $\text{Ai}'(z)$ as $z \rightarrow \infty$. Using the above Airy function identities, we find that

$$\begin{aligned} \int_{\bar{C}}^{\infty} \frac{\text{Ai}(\xi) \text{Ai}'(\eta) - \text{Ai}'(\xi) \text{Ai}(\eta)}{\xi - \eta} e^{\alpha\eta} d\eta &= \int_{\bar{C}}^{\infty} \int_0^{\infty} \text{Ai}(\xi + t) \text{Ai}(\eta + t) dt e^{\alpha\eta} d\eta \\ &= \int_0^{\infty} \text{Ai}(\xi + t) e^{-\alpha t} \left(\int_{\bar{C}+t}^{\infty} \text{Ai}(\bar{\eta}) e^{\alpha\bar{\eta}} d\bar{\eta} \right) dt. \end{aligned} \quad (244)$$

Similarly,

$$\int_{\bar{C}}^{\omega \cdot \infty} \frac{\omega^2 \text{Ai}(\xi) \text{Ai}'(\omega^2 \eta) - \text{Ai}'(\xi) \text{Ai}(\omega^2 \eta)}{\xi - \eta} e^{\alpha\eta} d\eta = \int_0^{\infty} \text{Ai}(\xi + t) e^{-\alpha t} \left(\int_{\bar{C}+t}^{\omega \cdot \infty} \text{Ai}(\omega^2 \bar{\eta}) e^{\alpha\bar{\eta}} d\bar{\eta} \right) dt, \quad (245)$$

and

$$\int_{\omega^2 \cdot \infty}^{\bar{C}} \frac{\text{Ai}(\xi) \text{Ai}'(\omega\eta) - \omega^2 \text{Ai}'(\xi) \text{Ai}(\omega\eta)}{\xi - \eta} e^{\alpha\eta} d\eta = \omega^2 \int_0^{\infty} \text{Ai}(\xi + t) e^{-\alpha t} \left(\int_{\omega^2 \cdot \infty}^{\bar{C}+t} \text{Ai}(\omega\bar{\eta}) e^{\alpha\bar{\eta}} d\bar{\eta} \right) dt. \quad (246)$$

From these results and (226), (233), (234) and (240), we find that for $x \in E_{T,\epsilon/2}$,

$$\begin{aligned} \mathbf{\Gamma}_{n-j}(a) \tilde{\psi}_{n-j}(x) &= e^{n(ax-V(x)/2)} - Q_n \left[\int_0^{\infty} \text{Ai}(\xi + t) e^{-\alpha t} \left(\int_{\bar{C}+t}^{\infty} \text{Ai}(\bar{\eta}) e^{\alpha\bar{\eta}} d\bar{\eta} \right. \right. \\ &\quad \left. \left. + \omega^2 \int_{\bar{C}+t}^{\omega \cdot \infty} \text{Ai}(\omega^2 \bar{\eta}) e^{\alpha\bar{\eta}} d\bar{\eta} + \omega \int_{\bar{C}+t}^{\omega^2 \cdot \infty} \text{Ai}(\omega\bar{\eta}) e^{\alpha\bar{\eta}} d\bar{\eta} \right) dt + o(1) \right]. \end{aligned} \quad (247)$$

Now the sum of three integrals inside the parentheses equals $e^{\alpha^3/3}$ (cf. (223)), for all t . In order to see this, first note that the sum is independent of t since its derivative with respect to t equals 0 from the Airy function identity

$$\text{Ai}(z) + \omega \text{Ai}(\omega z) + \omega^2 \text{Ai}(\omega^2 z) = 0, \quad z \in \mathbb{C}. \quad (248)$$

Then set $t = 0$ and call the sum $S(\alpha)$. Taking the derivative of $S(\alpha)$ with respect to α and using (248) and the differential equation for the Airy function, we find that $S'(\alpha) = \alpha^2 S(\alpha)$. Now by noting that $S(0) = 1$ since $\int_0^{\infty} \text{Ai}(\bar{\eta}) d\bar{\eta} = 1/3$, we obtain that $S(\alpha) = e^{\alpha^3/3}$. Hence

$$\mathbf{\Gamma}_{n-j}(a) \tilde{\psi}_{n-j}(x) = Q_n \left[\frac{1}{Q_n} e^{ax - \frac{1}{2}V(x)} - e^{\alpha^3/3} \int_0^{\infty} \text{Ai}(\xi + t) e^{-\alpha t} dt + o(1) \right] \quad (249)$$

uniformly for $x \in E_{T,\epsilon/2}$. Therefore using (197) we find that

$$\begin{aligned} \tilde{\psi}_{n-j}(x) &= \frac{\beta\sqrt{n}}{\mathcal{B}_{j,n}(\mathbf{e})} \left[e^{-\alpha^3/3} \frac{1}{Q_n} e^{ax - \frac{1}{2}V(x)} - \int_0^{\infty} \text{Ai}(\xi + t) e^{-\alpha t} dt + o(1) \right] \\ &= \frac{\beta\sqrt{n}}{\mathcal{B}_{j,n}(\mathbf{e})} \left[C_{-\alpha}(\xi) + e^{-\alpha^3/3} \left(\frac{1}{Q_n} e^{ax - \frac{1}{2}V(x)} - e^{\alpha\xi} \right) + o(1) \right] \end{aligned} \quad (250)$$

uniformly for $x = \mathbf{e} + \beta^{-1}n^{-2/3}\xi \in E_{T,\epsilon/2}$. In the last line, we used the identity

$$e^{-\alpha^3/3}e^{\alpha\xi} - \int_0^\infty \text{Ai}(\xi+t)e^{-\alpha t}dt = C_{-\alpha}(\xi). \quad (251)$$

□

Proof of (225). Let $x \geq \mathbf{e} + \epsilon/2$. Using (338) and (214), a straightforward estimate implies that

$$\frac{1}{Q_n} \int_{\mathbf{e}+C_n}^{\mathbf{e}+\epsilon/4} K_{n-j,n}(x,y)e^{n(ay-\frac{1}{2}V(y))}dy = O(n^{-1/2}(1+|x|)^{-j}e^{n(\mathbf{G}(x;a)-\mathbf{H}(x;a))/2}). \quad (252)$$

Similarly, we obtain using (336)

$$\frac{1}{Q_n} \int_{\mathbf{e}+\epsilon/4}^\infty K_{n-j,n}(x,y)e^{n(ay-\frac{1}{2}V(y))}dy = O((1+|x|)^{-j}e^{n(\mathbf{G}(x;a)-\mathbf{H}(x;a))/2}). \quad (253)$$

For the integral on $\Sigma_\pm + C_n$, the calculation is easier than the proof of (224) since $|x-z| \geq \epsilon/2$. Straightforward estimates using Proposition 6.1 imply that

$$\frac{1}{Q_n} \int_{\Sigma_\pm + C_n} \mathbb{K}_{n-j,n}(x,z)e^{naz}dz = O(n^{-1/2}(1+|x|)^{-j}e^{n(\mathbf{G}(x;a)-\mathbf{H}(x;a))/2}). \quad (254)$$

Therefore, from (226),

$$\begin{aligned} \frac{1}{Q_n} \mathbf{\Gamma}_{n-j}(a)\tilde{\psi}_{n-j}(x) &= \frac{1}{Q_n} e^{n(-\frac{1}{2}V(x)+ax)} + O((1+|x|)^{-j}e^{n(\mathbf{G}(x;a)-\mathbf{H}(x;a))}) \\ &= e^{n(\mathbf{G}(x;a)+\mathbf{H}(x;a)-2\mathbf{G}(\mathbf{e};a))/2} (1 + O((1+|x|)^{-j}e^{-n(\mathbf{H}(x;a)-\mathbf{G}(x;a))})) \\ &= e^{n(\mathbf{G}(x;a)+\mathbf{H}(x;a)-2\mathbf{G}(\mathbf{e};a))/2} (1 + o(1)). \end{aligned} \quad (255)$$

Hence using (197), we obtain (225).

