

Perturbing Isoradial Triangulations

Jeanne Scott (+ François David)
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heuristic

discrete world

moduli-space of
Delaunay triangulations
 $\subset \text{Gr}_{\mathbb{C}}^{(2, \infty)} / \infty\text{-torus}$

Ψ
isoradial triangulations

correspondence
 \longleftrightarrow

\longleftrightarrow

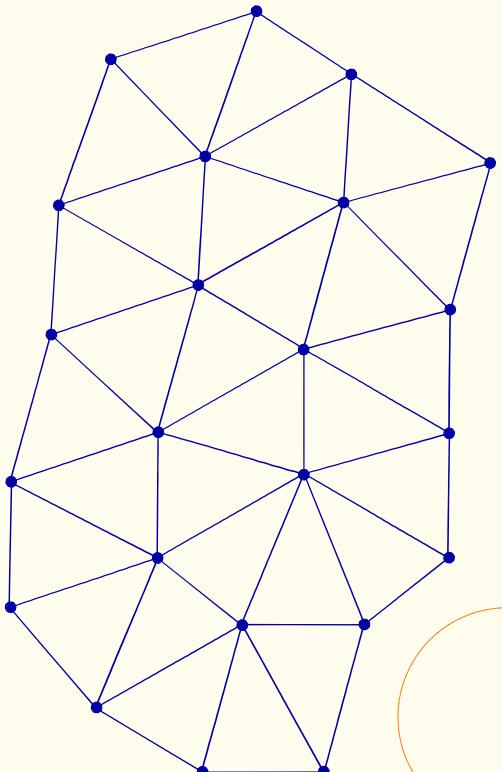
continuous world

space of metrics
 $g = (g_{\mu\nu})$ on the plane

Ψ

flat metrics

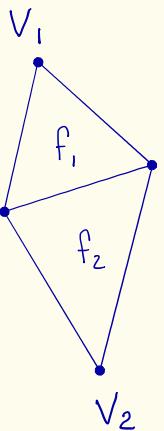
(fragment of an) infinite
Delaunay triangulation T



i.e infinite planar graph
embedded in \mathbb{C} with
vertices, edges, faces

$$V(T) \quad E(T) \quad F(T)$$

- each face $f \in F(T)$ triangle
- faces cover \mathbb{C} , locally finite
- for all abutting triangles f_1, f_2



$v_1 \notin$ interior of the disk
of circumcircle C_2

$v_2 \notin$ interior of the disk
of circumcircle C_1

Notation / set-up :

for $u \in V(T)$, $z(u) \in \mathbb{C}$ coordinate

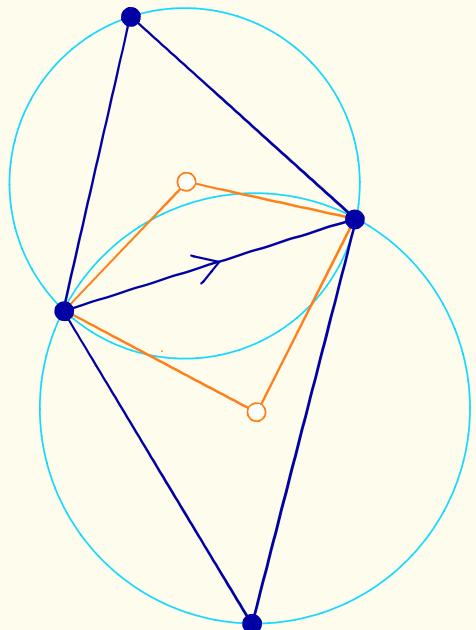
for an oriented edge \vec{uv} have

- north and south triangles X_n, X_s
- circumcenters x_n, x_s
- circumradii R_n, R_s
- angles θ_n, θ_s

def the conformal angle

$$\theta(uv) = \frac{1}{2}(\theta_n + \theta_s)$$

(symmetric, $SL_2(\mathbb{C})$ - invariant)



three operators Θ acting
on $\mathbb{C}[\tau] = \{\phi : V(\tau) \rightarrow \mathbb{C}\}$
each having the form

$$\Theta\phi(u) = \sum_{u \in V(\tau)} c_\Theta(\vec{uv}) [\phi(u) - \phi(v)]$$

(1) Beltrami-Laplace operator Δ

$$c_\Delta(\vec{uv}) = \frac{1}{2} \left[\tan \theta_n(\vec{uv}) + \tan \theta_s(\vec{uv}) \right]$$