□

5.1.3 Proof

Recall the outline (b2) in Section 2.

From Lemma 5.2 and Proposition 6.1, and using (214) to estimate $\frac{1}{Q_n}e^{n(ax-V(x)/2)} - e^{\alpha\xi}$, we obtain

$$\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{I_n^T} \rightarrow \int_T^\infty C_{-\alpha}(\xi) \text{Ai}(\xi) d\xi. \quad (256)$$

Now we evaluate $u_{j,n}(\xi) := \frac{1}{\sqrt{n}}(K_{n-j,n}\chi_{I_n^T}\tilde{\psi}_{n-j})(\mathbf{e} + \beta^{-1}n^{-2/3}\xi)$ defined in (78) asymptotically in $L^2([T, \infty))$. Using Lemma 5.2, and then estimating as in Subsubsection 5.1.2, we obtain

$$\int_{I_n^T} K_{n-j,n}(x,y)\tilde{\psi}_{n-j}dy = \frac{\beta\sqrt{n}}{\mathcal{B}_{j,n}(\mathbf{e})} (K_{\text{Airy}}\chi_{[T,\infty)}C_{-\alpha})(\xi) + O(n^{-1/3}e^{-\frac{1}{2}|\xi|^{3/2}}) \quad (257)$$

for $x = \mathbf{e} + \beta n^{2/3} \xi \in E_{T, \epsilon/2}$. For $x \geq \mathbf{e} + \epsilon/2$,

$$\int_{I_n^T} K_{n-j, n}(x, y) \tilde{\psi}_{n-j} dy = O((1 + |x|)^{-j} e^{n(\mathbf{G}(x; a) - \mathbf{H}(x; a))/2}). \quad (258)$$

The calculation is similar to Subsubsection 5.1.2 and we skip the details. Thus

$$u_{j, n} - \frac{\beta}{\mathcal{B}_{j, n}(\mathbf{e})} (K_{\text{Airy}} \chi_{[T, \infty)} C_{-\alpha}) \rightarrow 0, \quad \text{in } L^2[T, \infty). \quad (259)$$

Therefore, from (79),

$$\begin{aligned} \mathbb{P}_{n-j+1, n}(a; I_n^T) &= F_0(T) \cdot (1 - \langle C_{-\alpha}, \text{Ai} \rangle_{[T, \infty)}) \\ &\quad - \langle (1 - \chi_{[T, \infty)} K_{\text{Airy}} \chi_{[T, \infty)})^{-1} K_{\text{Airy}} \chi_{[T, \infty)} C_{-\alpha}, \text{Ai} \rangle_{[T, \infty)} \\ &= F_1(T; -\alpha). \end{aligned} \quad (260)$$

Hence Theorem 1.1(b) and Theorem 1.3(a) are proved.

5.2 Proof of Theorem 1.5

Note that $\mathbf{G}(x_0(\mathbf{a}_c); \mathbf{a}_c) = \mathbf{G}(\mathbf{e}; \mathbf{a}_c) > \mathbf{G}(x; \mathbf{a}_c)$ for all $x \in (\mathbf{e}, \infty) \setminus \{x_0(\mathbf{a}_c)\}$. For a given in either (54) or (56), let $x_0(a)$ be the point near $x_0(\mathbf{a}_c)$ such that $\mathbf{G}(x; a)$ achieves its local maximum. The point $x_0(a)$ is well defined as long as $|a - \mathbf{a}_c|$ is small enough. Note that for $a \geq \mathbf{a}_c$, $x_0(a)$ is same as in the definition of $x_0(a)$ in Lemma 1.3. However, for $a < \mathbf{a}_c$, $x_0(a)$ is not defined in Lemma 1.3. We extend the definition of $x_0(a)$ here for $a < \mathbf{a}_c$ when $a - \mathbf{a}_c$ is small enough.

5.2.1 Proof of Theorem 1.5(a)

We consider the double scaling situation

$$a = \mathbf{a}_c + \frac{\beta \alpha}{n^{1/3}} \quad (261)$$

where a is in a compact subset of $(-\infty, 0)$. Since we assume $G''(x_0(\mathbf{a}_c); \mathbf{a}_c) \neq 0$, we have $x_0(a) - x_0(\mathbf{a}_c) = O(|a - \mathbf{a}_c|) = O(n^{-1/3})$. We also have, as in (97), using $\frac{\partial}{\partial a} \mathbf{G}(x; a) = x$,

$$\begin{aligned} \mathbf{G}(x_0(a); a) &= \mathbf{G}(x_0(\mathbf{a}_c); \mathbf{a}_c) + x_0(\mathbf{a}_c)(a - \mathbf{a}_c) + O(|a - \mathbf{a}_c|^2) \\ &= \mathbf{G}(x_0(\mathbf{a}_c); \mathbf{a}_c) + \frac{\beta \alpha}{n^{1/3}} x_0(\mathbf{a}_c) + O(|a - \mathbf{a}_c|^2). \end{aligned} \quad (262)$$

Hence since $\mathbf{G}(\mathbf{e}; a) = \mathbf{G}(\mathbf{e}; \mathbf{a}_c) + (a - \mathbf{a}_c)\mathbf{e}$ by the definition of \mathbf{G} ,

$$\mathbf{G}(x_0(a); a) - \mathbf{G}(\mathbf{e}; a) = \frac{\beta \alpha}{n^{1/3}} (x_0(\mathbf{a}_c) - \mathbf{e}) + O(n^{-2/3}). \quad (263)$$

We first evaluate $\mathbf{\Gamma}_{n-j}(a)$ as in Lemma 5.1. Note that in Subsection 5.1.1, we used properties of $\mathbf{H}(z; a)$ for the integrals over Σ_{\pm} and properties of $\mathbf{G}(x; a)$ for the integral over (\mathbf{e}, ∞) . Since there is no change in the properties of $\mathbf{H}(z; a)$, the integrals over Σ_+ and Σ_- are computed exactly

the same as given by (219) and (220). For the integral over (\mathbf{e}, ∞) , note that the main contribution to (221) was from the part of y near \mathbf{e} since $\mathbf{G}(y; a)$ takes its maximum for y near \mathbf{e} . However, now due to (263) we need to add a contribution from y near $x_0(a)$. By using the standard Laplace's method as in (99), the contribution to the integral near $x_0(a)$ equals

$$\sqrt{\frac{2\pi}{-n\mathbf{G}''(x_0(\mathbf{a}_c); \mathbf{a}_c)}} \mathcal{M}_{j,n}(x_0(\mathbf{a}_c)) e^{n(\mathbf{G}(x_0(a); a) - \ell/2)} (1 + o(1)), \quad (264)$$

Adding (264) and (221),

$$\int_{\mathbf{e}}^{\infty} \varphi_{n-j}(y) e^{nay} dy = \frac{Q_n}{\beta\sqrt{n}} \left(\mathcal{B}_{j,n}(\mathbf{e}) \int_0^{\infty} \text{Ai}(\xi) e^{\alpha\xi} d\xi + \beta \sqrt{\frac{2\pi}{-\mathbf{G}''(x_0(\mathbf{a}_c); \mathbf{a}_c)}} \mathcal{M}_{j,n}(x_0(\mathbf{a}_c)) e^{n(\mathbf{G}(x_0(a); a) - \mathbf{G}(\mathbf{e}; a))} + o(1) \right). \quad (265)$$

But (263) implies that

$$n(\mathbf{G}(x_0(a); a) - \mathbf{G}(\mathbf{e}; a)) = n^{2/3} \beta \alpha (x_0(\mathbf{a}_c) - \mathbf{e}) + O(n^{1/3}) \ll 0 \quad (266)$$

for all large enough n since $\alpha < 0$. Hence we find that

$$\int_{\mathbf{e}}^{\infty} \varphi_{n-j}(y) e^{nay} dy = \frac{Q_n}{\beta\sqrt{n}} \left(\mathcal{B}_{j,n}(\mathbf{e}) \int_0^{\infty} \text{Ai}(\xi) e^{\alpha\xi} d\xi + o(1) \right), \quad (267)$$

as in Lemma 5.1. Adding the integrals on the contours Σ_{\pm} , we obtain

$$\Gamma_{n-j}(a) = \frac{Q_n}{\beta\sqrt{n}} e^{\alpha^3/3} (\mathcal{B}_{j,n}(\mathbf{e}) + o(1)), \quad (268)$$

which is same as (197).

The evaluation of $\tilde{\psi}_{n-j}(x)$ is similar. We use the formula (226) as in Lemma 5.2. The evaluation of the integral over $(\Sigma_+ + C_n) \cup (\Sigma_- + C_n)$ is exactly the same as in Subsubsection 5.1.2. For the evaluation of the integral over $(\mathbf{e} + C_n, \infty)$ when $x \in E_{T, \epsilon/2}$, we find that (236) is unchanged. For the integral over $(\mathbf{e} + \epsilon, \infty)$, the contribution near $y = x_0(a)$ implies that (238) becomes

$$\int_{\mathbf{e}+\epsilon}^{\infty} K_{n-j,n}(x, y) e^{n(ay - \frac{1}{2}V(y))} dy = Q_n O(n^{1/6} e^{-\frac{3}{5}|\xi|^{3/2}} \frac{e^{n(\mathbf{G}(x_0(a); a) - \mathbf{G}(\mathbf{e}; a))}}{\sqrt{n}}). \quad (269)$$

However, due to (266), this is again $Q_n O(n^{-1/3} e^{-\frac{3}{5}|\xi|^{3/2}})$ as in (238). Therefore the result (224) still holds for $x \in E_{T, \epsilon/2}$. For $x \in (\mathbf{e} + \epsilon/2, \infty)$, the estimates (254) and (252) hold without any change. Moreover, it is straightforward to check that (253) still holds. Therefore (225) holds for $x \geq \mathbf{e} + \epsilon/2$. Therefore, Lemma 5.2 holds without any changes.

Now we proceed as in Subsection 5.1.3. In evaluating $\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{I_n^T}$, the integral over $(\mathbf{e} + \epsilon/2)$ becomes, due to the contribution near $x_0(a)$,

$$\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{(\mathbf{e}+\epsilon, \infty)} = \int_{\mathbf{e}+\epsilon}^{\infty} O(\sqrt{n} e^{n(\mathbf{G}(x; a) - \mathbf{G}(\mathbf{e}; a))}) dx = O(e^{n(\mathbf{G}(x_0(a); a) - \mathbf{G}(\mathbf{e}; a))}) \rightarrow 0 \quad (270)$$

by (266). This implies that (256) holds without a change. Similarly, it is straightforward to check that (259) holds. Therefore, we obtain (260) and Theorem 1.5(a) is proved.

5.2.2 Proof of Theorem 1.5(b)

We consider the double scaling situation

$$a = \mathbf{a}_c + \frac{\alpha'}{n} \quad (271)$$

where α' is in a compact subset of \mathbb{R} . In this case, (262) and (263) are changed to

$$\mathbf{G}(x_0(a); a) = \mathbf{G}(x_0(\mathbf{a}_c); \mathbf{a}_c) + x_0(\mathbf{a}_c) \frac{\alpha'}{n} + O(n^{-2}), \quad (272)$$

and

$$\mathbf{G}(x_0(a); a) - \mathbf{G}(\mathbf{e}; a) = \frac{\alpha'}{n} (x_0(\mathbf{a}_c) - \mathbf{e}) + O(n^{-2}). \quad (273)$$

First we consider $\mathbf{\Gamma}_{n-j}(a)$. There are two changes from the previous subsection. The first is that since α defined in (261) and α' defined in (271) are related as $\alpha = \beta^{-1} n^{-2/3} \alpha'$, we have $\alpha \rightarrow 0$ and hence in (219), (220) and (265), we have $\alpha = 0$ in the integrals involving the Airy function. The second is that (267) does not follow from (265) since (266) no longer holds. Instead, due to (273), (265) implies that

$$\begin{aligned} \int_{\mathbf{e}}^{\infty} \varphi_{n-j}(y) e^{n a y} dy &= \frac{Q_n}{\beta \sqrt{n}} \left(\mathcal{B}_{j,n}(\mathbf{e}) \int_0^{\infty} \text{Ai}(\xi) d\xi \right. \\ &\quad \left. + \beta \sqrt{\frac{2\pi}{-\mathbf{G}''(x_0(\mathbf{a}_c); \mathbf{a}_c)}} \mathcal{M}_{j,n}(x_0(\mathbf{a}_c)) e^{\alpha' (x_0(\mathbf{a}_c) - \mathbf{e})} + o(1) \right). \end{aligned} \quad (274)$$

Hence adding (219) and (220) (with $\alpha = 0$), we obtain

$$\mathbf{\Gamma}_{n-j}(a) = \frac{Q_n}{\beta \sqrt{n}} \mathcal{B}_{j,n}(\mathbf{e}) \left(\frac{D_0 + D_1(\alpha')}{D_0} + o(1) \right), \quad (275)$$

where

$$D_0(\alpha') = \frac{\mathcal{B}_{j,n}(\mathbf{e})}{\beta}, \quad D_1(\alpha') = \sqrt{\frac{2\pi}{-\mathbf{G}''(x_0(\mathbf{a}_c); \mathbf{a}_c)}} \mathcal{M}_{j,n}(x_0(\mathbf{a}_c)) e^{\alpha' (x_0(\mathbf{a}_c) - \mathbf{e})}. \quad (276)$$

Now consider $\tilde{\psi}_{n-j}(x)$. Like in Subsubsection 5.2.1, most of the estimates of Subsubsection 5.1.2 remain the same. The only changes are the contribution near $x_0(\mathbf{a}_c)$. Substituting (273) into (269) with $a = \beta^{-1} n^{-2/3} \alpha'$, we have

$$\int_{\mathbf{e} + \epsilon}^{\infty} K_{n-j,n}(x, y) e^{n(a y - \frac{1}{2} V(y))} dy = Q_n O(n^{-1/3} e^{-\frac{3}{5} |\xi|^{3/2}}) \quad (277)$$

for $x \in E_{T, \epsilon/2}$. This is similar to (238) except that $e^{-\frac{\epsilon'}{2} n}$ is replaced by $n^{-1/3}$. However, it is easy to check that this bound is enough for the rest of the analysis. The rest of the analysis of Subsubsection 5.1.2 continues without changes and we obtain the asymptotics (249) of $\mathbf{\Gamma}_{n-j}(a) \tilde{\psi}_{n-j}(x)$ for $x \in E_{T, \epsilon/2}$ and (255) for $x \geq \mathbf{e} + \epsilon/2$. Using (275), we find that

$$\tilde{\psi}_{n-j}(x) = p_{j,n}(\alpha') \frac{\beta \sqrt{n}}{\mathcal{B}_{j,n}(\mathbf{e})} \left[C_0(\xi) + e^{-\alpha^3/3} \left(\frac{1}{Q_n} e^{a n - \frac{1}{2} V(x)} - e^{\alpha \xi} \right) + o(1) \right] \quad (278)$$

for $x \in E_{T,\epsilon/2}$ and

$$\tilde{\psi}_{n-j}(x) = p_{j,n}(\alpha') \frac{\beta\sqrt{n}}{\mathcal{B}_{j,n}(\mathbf{e})} e^{n(\mathbf{G}(x;a) + \mathbf{H}(x;a) - 2\mathbf{G}(\mathbf{e};a))/2} (1 + o(1)) \quad (279)$$

for $x \geq \mathbf{e} + \epsilon/2$, where

$$p_{j,n}(\alpha') := \frac{D_0}{D_0 + D_1(\alpha')}. \quad (280)$$

The formulas (278) and (279) are different from (224) and (225) only by the factor $p_{j,n}(\alpha')$.