(2) discrete Kähler operator \mathcal{D}

$$c_{\mathcal{D}}(\vec{uv}) = \frac{1}{2} \left[\frac{\tan \theta_n(\vec{uv}) + i}{R_n^2} + \frac{\tan \theta_s(\vec{uv}) - i}{R_s^2} \right]$$

(3) conformal Laplacian Δ_{conf}

$$c_{\Delta_{\text{conf}}}(\vec{uv}) = \tan \theta(\vec{uv})$$

def a Delaunay triangulation T is isoradial if all circumradii are equal; $R_f = R_{cr} > 0 \quad \forall f \in F(T)$

in this case $\Delta = \mathcal{D} = \Delta_{\text{conf}}$
 (and $R_{cr}=1$) "critical laplacian" Δ_{cr}

Green's function Δ_{cr}^{-1} (studied by R. Kenyon '02)
 characterised by properties

$$(1) \Delta_{cr} \Delta_{cr}^{-1} = \mathbb{1} \quad \begin{array}{l} \text{(means that} \\ v \mapsto [\Delta_{cr}^{-1}]_{uv} \end{array}$$

is harmonic

$$(2) [\Delta_{cr}^{-1}]_{uv} = O(\log |z_u - z_v|)$$

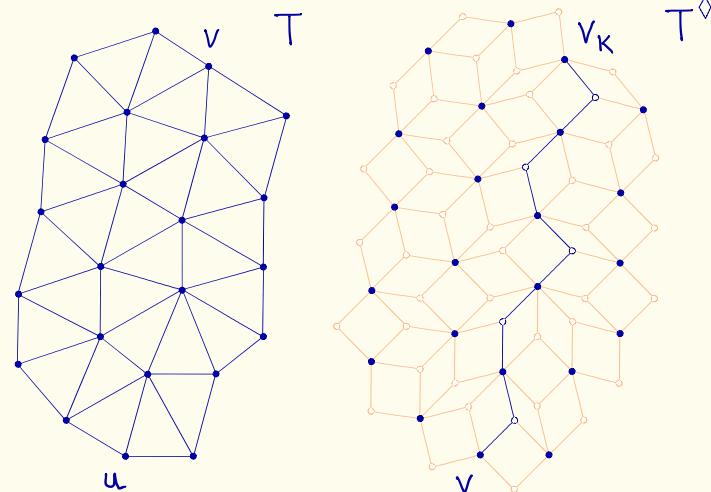
for $u, v \in V(T), |z_u - z_v| \gg 0$

$$(3) [\Delta_{cr}^{-1}]_{u,u} = 0$$

by [Kenyon 2002]

$$(1) [\Delta_{cr}^{-1}]_{uv} = \frac{1}{2\pi} \operatorname{Re} \int_0^1 \frac{dt}{t} (E_{\underline{\theta}}(-t) - 1)$$

$$\text{where } E_{\underline{\theta}}(z) = \prod_{j=1}^k \frac{z + e^{i\theta_j}}{z - e^{i\theta_j}}$$



$v_0 = u$ and $v_K = v$

$$e^{i\theta_j} = z(v_j) - z(v_{j-1})$$

rhombic graph

v_0, \dots, v_K path

(2) Asymptotics $[\Delta_{\text{cr}}^{-1}]_{uv} =$

$$-\frac{1}{2\pi} \left(\log 2D + \gamma_{\text{Euler}} + \frac{\operatorname{re}[P_3(u,v)]}{6D^3} + O(D^{-4}) \right)$$

here $D = |z(u) - z(v)|$ and $|P_3(u,v)| \leq 3D$

for $\epsilon_1, \dots, \epsilon_m > 0$ small
T remains Delaunay
(for all values $0 < l \leq \infty$)