We now prove the theorem. First, consider (58). From (279) and (330), we obtain

$$\begin{aligned} \langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{J_n^T(x_0(\mathbf{a}_c))} &= p_{j,n}(\alpha') \frac{\beta\sqrt{n}}{\mathcal{B}_{j,n}(\mathbf{e})} \int_{J_n^T(x_0(\mathbf{a}_c))} M_{j,n}(x) e^{n(\mathbf{G}(x;a) - \mathbf{G}(\mathbf{e};a))} (1 + o(1)) dx \\ &= p_{j,n}(\alpha') \frac{\beta\sqrt{n}}{\mathcal{B}_{j,n}(\mathbf{e})} \mathcal{M}_{j,n}(x_0(\mathbf{a}_c)) \frac{e^{n(\mathbf{G}(x_0(a);a) - \mathbf{G}(\mathbf{e};a))}}{\sqrt{-n\mathbf{G}''(x_0(a);a)/2\pi}} \int_T^\infty e^{-\frac{1}{2}\xi^2} d\xi (1 + o(1)) \\ &= (1 - p_{j,n}(\alpha'))(1 - G(T))(1 + o(1)), \end{aligned} \quad (281)$$

by using (272) and (273). Also, for $x \in J_n^T(x_0(\mathbf{a}_c))$, by (336) and (279)

$$\begin{aligned} (K_{n-j,n} \chi_{J_n^T(x_0(\mathbf{a}_c))} \tilde{\psi}_{n-j})(x) &= \int_{J_n^T(x_0(\mathbf{a}_c))} K_{n-j,n}(x, y) \tilde{\psi}_{n-j}(y) dy \\ &= O(e^{n(\mathbf{G}(x;a) - \mathbf{H}(x;a))/2}) \int_{J_n^T(x_0(\mathbf{a}_c))} O(\sqrt{n} e^{n(\mathbf{G}(y;a) - \mathbf{G}(\mathbf{e};a))}) dy \\ &= O(e^{n(\mathbf{G}(x;a) - \mathbf{H}(x;a))/2}). \end{aligned} \quad (282)$$

Hence $K_{n-j,n} \chi_{J_n^T(x_0(\mathbf{a}_c))} \tilde{\psi}_{n-j} \rightarrow 0$ in $L^2(J_n^T(x_0(\mathbf{a}_c)))$. Therefore, we find from (76) that

$$\mathbb{P}_{n-j+1,n}(a; J_n^T(x_0(\mathbf{a}_c))) = 1 - (1 - p_{j,n}(\alpha'))(1 - G(T))(1 + o(1)) \quad (283)$$

and (58) is proved.

Second, consider (57). We proceed as in the outline (b2) in Section 2. We first evaluate $\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{I_n^T} = \langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{E_{T,\epsilon/2}} + \langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{(\mathbf{e} + \epsilon/2, \infty)}$. For second term, a computation as in (281) yields that

$$\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{(\mathbf{e} + \epsilon/2, \infty)} = (1 - p_{j,n}(\alpha'))(1 + o(1)). \quad (284)$$

The first term is calculated as in Subsubsection 5.2.1 with the only change that the prefactor $p_{j,n}(\alpha')$ is multiplied:

$$\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{E_{T,\epsilon/2}} = p_{j,n}(\alpha') \int_T^\infty C_0(\xi) \text{Ai}(\xi) d\xi (1 + o(1)). \quad (285)$$

Hence

$$\langle \tilde{\psi}_{n-j}, \psi_{n-j} \rangle_{I_n^T} = p_{j,n}(\alpha') \int_T^\infty C_0(\xi) \text{Ai}(\xi) d\xi + (1 - p_{j,n}(\alpha')) + o(1). \quad (286)$$

Now we evaluate $u_{j,n}(\xi) := n^{-1/2}K_{n-j,n}\chi_{I_n^T}\tilde{\psi}_{n-j}(x)$, $x = \mathbf{e} + \beta^{-1}n^{-2/3}\xi \in I_n^T$. But it is straightforward to check that, as in Subsubsection 5.2.1,

$$u_{j,n} - p_{j,n}(\alpha') \frac{\beta}{\mathcal{B}_{j,n}(\mathbf{e})} (K_{\text{Airy}}\chi_{[T,\infty)}C_0) \rightarrow 0 \quad \text{in } L^2([T, \infty)). \quad (287)$$

Therefore, from (79),

$$\begin{aligned} \mathbb{P}_{n-j+1,n}(a; I_n^T) &= F_0(T) \cdot (1 - (p_{j,n}(\alpha')\langle C_0, \text{Ai} \rangle_{[T,\infty)} + 1 - p_{j,n}(\alpha')) \\ &\quad - p_{j,n}(\alpha')\langle (1 - \chi_{[T,\infty)}K_{\text{Airy}}\chi_{[T,\infty)})^{-1}(K_{\text{Airy}}\chi_{[T,\infty)}C_0), \text{Ai} \rangle_{[T,\infty)}) + o(1) \\ &= p_{j,n}(\alpha')F_1(T; 0) + o(1). \end{aligned} \quad (288)$$

Thus (57) is proved.

6 Summary of asymptotics of orthogonal polynomials and the Christoffel-Darboux kernel

Define the matrix-valued function

$$Y_k(z; n) = \begin{pmatrix} \gamma_k^{-1}(n)p_k(z; n) & \gamma_k^{-1}(n)(C\varphi_k)(z; n) \\ -2\pi i\gamma_{k-1}(n)p_{k-1}(z; n) & -2\pi i\gamma_{k-1}(n)(C\varphi_{k-1})(z; n) \end{pmatrix}, \quad (289)$$

for $z \in \mathbb{C} \setminus \mathbb{R}$. Here $(C\varphi_k)(z; n)$ is the Cauchy transform of $\varphi_k(x; n) = p_k(x; n)e^{-nV(x)}$. The matrix $Y_k(z; n)$ is the solution to the Riemann-Hilbert problem for orthogonal polynomials with weight $e^{-nV(z)}$ (see [20]). We are interested in the asymptotics of $Y_{n-j}(z; n)$ when $j = O(1)$ and $n \rightarrow \infty$. We indicate the changes from the analysis of [17] and state the results. A similar derivation for the discrete weight can be found in [4].

Fix $\delta > 0$ small enough. Let (see Figure 10)

$$A_\delta := \{z \in \mathbb{C} \mid |z - \mathbf{e}| < \delta\}, \quad (290)$$

$$B_\delta := \{z \in \mathbb{C} \mid \Re z \geq \mathbf{e} \text{ and } |z - \mathbf{e}| > \delta\} \cup \{z \in \mathbb{C} \mid \Re z < \mathbf{e} \text{ and } |\Im z| > \delta\}, \quad (291)$$

where \mathbf{e} is the rightmost end-point of the support of the equilibrium measure. Comparing with notations in [17], A_δ is the circle $D_{\epsilon, a_{N+1}}$ with δ corresponding to the radius ϵ . in [17, Figure 1.4]. As in [17, Figure 1.4], A_δ is divided into four regions I, II, III and IV. Let Σ_R be the contour in [17, Figure 4.9]. We assume that the boundary of A_δ is a part of Σ_R and B_δ is outside of the lens-shaped regions, cf. [17, Formula (4.116)].

Several notations from [17] are used in this section, and we summarize them in Table 1. Other notations may be slightly different but should be clear.

By following the procedure of [17], we find asymptotics of Y . Noting the symmetry

$$Y(z) = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \overline{Y(\bar{z})} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad (292)$$

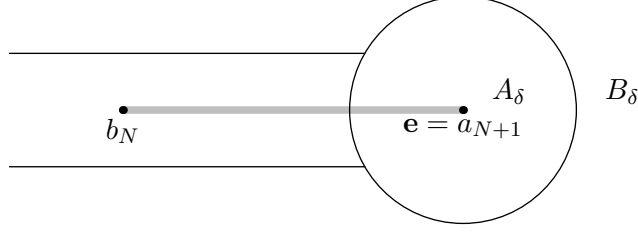


Figure 10: A_δ and B_δ .

| Notation | Definition in [17] |
|-------------------------|--|
| $\vec{\Omega}$ | defined in Formula (1.21) |
| θ | defined in Formula (1.24) |
| γ | defined in Formula (1.26) |
| u | defined in Formula (1.29) |
| d | defined in Formula (1.30) |
| $u_+(\infty)$ | explained below Theorem 1.1 |
| $\Phi = \Phi_{a_{N+1}}$ | defined in Formula (1.34) |
| P | defined in Formulas (1.38)–(1.40) |
| σ_3 | mentioned between Formulas (1.40) and (1.41) |
| $\Sigma^{(1)}$ | shown in Figure 1.5 |
| $v^{(\infty)}$ | defined in Formulas (1.104)–(1.107) |
| G | defined in Formula (3.44) |
| Σ_R | shown in Figure 4.9 |
| v_R | defined in Formula (4.108) |

Table 1: Notations taken from [17]

it is enough to consider $z \in \mathbb{C}_+$. The asymptotic formulas of $Y(z)$ are different in B_δ and A_δ . For $z \in B_\delta \cap \mathbb{C}_+$, we have (cf. [17, Formula 4.116])

$$Y_{n-j}(z; n) = e^{\frac{n\ell}{2}\sigma_3} R_{j,n}(z) M_{j,n}^{(\infty)}(z) e^{n(\mathbf{g}(z) - \frac{\ell}{2})\sigma_3}. \quad (293)$$

The remainder $R_{j,n}(z)$ solves the Riemann-Hilbert problem similar to the Riemann-Hilbert problem for R in [17, Subsection 4.6] (cf. [17, Formulas (4.106)–(4.108)]) with the jump matrix which has the same estimate as in v_R shown in [17, Figure 1.4]. This implies that uniformly for all $z \in \mathbb{C} \setminus \Sigma_R$ and n (cf. [17, Formula 4.115])

$$R_{j,n}(z) = I + O(n^{-1}). \quad (294)$$

The outer parametrix $M_{j,n}^{(\infty)}(z)$ solves the Riemann-Hilbert problem (cf. [17, Formulas (4.24)–

(4.26)]

$$M_{j,n}^{(\infty)}(z) \text{ is analytic in } \mathbb{C} \setminus \Sigma^{(1)} \quad (295)$$

$$(M_{j,n}^{(\infty)})_+(z) = (M_{j,n}^{(\infty)})_-(z)v^{(\infty)}(z), \quad z \in \Sigma^{(1)}, \quad (296)$$

$$M_{j,n}^{(\infty)}(z) = \left(I + O\left(\frac{1}{z}\right) \right) \begin{pmatrix} z^{-j} & 0 \\ 0 & z^j \end{pmatrix}, \quad z \rightarrow \infty, \quad z \in \mathbb{C} \setminus \mathbb{R}. \quad (297)$$

Note that the dependence on j in the asymptotics as $z \rightarrow \infty$. The solution of this Riemann-Hilbert problem can be solved as in [17, Lemma 4.3]). Setting

$$\Theta(\infty) := \frac{\theta(u_+(\infty) + d)}{\theta(u_+(\infty) - d)}, \quad (298)$$

$$\tilde{\Theta}(z) := \frac{\gamma(z) - \gamma(z)^{-1} \theta(u(z) + d)}{\gamma(z) + \gamma(z)^{-1} \theta(u(z) - d)}, \quad (299)$$

it is straightforward to check that for $z \in B_\delta \cap \mathbb{C}_+$,

$$\begin{aligned} M_{j,n}^{(\infty)}(z) &= \text{diag} \left(\left(\frac{\sum_{l=0}^N a_{l+1} - b_l}{4} \right)^{-j}, \left(\frac{\sum_{l=0}^N a_{l+1} - b_l}{4} \right)^j \right) \\ &\times \text{diag} \left(\frac{\Theta(\infty)^{-j} \theta(u_+(\infty) + d)}{\theta(u_+(\infty) - \frac{n}{2\pi} \vec{\Omega} - (2j-1)d)}, \frac{\Theta(\infty)^j \theta(u_+(\infty) + d)}{\theta(u_+(\infty) + \frac{n}{2\pi} \vec{\Omega} + (2j+1)d)} \right) \\ &\times \begin{pmatrix} \tilde{\Theta}(z)^j \frac{\gamma + \gamma^{-1}}{2} \frac{\theta(u(z) - \frac{n}{2\pi} \vec{\Omega} - (2j-1)d)}{\theta(u(z)+d)} & \tilde{\Theta}(z)^{-j} \frac{\gamma - \gamma^{-1}}{-2i} \frac{\theta(u(z) + \frac{n}{2\pi} \vec{\Omega} + (2j-1)d)}{\theta(u(z)-d)} \\ \tilde{\Theta}(z)^j \frac{\gamma - \gamma^{-1}}{2i} \frac{\theta(u(z) - \frac{n}{2\pi} \vec{\Omega} - (2j+1)d)}{\theta(u(z)-d)} & \tilde{\Theta}(z)^{-j} \frac{\gamma + \gamma^{-1}}{2} \frac{\theta(u(z) + \frac{n}{2\pi} \vec{\Omega} + (2j+1)d)}{\theta(u(z)+d)} \end{pmatrix}, \end{aligned} \quad (300)$$

and for $z \in B_\delta \cap \mathbb{C}_-$,

$$\begin{aligned} M_{j,n}^{(\infty)}(z) &= \left(\left(\frac{\sum_{l=0}^N a_{l+1} - b_l}{4} \right)^{-j}, \left(\frac{\sum_{l=0}^N a_{l+1} - b_l}{4} \right)^j \right) \\ &\times \text{diag} \left(\frac{\Theta(\infty)^{-j} \theta(u_+(\infty) + d)}{\theta(u_+(\infty) - \frac{n}{2\pi} \vec{\Omega} - (2j-1)d)}, \frac{\Theta(\infty)^j \theta(u_+(\infty) + d)}{\theta(u_+(\infty) + \frac{n}{2\pi} \vec{\Omega} + (2j+1)d)} \right) \\ &\times \begin{pmatrix} \tilde{\Theta}(z)^{-j} \frac{\gamma - \gamma^{-1}}{-2i} \frac{\theta(u(z) + \frac{n}{2\pi} \vec{\Omega} + (2j-1)d)}{\theta(u(z)+d)} & \tilde{\Theta}(z)^j \frac{\gamma + \gamma^{-1}}{-2} \frac{\theta(u(z) - \frac{n}{2\pi} \vec{\Omega} - (2j-1)d)}{\theta(u(z)+d)} \\ \tilde{\Theta}(z)^{-j} \frac{\gamma + \gamma^{-1}}{2} \frac{\theta(u(z) + \frac{n}{2\pi} \vec{\Omega} + (2j+1)d)}{\theta(u(z)+d)} & \tilde{\Theta}(z)^j \frac{\gamma - \gamma^{-1}}{-2i} \frac{\theta(u(z) - \frac{n}{2\pi} \vec{\Omega} - (2j+1)d)}{\theta(u(z)-d)} \end{pmatrix}. \end{aligned} \quad (301)$$

Note that both formulas (300) and (301) can be extended to $z \in B_\delta \cap \mathbb{R}$.

The asymptotics (293) especially implies the asymptotics of γ_{n-j} . Since (cf. [16, Formulas (3.10) and (3.11)])

$$\gamma_{n-j} = \left(-2\pi i \lim_{|z| \rightarrow \infty} z^{-n+j+1} (R_{j,n}(z) M_{j,n}^{(\infty)}(z))_{12} e^{-n(\mathbf{g}(z)+\ell)} \right)^{-1/2}, \quad (302)$$

we have from (294) and (300) that

$$\gamma_{n-j} = \hat{\gamma}_{n-j} (1 + O(n^{-1})), \quad (303)$$

where

$$\hat{\gamma}_{n-j}^2 = \left(\frac{\pi \left(\sum_{l=0}^N a_{l+1} - b_l \right) \theta(u_+(\infty) + d) \theta(u_+(\infty) + \frac{n}{2\pi} \vec{\Omega} + (2j-1)d)}{2 \theta(u_+(\infty) - \frac{n}{2\pi} \vec{\Omega} - (2j-1)d) \theta(u_+(\infty) - d)} \right)^{-1} \times \left(\frac{\sum_{l=0}^N a_{l+1} - b_l}{4} \Theta(\infty) \right)^{2j} e^{-n\ell}. \quad (304)$$

See [17, Formulas (1.62) and (1.63)] for the cases $j = 1$ and $j = 0$.