↑ in general not isoradial

Perturbations:

Fix an isoradial triangulation T_{cr}
deform embedding; for $v \in V(T_{\text{cr}})$

But keep incidence relations !
(i.e. same vertex,
edge, face sets)

$$z(v) \mapsto z(v) + \epsilon_1 l F_1(z(v)/l) + \dots + \epsilon_m l F_m(z(v)/l)$$

$\epsilon_1, \dots, \epsilon_m > 0$ deformation parameters

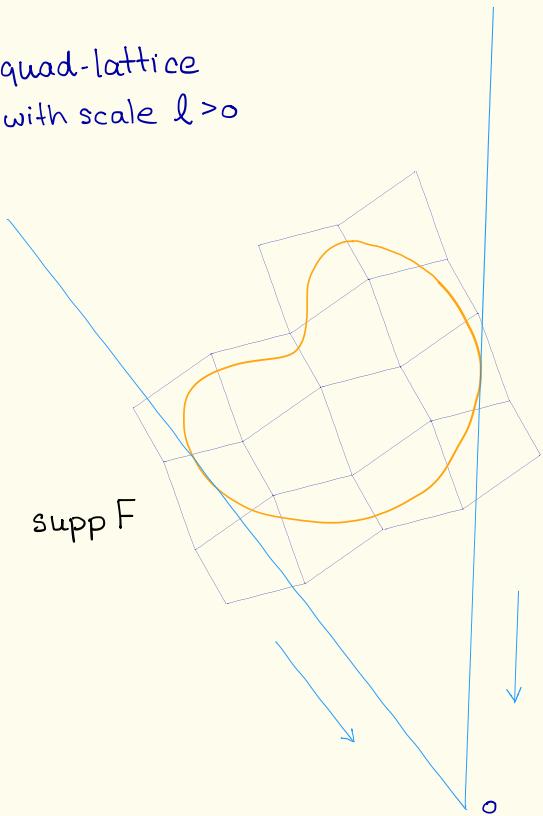
$l > 0$ scaling parameter

$F_1, \dots, F_m : \mathbb{C} \rightarrow \mathbb{C}$ smooth, compact supports

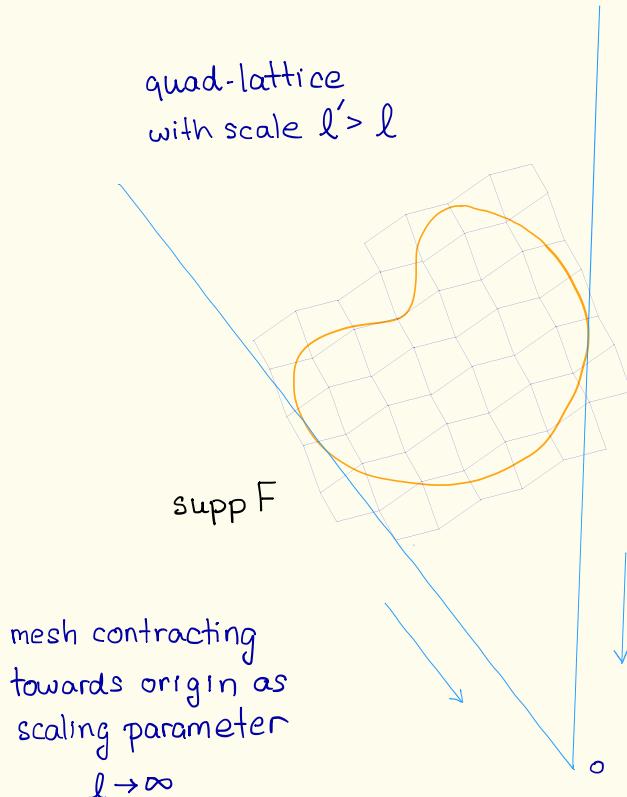
$\operatorname{supp} F_i \cap \operatorname{supp} F_j = \emptyset$ for all i, j