Now consider $z \in A_\delta$. Then the analysis of the local parametrix as in [17, Section 4.3] implies that (cf. [17, (4.119)–(4.121)])

$$Y_{n-j}(z; n) = e^{\frac{n\ell}{2}\sigma_3} R_{j,n}(z) (M_{j,n})_p(z) e^{n(\mathbf{g}(z) - \frac{\ell}{2})\sigma_3} \quad (305)$$

for z is in regions I and IV in [17, Figure 1.4],

$$Y_{n-j}(z; n) = e^{\frac{n\ell}{2}\sigma_3} R_{j,n}(z) (M_{j,n})_p(z) \begin{pmatrix} 1 & 0 \\ e^{-nG(z)} & 1 \end{pmatrix} e^{n(\mathbf{g}(z) - \frac{\ell}{2})\sigma_3} \quad (306)$$

for z is in regions II in [17, Figure 1.4], and

$$Y_{n-j}(z; n) = e^{\frac{n\ell}{2}\sigma_3} R_{j,n}(z) (M_{j,n})_p(z) \begin{pmatrix} 1 & 0 \\ -e^{nG(z)} & 1 \end{pmatrix} e^{n(\mathbf{g}(z) - \frac{\ell}{2})\sigma_3} \quad (307)$$

for z is in regions III in [17, Figure 1.4]. Here the local parametrix $(M_{j,n})_p$ is given by (cf. [17, Formulas (4.75) and (4.76)])

$$(M_{j,n})_p(z) := M_{j,n}^{(\infty)}(z) \frac{1}{\sqrt{2i}} \begin{pmatrix} i & -i \\ 1 & 1 \end{pmatrix} (\Phi(z))^{\sigma_3/4} P(\Phi(z)). \quad (308)$$

where $\Phi(z)$ denotes $\Phi_{a_{N+1}}(z)$ in [17]. We note that by definition

$$\Phi(z) = \left[-\frac{3n}{4} (2\mathbf{g}(z) - V(z) - \ell) \right]^{2/3}, \quad (309)$$

so that

$$\Phi(z) = \beta n^{2/3} (z - \mathbf{e}) (1 + O(|z - \mathbf{e}|)) \quad \text{as } z \rightarrow \mathbf{e} \quad (310)$$

with β defined in (22) (see [17, Equations (1.34), (1.35), (4.74)]).

We now summarize the asymptotics the orthonormal polynomials and their Cauchy transformations. For notational convenience, we denote for $z \in B_\delta$

$$\mathcal{M}_{j,n}(z) := \hat{\gamma}_{n-j} e^{\frac{n\ell}{2}} (M_{j,n}^{(\infty)})_{11}(z), \quad (311)$$

$$\tilde{\mathcal{M}}_{j,n}(z) := \hat{\gamma}_{n-j} e^{\frac{n\ell}{2}} (M_{j,n}^{(\infty)})_{12}(z), \quad (312)$$

and for $z \in A_\delta$

$$\mathcal{B}_{j,n}(z) := \sqrt{\pi} n^{-1/6} \hat{\gamma}_{n-j} e^{\frac{n\ell}{2}} \left((M_{j,n}^{(\infty)})_{11}(z) - i (M_{j,n}^{(\infty)})_{12}(z) \right) \Phi(z)^{1/4}, \quad (313)$$

$$\mathcal{D}_{j,n}(z) := \sqrt{\pi} n^{1/6} \hat{\gamma}_{n-j} e^{\frac{n\ell}{2}} \left(-(M_{j,n}^{(\infty)})_{11}(z) - i (M_{j,n}^{(\infty)})_{12}(z) \right) \Phi(z)^{-1/4}. \quad (314)$$

Remark 6.1. The formulas (300) and (301) contain the N -variable Theta function θ . If $N = 0$, i.e., the equilibrium measure is supported on one interval, then $\theta \equiv 1$, and the expressions of $\hat{\gamma}_{n-j}$, $\mathcal{M}_{j,n}(z)$, $\tilde{\mathcal{M}}_{j,n}(z)$, $\mathcal{B}_{j,n}(z)$ and $\mathcal{D}_{j,n}(z)$ are much simplified. In particular $\mathcal{M}_{j,n}(z)$, $\tilde{\mathcal{M}}_{j,n}(z)$, $\mathcal{B}_{j,n}(z)$ and $\mathcal{D}_{j,n}(z)$ do not depend on n when $N = 0$. For example, for $z \in \mathbb{C}_+$

$$\mathcal{M}_{j,n}(z) = \sqrt{\frac{2}{\pi(a_1 - b_0)}} \frac{\gamma + \gamma^{-1}}{2} \left(\frac{\gamma - \gamma^{-1}}{\gamma + \gamma^{-1}} \right)^j, \quad (315)$$

$$\tilde{\mathcal{M}}_{j,n}(z) = \sqrt{\frac{2}{\pi(a_1 - b_0)}} \frac{\gamma - \gamma^{-1}}{-2i} \left(\frac{\gamma - \gamma^{-1}}{\gamma + \gamma^{-1}} \right)^{-j}, \quad (316)$$

and

$$\mathcal{B}_{j,n}(\mathbf{e}) = \sqrt{2}(a_1 - b_0)^{-1/4} \beta^{1/4}. \quad (317)$$

Proposition 6.1. *There exists $\delta_0 > 0$ such that for each fixed $\delta \in (0, \delta_0]$, the following holds as $n \rightarrow \infty$ and $j = O(1)$.*

(a) For $z \in B_\delta$,

$$\varphi_{n-j}(z; n) = M_{j,n}(z) e^{n(\mathbf{g}(z) - V(z) - \ell/2)}, \quad (318)$$

where $M_{j,n}(z)$ is an analytic function in B_δ and

$$M_{j,n}(z) = \mathcal{M}_{j,n}(z)(1 + O(n^{-1})) \quad (319)$$

uniformly in z and n . The function $M_{j,n}$ satisfies that (i) $M_{j,n}(z) = O(z^{-j})$ uniformly in n as $z \rightarrow \infty$, (ii) in any compact subset $K \in B_\delta$, $M_{j,n}(z)$, $M'_{j,n}(z)$ and $1/M_{j,n}(z)$ are $O(1)$ uniformly in n and $z \in K$, and (iii) $\mathcal{M}_{j,n}(x) > 0$ and $M_{j,n}(x) > 0$ for all real $x > \mathbf{e}$.

(b) For $z \in B_\delta$,

$$(C\varphi_{n-j})(z; n) = \tilde{M}_{j,n}(z) e^{n(-\mathbf{g}(z) + \ell/2)}, \quad (320)$$

where $\tilde{M}_{j,n}(z)$ is analytic in $B_\delta \setminus \mathbb{R}$, continuous up to the boundary, and

$$\tilde{M}_{j,n}(z) = \tilde{\mathcal{M}}_{j,n}(z)(1 + O(n^{-1})) \quad (321)$$

uniformly in z and n where $\tilde{\mathcal{M}}_{j,n}(z)$ defined in (312) is analytic in B_δ . The function $\tilde{M}_{j,n}$ satisfies that (i) $\tilde{M}_{j,n}(z) = O(z^j)$ uniformly in n as $z \rightarrow \infty$, (ii) in any compact subset $K \in B_\delta$, $\tilde{M}_{j,n}(z)$, $\tilde{M}'_{j,n}(z)$ and $1/\tilde{M}_{j,n}(z)$ are $O(1)$ uniformly in n and $z \in K \setminus \mathbb{R}$ and (iii) $-i\tilde{\mathcal{M}}_{j,n}(x) > 0$ for $x > \mathbf{e}$.

(c) For $z \in A_\delta$,

$$\varphi_{n-j}(z) = \left(n^{1/6} \text{Ai}(\Phi(z)) B_{j,n}(z) + n^{-1/6} \text{Ai}'(\Phi(z)) D_{j,n}(z) \right) e^{-\frac{n}{2}V(z)}, \quad (322)$$

where $B_{j,n}(z)$ and $D_{j,n}(z)$ are analytic functions in A_δ and

$$B_{j,n}(z) = \mathcal{B}_{j,n}(z)(1 + O(n^{-1})), \quad D_{j,n}(z) = \mathcal{D}_{j,n}(z)(1 + O(n^{-1})) \quad (323)$$

uniformly in z and n . The functions $B_{j,n}$ and $D_{j,n}$ satisfy (i) $B_{j,n}(z)$, $D_{j,n}(z)$, $B'_{j,n}(z)$, $D'_{j,n}(z)$, $1/B_{j,n}(z)$ and $1/D_{j,n}(z)$ are $O(1)$ uniformly in n and $z \in A_\delta$ and (ii) $\mathcal{B}_{j,n}(x) > 0$ and $\mathcal{D}_{j,n}(x) > 0$ for $x \in (\mathbf{e} - \delta, \mathbf{e} + \delta)$.

(d) For $z \in A_\delta \cap \mathbb{C}_+$, we have

$$(C\varphi_{n-j})(z) = e^{\pi i/3} \left(n^{1/6} \text{Ai}(\omega^2 \Phi(z)) B_{j,n}(z) + n^{-1/6} \omega^2 \text{Ai}'(\omega^2 \Phi(z)) D_{j,n}(z) \right) e^{-\frac{n}{2}V(z)}, \quad (324)$$

and for $z \in A_\delta \cap \mathbb{C}_-$, we have

$$(C\varphi_{n-j})(z) = -e^{\pi i/3} \left(n^{1/6} \omega^2 \text{Ai}(\omega \Phi(z)) B_{j,n}(z) + n^{-1/6} \text{Ai}'(\omega \Phi(z)) D_{j,n}(z) \right) e^{-\frac{n}{2}V(z)}, \quad (325)$$

where $\omega = e^{2\pi i/3}$ and $B_{j,n}(z)$ and $D_{j,n}(z)$ are the same functions in (322). The formulas (324) and (325) hold up to the boundary $z \in A_\delta \cap \mathbb{R}$.

Proof. By formulas (300) and (301) of $M_{j,n}^{(\infty)}(z)$, the properties of the theta function θ and the definition of d , we have that for $p, q = 1, 2$, the functions $(M_{j,n}^{(\infty)})_{pq}(z)$ and $(M_{j,n}^{(\infty)})_{pq}(z)^{-1}$ are uniformly bounded for z in any compact subset $K \subset B_\delta$ and the functions

$$\begin{aligned} & \left(i(M_{j,n}^{(\infty)})_{p1}(z) + (M_{j,n}^{(\infty)})_{p2}(z) \right) (z - \mathbf{e})^{1/4}, \\ & \left(i(M_{j,n}^{(\infty)})_{p1}(z) + (M_{j,n}^{(\infty)})_{p2}(z) \right)^{-1} (z - \mathbf{e})^{-1/4}, \\ & \left(-i(M_{j,n}^{(\infty)})_{p1}(z) + (M_{j,n}^{(\infty)})_{p2}(z) \right) (z - \mathbf{e})^{-1/4}, \\ & \left(i(M_{j,n}^{(\infty)})_{p1}(z) + (M_{j,n}^{(\infty)})_{p2}(z) \right)^{-1} (z - \mathbf{e})^{1/4} \end{aligned}$$

are uniformly bounded for $z \in A_\delta$ and n . We note that although $M_{j,n}^{(\infty)}(z)$ is not well defined on J , the functions considered above are well defined on $A_\delta \cap J$. Plugging in these estimates and the estimate (294) of $R_{j,n}(z)$ into (293), (305), (306) and (307) we obtain the estimates of $M_{j,n}(z)$, $1/M_{j,n}(z)$, $\tilde{M}_{j,n}(z)$, $1/\tilde{M}_{j,n}(z)$, $B_{j,n}(z)$, $1/B_{j,n}(z)$, $D_{j,n}(z)$ and $1/D_{j,n}(z)$. With the help of the Cauchy's integral formula, we further obtain the estimates of $M'_{j,n}(z)$, $\tilde{M}'_{j,n}(z)$, $B'_{j,n}(z)$ and $D'_{j,n}(z)$. From the formula of $M_{j,n}^{(\infty)}(z)$ we also derive the positivity of $\mathcal{M}_{j,n}(x)$, $M_{j,n}(x)$, $-i\tilde{\mathcal{M}}_{j,n}(x)$, $\mathcal{B}_{j,n}(x)$ and $B_{j,n}(x)$. \square

Remark 6.2. We use the following identity in the analysis. It is straightforward to derive from the Riemann-Hilbert problem of $Y_k(z; n)$ that $\det Y_k(z; n) \equiv 1$. This implies that

$$\frac{\gamma_k}{\gamma_{k-1}} = \frac{-1}{2\pi i} (p_k(z) C\varphi_{k-1}(z) - p_{k-1}(z) C\varphi_k(z)). \quad (326)$$

Taking $k = n - j$ and using asymptotic formulas (322), (324) and (325) in (326), with the help of [1, 10.4.11 and 10.4.12]

$$\text{Ai}(z)\omega \text{Ai}'(\omega z) - \text{Ai}'(z) \text{Ai}(\omega z) = \frac{e^{-\pi i/6}}{2\pi}, \quad (327)$$

$$\text{Ai}(z)\omega^2 \text{Ai}'(\omega^2 z) - \text{Ai}'(z) \text{Ai}(\omega^2 z) = \frac{e^{\pi i/6}}{2\pi}, \quad (328)$$

we find

$$B_{j,n}(z) D_{j+1,n}(z) - B_{j+1,n}(z) D_{j,n}(z) = \frac{\gamma_{n-j}}{\gamma_{n-j-1}}. \quad (329)$$

Proposition 6.1 implies the following asymptotic properties of ψ_{n-j} . These are used in the main analysis extensively.

Corollary 6.1. *Fix $T \in \mathbb{R}$. There exists $\delta_0 > 0$ such that for each fixed $\epsilon \in (0, \delta_0]$, the following holds as $n \rightarrow \infty$ and $j = O(1)$.*

(a) For $x \geq \mathbf{e} + \epsilon$,

$$\psi_{n-j}(x) = M_{j,n}(x)e^{n(\mathbf{G}(x) - \mathbf{H}(x))/2} = O(e^{n(\mathbf{G}(x) - \mathbf{H}(x))/2}(1 + |x|)^{-j}), \quad (330)$$

for $M_{j,n}(x)$ in (318).

(b) Let $E_{T,\epsilon} := [\mathbf{e} + \beta^{-1}n^{-2/3}T, \mathbf{e} + \epsilon]$ be the interval defined in (116). For $x \in E_{T,\epsilon}$,

$$\psi_{n-j}(x) = O(n^{1/6}e^{-\frac{2}{3}|\xi|^{3/2}}), \quad \xi := \beta n^{2/3}(x - \mathbf{e}). \quad (331)$$

(c) Let $I_n^T := [\mathbf{e} + \beta^{-1}n^{-2/3}T, \infty)$ be the interval defined in (23). Then

$$\|\psi_{n-j}\|_{L^2(I_n^T)} = O(n^{-1/6}). \quad (332)$$

Also for every \bar{x} in $(\mathbf{e} + \epsilon, \infty)$, there is $\epsilon' > 0$ such that

$$\|\psi_{n-j}\|_{L^2([\bar{x}, \infty))} = O(e^{-\epsilon'n}). \quad (333)$$

(d) As $n \rightarrow \infty$, $v_{j,n}(\xi) := \psi_{n-j}(\mathbf{e} + \beta^{-1}n^{-2/3}\xi)$ satisfies

$$v_{j,n}(\xi) \rightarrow \mathcal{B}_{j,n}(\mathbf{e}) \text{Ai}(\xi), \quad \text{in } L^2([T, \infty)). \quad (334)$$

Proof. The result (a) follows from (318) and noting that $2\mathbf{g}(z) - V(z) - \ell = \mathbf{G}(z; a) - \mathbf{H}(z; a)$.