$m = 1$, i.e. one zone

quad-lattice
with scale $l > 0$



quad-lattice
with scale $l' > l$



mesh contracting
towards origin as
scaling parameter
 $l \rightarrow \infty$

set $\delta \Theta = \Theta - \Delta_{\text{cr}}$ and then formally expand

$$\log \det \Theta = \log \det \Delta_{\text{cr}} + \text{tr} [\delta \Theta \Delta_{\text{cr}}^{-1}] - \frac{1}{2} \text{tr} [\delta \Theta \Delta_{\text{cr}}^{-1}]^2 + \dots$$

Thm [F.David + J.Scott] let $m=2$ (i.e. bilocal perturbation)

and $\Theta = \Delta$ or \mathcal{D} then the $\epsilon_1 \epsilon_2$ -crossterm of $\text{tr} [\delta \Theta \Delta_{\text{cr}}^{-1}]^2$ is

$$-\frac{2}{2\pi^2} \sum_{\substack{\text{triangles} \\ (X_1, X_2)}} A(X_1) A(X_2) \text{re} \left[\frac{\bar{\partial} F_1(x_1/\ell) \bar{\partial} F_2(x_2/\ell)}{(x_1 - x_2)^4} \right] + O(D_\ell^{-5})$$

(x_1, x_2)

where x_i is the center of X_i

$$D_\ell = \text{dist}(\Omega_1(\ell), \Omega_2(\ell))$$

$$\text{and } \Omega_i(\ell) = \{z \in \mathbb{C} \mid z/\ell \in \text{supp } F_i\}$$

connections with CFT :

heuristic	discrete world	continuous world
	moduli-space of Delaunay triangulations $\subset \text{Gr}_{\mathbb{C}}^{(2, \infty)} / \infty\text{-torus}$ $\stackrel{\psi}{\longleftarrow}$ isotradual triangulations	\longleftrightarrow moduli-space of metrics $g = (g_{\mu\nu})$ on the plane
		ψ \longleftrightarrow flat metrics
		$\delta g = J^+ g + g J^-$ $J = \text{jac.}(F)$ $F: M \rightarrow M$
<p>general riemannian manifold (M, g) (g flat) equipped with Beltrami-Laplace Δ_g (self-adjoint) with respect to pairing $\langle \phi, \psi \rangle = \int_M dx g ^{1/2} \phi(x) \psi(x)$ by analogy with f.d. Gaussian integrals "define"</p> $\det^{-1/2} \Delta_g = \int D[\phi] \exp -\frac{1}{2} \langle \phi, \Delta_g \phi \rangle =: Z_g$ <p>i.e. $-\frac{1}{2} \log \det \Delta_g = \log Z_g$</p>	<p>perturb $g \mapsto g + \epsilon \delta g$ do functional Taylor expansion use <u>moment formulae</u> of the <u>formal gaussian integral</u> to get</p> $\log Z_{g+\epsilon \delta g} = \log Z_g + \epsilon \log Z_g^{(1)} + \epsilon^2 \log Z_g^{(2)} + \dots$ <p>by CFT, $M = \mathbb{C}$</p> $Z_g^{(2)} = \frac{-2}{4\pi^2} \iint_{\mathbb{C} \times \mathbb{C}} du dv \, re \left\{ \frac{\bar{\partial} F(u) \bar{\partial} F(v)}{(u-v)^4} \right\}$	

opening remarks :

this is joint work with François David at IPhT, Saclay
project is born out of the heuristic which sees
infinite triangulation of the plane (specifically Delaunay
triangulations) as discretisations of continuous metrics
on the plane: Just as a metric g determines
a Laplace-Beltrami operator Δ_g , likewise a Delaunay
triangulation carries a Laplace-Beltrami operator. And
just as one can formalise taking the determinant
 $\det \Delta_g$ of a continuous LB-operator, a similar
analysis can be used for triangulations.

side-step
by looking
at the
difference.

Avoid explaining the mathematical scaffolding needed
to make sense of determinants of ∞ -dim'l operators.