For (b), note that $T \leq \xi \leq \epsilon n^{2/3}$. Thus, $|\text{Ai}(\xi)| \leq Ce^{-\frac{2}{3}|\xi|^{3/2}}$ and $|\text{Ai}'(\xi)| \leq C(|\xi|^{1/4} + 1)e^{-\frac{2}{3}|\xi|^{3/2}} \leq C'n^{1/6}e^{-\frac{2}{3}|\xi|^{3/2}}$ for some constants $C, C' > 0$. From (322) and the behavior of $\Phi(z)$ in A_ϵ , we obtain the estimate with the factor $\frac{2}{3}$ changed to a smaller constant which can be made arbitrarily close to $\frac{2}{3}$ if we take ϵ smaller. To be definite, we fix this constant as $\frac{3}{5}$.

For (c), by the asymptotics (330) and (331) of ψ_{n-1} , we find

$$\begin{aligned} \|\psi_{n-j}\|_{L^2(I_n^T)}^2 &\leq 2 \left(\|\psi_{n-j}\|_{L^2(E_{T,\epsilon})}^2 + \|\psi_{n-j}\|_{L^2([\mathbf{e} + \epsilon, \infty))}^2 \right) \\ &= 2 \left[\int_T^{\beta n^{2/3}\epsilon} O\left(n^{1/3}e^{-2\frac{3}{5}|\xi|^{3/2}}\right) \frac{d\xi}{\beta n^{2/3}} + \int_{\mathbf{e} + \epsilon}^\infty O\left(e^{n(\mathbf{G}(x) - \mathbf{H}(x))}(1 + |x|)^{-2j}\right) dx \right] \\ &= O(n^{-1/3}). \end{aligned} \quad (335)$$

The estimate (333) is similar.

Item (d) follows from Proposition 6.1 (c). □

The above asymptotics for ψ_{n-j} yields the asymptotics for the Christoffel-Darboux kernel $K_{n-j,n}(x, y)$.

Corollary 6.2. *Let $T \in \mathbb{R}$ be fixed. There exists $\delta_0 > 0$ such that for each fixed $\epsilon \in (0, \delta_0]$, the followings hold as $n \rightarrow \infty$ and $j = O(1)$. Let $E_{T,\epsilon}$ be the interval defined in (116).*

(a) For $x, y \in (\mathbf{e} + \epsilon/2, \infty)$,

$$K_{n-j,n}(x, y) = O(e^{n(\mathbf{G}(x) - \mathbf{H}(x) + \mathbf{G}(y) - \mathbf{H}(y))/2} (1 + |x|)^{-j} (1 + |y|)^{-j}). \quad (336)$$

(b) For $x, y \in E_{T,\epsilon}$,

$$K_{n-j,n}(x, y) = O(n^{2/3} e^{-\frac{3}{5}(|\xi|^{3/2} + |\eta|^{3/2})}), \quad (337)$$

where $\xi := \beta n^{2/3}(x - \mathbf{e})$ and $\eta := \beta n^{2/3}(y - \mathbf{e})$.

(c) For $x \in (\mathbf{e} + \epsilon, \infty)$ and $y \in E_{T,\epsilon/2}$,

$$K_{n-j,n}(x, y) = O(n^{1/6} e^{n(\mathbf{G}(x) - \mathbf{H}(x))/2} e^{-\frac{3}{5}|\eta|^{3/2}} (1 + |x|)^{-j}). \quad (338)$$

(d) For $x \in E_{T,\epsilon/2}$ and $y \in (\mathbf{e} + \epsilon, \infty)$,

$$K_{n-j,n}(x, y) = O(n^{1/6} e^{-\frac{3}{5}|\xi|^{3/2}} e^{n(\mathbf{G}(y) - \mathbf{H}(y))/2} (1 + |y|)^{-j}). \quad (339)$$

All estimates above are uniform in x, y in their domains and in n .

Proof. Item (c) and (d) follow directly the asymptotics (330) and (331) of ψ_{n-j} , and the Christoffel-Darboux formula (69) of $K_{n-j,n}(x, y)$, noting that $x - y$ never vanishes.

For $x, y \in (\mathbf{e} + \epsilon/2, \infty)$, (330) implies that

$$K_{n-j,n}(x, y) = e^{n(\mathbf{G}(x) - \mathbf{H}(x) + \mathbf{G}(y) - \mathbf{H}(y))/2} \frac{M_{j,n}(x)M_{j+1,n}(y) - M_{j+1,n}(x)M_{j,n}(y)}{x - y}. \quad (340)$$

Since $M_{j,n}$ and its derivatives are uniformly bounded, we obtain (a).

Item (b) follows from a similar calculation but using the asymptotics (322). The calculation is direct and is the same as [14, Formula (3.8)]. \square

We also need the following results for the Christoffel-Darboux kernel.

Corollary 6.3. *Fix $T \in \mathbb{R}$ and let I_n^T be the interval defined in (23). Then we have the following:*

(a) For any fixed C , we have

$$\frac{1}{\beta n^{2/3}} K_{n-j,n}(x, y) = K_{\text{Airy}}(\xi, \eta) + o(e^{-C(|\xi| + |\eta|)}) \quad (341)$$

for all $x, y \in I_n^T$ as $n \rightarrow \infty$, where $\xi := (x - \mathbf{e})\beta n^{2/3}$ and $\eta := (y - \mathbf{e})\beta n^{2/3}$.

(b) Define the operator $\mathcal{K}_{n-j,n}$ by kernel

$$\mathcal{K}_{n-j,n}(\xi, \eta) := \frac{1}{\beta n^{2/3}} K_{n-j,n} \left(\mathbf{e} + \frac{\xi}{\beta n^{2/3}}, \mathbf{e} + \frac{\eta}{\beta n^{2/3}} \right). \quad (342)$$

Then

$$\left(1 - \chi_{[T,\infty)} \mathcal{K}_{n-j,n} \chi_{[T,\infty)} \right)^{-1} \rightarrow \left(1 - \chi_{[T,\infty)} K_{\text{Airy}} \chi_{[T,\infty)} \right)^{-1} \quad (343)$$

in trace norm as $n \rightarrow \infty$.

(c) The operator norms of $\left(1 - \chi_{I_n^T} K_{n-j,n} \chi_{I_n^T} \right)^{-1}$ are bounded uniformly in n . As a corollary, The operator norms of

$$\left(1 - \chi_{[\bar{x}_n, \infty)} K_{n-j,n} \chi_{[\bar{x}_n, \infty)} \right)^{-1}, \quad (344)$$

are also bounded uniformly in n and in \bar{x}_n as long as \bar{x}_n are in a compact subset of (\mathbf{e}, ∞) .

(d) We have

$$\lim_{n \rightarrow \infty} \det \left(1 - \chi_{I_n^T} K_{n-j,n} \chi_{I_n^T} \right) = F_0(T), \quad (345)$$

and

$$\lim_{n \rightarrow \infty} \det \left(1 - \chi_{[\bar{x}, \infty)} K_{n-j,n} \chi_{[\bar{x}, \infty)} \right) = 1. \quad (346)$$

for any \bar{x} is in a compact subset of (\mathbf{e}, ∞) .

Proof. The proof of a result similar to (a) for the non-varying weight is given in [14, Formula (3.8)]. The varying weight case is proved in the same way. Note that our C is the c in [14, Formula (3.8)], which can be assumed to be an arbitrarily large positive number.

The proof of a result similar to (b) for the non-varying weight is given in the proof of the $\beta = 2$ case in [14, Corollary 1.4]. The varying weight case is proved in the same way.

Items (c) and (d) follow from (b). □

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